

# Monomer-dimer model and Neumann GFF

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# The dimer model

## Definition

$G$  = bipartite finite graph, planar

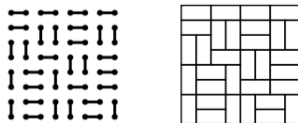
**Dimer configuration** = perfect matching on  $G$ :

each vertex incident to one edge

**Dimer model**: uniformly chosen configuration

More generally, weight  $w_e$  on each edge,

$$\mathbb{P}(\mathbf{m}) \propto \prod_{e \in \mathbf{m}} w_e.$$



On square lattice, equivalent to domino tiling.

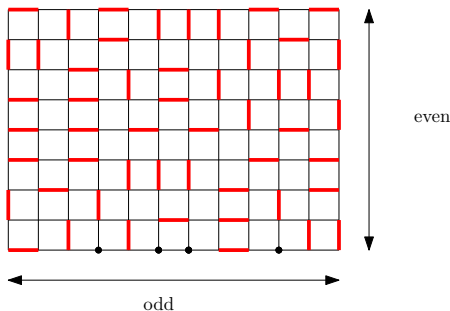
# Monomer-dimer model

Now allow monomers on a part of the boundary, call it  $\partial_m$ .  
Let  $z > 0$  and define

$$\mathbb{P}(\mathbf{m}) \propto z^{\#\text{monomers}}.$$

## Assumptions

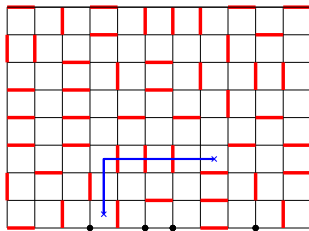
- (1)  $G$  is dimerable (so partition function is  $> 0$ ).
- (2)  $|\partial_m|$  is *odd* (a technical assumption).



Example:

# Height function

Paths between faces avoid monomers:



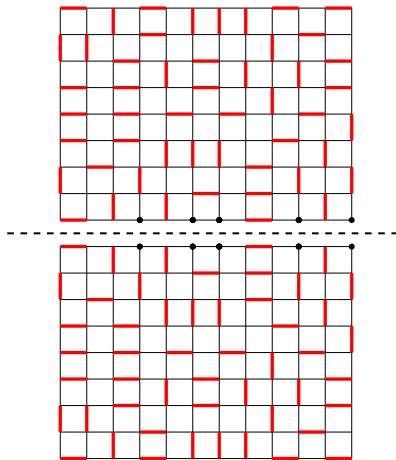
so height function still defined.

## Main question

What is scaling limit of centered height function?  
Conformal invariance?

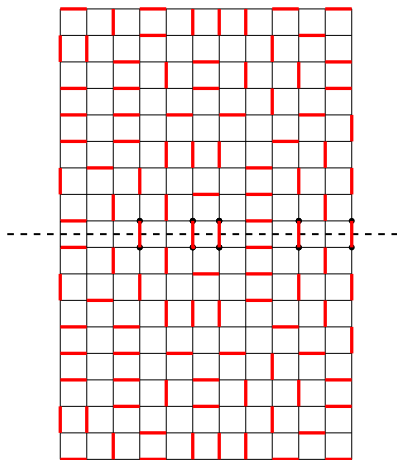
## Reflection symmetry

Suppose  $G \subset \mathbb{H}$  and  $\partial_m \subset \mathbb{R}$ . Then apply reflection:



Get a dimer configuration on  $G_{\text{double}}$ .

# Height function



MD-height function is restriction of dimer height function to  $\mathbb{H}$ .  
Note that height function is then *even*:  $h(z) = h(\bar{z})$ .

# Even height functions

## Conversely

Take dimer model on  $G_{\text{double}}$  with weight  $z > 0$  for  $\mathbb{R}$ -edges  
Condition to be symmetric (or height to be *even*)

Get monomer-dimer model by restricting to  $\mathbb{H}$

## Guessing the scaling limit

In  $\mathbb{C}$ , dimer height function  $\rightarrow$  **full plane GFF** (de Tiliere 2005):

### Full plane GFF

Consider  $\tilde{\mathcal{D}}$  = smooth test functions with compact support and

$$\int_{\mathbb{C}} \rho(z) dz = 0.$$

Scalar product

$$(\rho_1, \rho_2)_{\nabla} = \frac{1}{2\pi} \int_{\mathbb{C}} \nabla \rho_1 \cdot \nabla \rho_2,$$

$\tilde{H}$  = completion of  $\tilde{\mathcal{D}}$  under  $(\cdot, \cdot)_{\nabla}$ .  $\tilde{f}_n$  = orthonormal basis.

$$h_{\mathbb{C}} = \sum_n X_n \tilde{f}_n.$$

Can integrate  $\tilde{h}$  against fixed  $\rho \in \bar{\mathcal{D}}$ : defined up to constant.



# Even/odd decomposition

## Question

What is a (full plane) GFF conditioned to be even?

Any  $\rho \in \tilde{\mathcal{D}}$  can be written uniquely as

$$\rho = \rho_{\text{odd}} + \rho_{\text{even}}$$

where  $\rho_{\text{odd}}/\rho_{\text{even}}$  **Dirichlet/Neumann** boundary conditions.

Moreover

$$(\rho_{\text{odd}}, \rho_{\text{even}})_{\nabla} = 0$$

so

$$H = H_{\text{odd}} \oplus H_{\text{even}}.$$

and

$$h = h_{\text{odd}} + h_{\text{even}}$$

where  $h_{\text{odd}}, h_{\text{even}}$  are **independent** Dirichlet/Neumann GFF.

# Conjecture and main result

Hence “ $h_{\mathbb{C}}$  conditioned to be even”: simply a **Neumann GFF**.

## Conjecture:

The centered monomer-dimer height function on  $D$  converges to a **GFF with Neumann boundary conditions** on  $\partial_m$  and Dirichlet boundary conditions on  $\partial D \setminus \partial_m$ .

## Theorem (B.-Lis-Qian, 2019+)

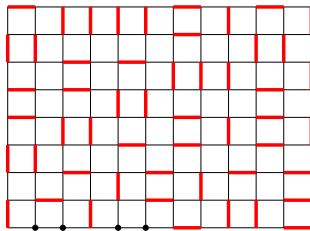
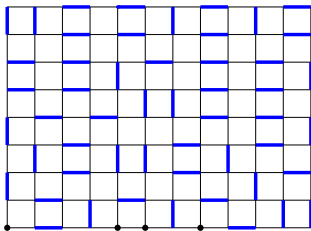
When  $D_n \uparrow \mathbb{H}$  there is a local (inf. volume) limit. Furthermore, in the scaling limit, the height function converges to Neumann GFF.

Remark: also true on infinite strips.

Note: first time the limit doesn't have Dirichlet b.c.

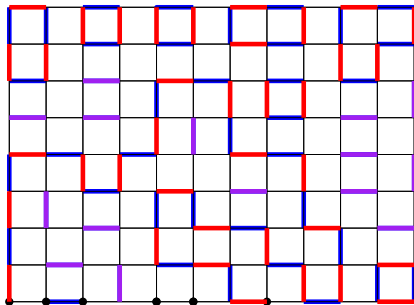
## Double monomer-dimer model

Superposition of two independent realisations of monomer-dimer model:



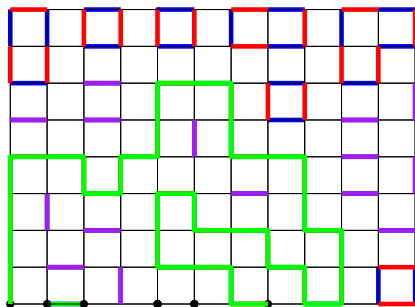
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# Double monomer-dimer model

Superposition of two independent realisations of monomer-dimer model:



Get a collection of Green arcs connecting  $\partial_m$  to  $\partial_m$ .

# Double monomer-dimer model

Question:

In the scaling limit, what is the law of these arcs?

Conjecture

Converges to  $\text{ALE}_4$  aka  $A_{-\lambda,\lambda}$  (cf. Aru–Lupu–Sepulveda)



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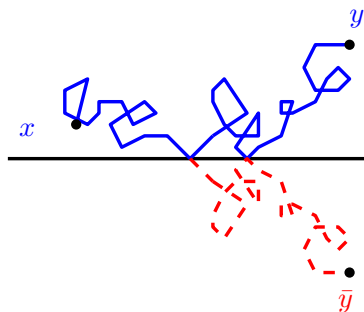
Indeed,  $\text{ALE}$  = boundary touching level lines of Neumann GFF  
(Qian–Werner, CMP 2018).

# More about Neumann GFF

Dirichlet GFF: “pointwise correlation” = Dirichlet Green function

$$\mathbb{E}[h^{\text{Dir}}(x)h^{\text{Dir}}(y)] = G^{\text{Dir}}(x,y) = \pi \int_0^\infty p_t^{\text{Dir}}(x,y)dt$$

In  $\mathbb{H}$  (more integrable):



$$p_t^{\text{Dir}}(x,y) = p_t^{\mathbb{C}}(x,y) - p_t^{\mathbb{C}}(x,\bar{y})$$

So

$$\begin{aligned} G^{\text{Dir}}(x,y) &= -\log|x-y| + \log|x-\bar{y}| \\ &= \log\left|\frac{x-\bar{y}}{x-y}\right|. \end{aligned}$$

# Correlation of Neumann GFF in $\mathbb{H}$

Neumann GFF == free boundary conditions:  
e.g. scaling limit of DGFF with free b.c.

$$\mathbb{E}[h^{\text{Neu}}(x)h^{\text{Neu}}(y)] = G^{\text{Neu}}(x, y) = \int_0^\infty p_t^{\text{Neu}}(x, y) dt$$

In  $\mathbb{H}$ :

$$p_t^{\text{Neu}}(x, y) = p_t^{\mathbb{C}}(x, y) + p_t^{\mathbb{C}}(x, \bar{y})$$

So

$$\begin{aligned} G^{\text{Neu}}(x, y) &= -\log|x - y| - \log|x - \bar{y}| \\ &= -\log|(x - \bar{y})(x - y)|. \end{aligned}$$



Only defined up to constant so  $G^{\text{Neu}}$  **not unique**. Instead:

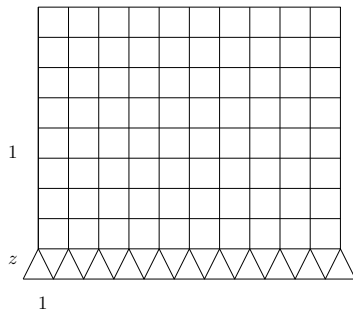
$$\mathbb{E}[(h(a_i) - h(b_i))(h(a_j) - h(b_j))] = \log \left| \frac{(a_i - a_j)(b_i - b_j)(\bar{a}_i - a_j)(\bar{b}_i - b_j)}{(a_i - b_j)(b_i - a_j)(\bar{a}_i - b_j)(\bar{b}_i - a_j)} \right|$$



# Sketch of proof of main result

## Bijection to non-bipartite dimer

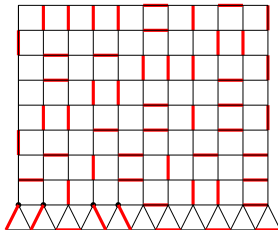
Giuliani, Jauslin, Lieb: **Pfaffian formula** for correlations.  
In fact, bijection dimer model



# Sketch of proof of main result

## Bijection to non-bipartite dimer

Giuliani, Jauslin, Lieb: **Pfaffian formula** for correlations.  
In fact, bijection dimer model



## Lemma

If  $|\partial_m|$  odd, and  $\#$  monomers even, then unique way to associate dimer configuration on augmented graph.

# Kasteleyn theory

## Problem:

Graph becomes non-bipartite.

## Kasteleyn theory

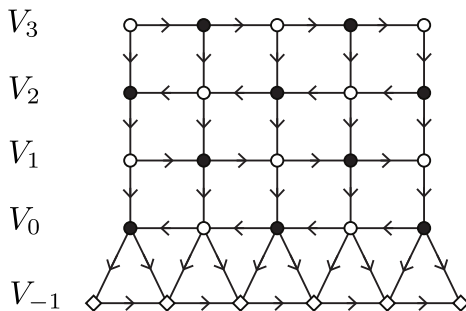
**Kasteleyn orientation:** going cclw on every face, **odd number** of clw arrows

**Gauge transform:** weight of every edge coming of a vertex  $v \rightarrow \times \lambda_v$ , with  $\lambda_v \in \mathbb{C}, |\lambda_v| = 1$ .

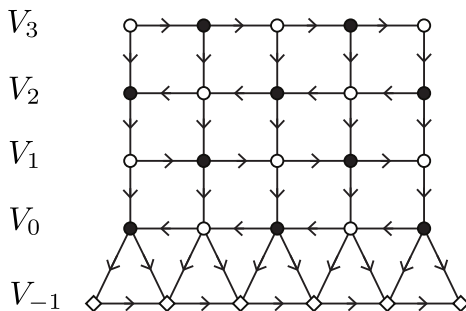
**Kasteleyn matrix:**  $K(u, v)$  = signed weight of edge  $(u, v)$  (so  $K$  antisymmetric).

Then correlations are given by  $Pf(K^{-1})$ .

# Kasteleyn orientation



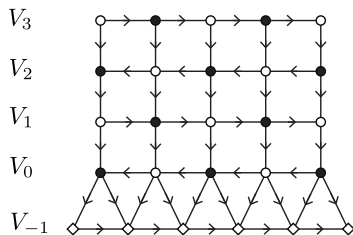
# Gauge transform



Even rows multiplied by  $i$ :

$\implies$  each edge  $e \in E_{\text{even}}$  multiplied by  $-1$ ;  
and each vertical edge has weight  $i$ .

# Kasteleyn matrix in bulk



**Kenyon:** consider  $D = K^* K$ .

$K = \text{n.n.}$  so  $D$  nonzero only from  $W \rightarrow W, B \rightarrow B$ .

Diagonal contributions vanish

So really  $W_0 \rightarrow W_0, \dots B_1 \rightarrow B_1$ .

Then  $D = \text{discrete Laplacian on each four sublattices.}$

Temperleyan b.c.:  $\implies D$  has Dirichlet b.c. on  $B_0$ .

## Scaling limit in bipartite setup

From the relation  $D = K^*K$  we get

$$D^{-1} = K^{-1}(K^*)^{-1}$$

and so

$$K^{-1} = D^{-1}K^*.$$

Moreover  $D^{-1} = \text{Green function}$  and  $K^* = \text{discrete derivative}$ .  
By Kasteleyn's theorem and since graph is bipartite,

$$\mathbb{P}(\mathbf{e}_1, \dots, \mathbf{e}_n \in \mathbf{m}) = \det(K^{-1}(\mathbf{e}_i, \mathbf{e}_j)_{1 \leq i, j \leq n})$$

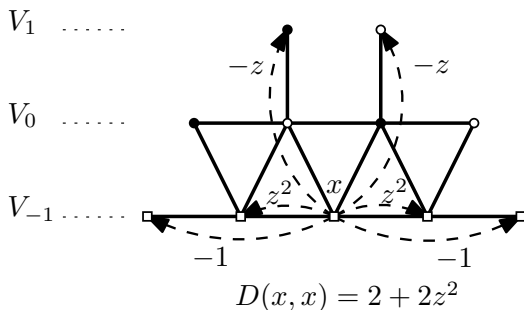
so leads to scaling limit for  $n$ -point correlation function.





# Kasteleyn matrix near monomers

At rows 1, 0, -1 the above analysis breaks down:



but  $-\sum_{y \sim x} D(x, y) = 2 - 2z^2 + 2z$ .

Diagonal terms still vanish.

# Dealing with negative rates

Let  $P(x, y) = -D(x, y)/D(x, x)$  : can be signed, don't sum to 1...

## Question

Can we still make sense of Green function?

If  $\|P\| < 1$  then

$$D^{-1}(x, y) = \frac{1}{D(y, y)} \sum_{\text{path } \pi: x \rightarrow y} w(\pi)$$

where

$$w(\pi) = \prod_{(u, v) \in \pi} P(u, v).$$

# Monomer excursions

Paths still restricted even  $\rightarrow$  even and odd  $\rightarrow$  odd rows.

Decompose in **excursions** into  $V_{-1}$  or  $V_0$ .

Eg odd case (harder): let  $u, v \in V_1$ .

Associated vertices  $u_-, u_+$  and  $v_-, v_+$  in  $V_{-1}$ , two steps away.

Let  $u_\bullet \in \{u_-, u_+\}$  and  $v_\bullet \in \{v_-, v_+\}$ .

Let  $\pi : u_\bullet \rightarrow v_\bullet$ . Parity fixed so

$$w(\pi) = (-1)^{v_\bullet - u_\bullet} \prod_{(x,y) \in \pi} p_{x,y}$$

where

$$p_{i,i\pm 1} = \frac{z^2}{2 + 2z^2} =: 1/2 - p, \quad p_{i,i\pm 2} = \frac{1}{2 + 2z^2} =: p.$$

$\implies$  an honest RW on  $V_{-1} \simeq \mathbb{Z}$  in limit!

## Odd monomer excursions

So

$$\sum_{\pi: u_{\bullet} \rightarrow v_{\bullet}; \pi \subset V_{-1}} w(\pi) = (-1)^{v_{\bullet} - u_{\bullet}} g_{u_{\bullet}, v_{\bullet}}$$

where  $g_{x,y} = 1d$  Green function (with certain b.c.).

Sum over  $u_{\bullet} \in \{u_{-}, u_{+}\}$ ,  $v_{\bullet} \in \{v_{-}, v_{+}\}$ , take local limit  $D_n \uparrow \mathbb{H}$ ,

$$\sum_{\pi: u \rightarrow v; \pi \subset V_{-1}} w(\pi) = C_z (-1)^k (2a_k - a_{k+1} - a_{k-1})$$

where  $a_k =$  **Potential kernel** of 1d walk;  $k = \text{Re}(v - u)$ .

# The miracle

## Lemma

$(-1)^k \Delta a_k \geq 0$  for all  $k \in \mathbb{Z}$

Moreover  $\sum_{k \in \mathbb{Z}} C_z (-1)^k (2a_k - a_{k+1} - a_{k-1}) = 1$ .

Gives an **effective random walk** on  $V_1 \cup V_3 \cup \dots \simeq \mathbb{H}$  !

## Lemma

$(-1)^k \Delta a_k$  decays exp. fast as  $k \rightarrow \infty$ .

So: reflection on boundary with jumps, but exponential tails!

# Proof of oscillations

$a_k$  solves a recurrence relation of order **four**.

Also by general theory [e.g. Lawler–Limic]:

$$a_x \sim \frac{|x|}{\sigma^2} \text{ as } |x| \rightarrow \infty$$

Hence

$$a_x = \frac{|x|}{\sigma^2} + A + B\gamma^{|x|}$$

where

$$1 = (1/2 - p)(\gamma + \gamma^{-1}) + p(\gamma^2 + \gamma^{-2}).$$

Hence let  $s = \gamma + \gamma^{-1}$

$$1 = (1/2 - p)s + p(s^2 - 2).$$

Can solve  $s$  so  $s = 2$  or  $s = -1 - 1/(2p)$ .

This implies  $\gamma \in (-1, 0)$  so oscillations.

## Towards scaling limit

Notice that  $D^{-1}(u, v)$  **not restricted** to  $B \rightarrow B, W \rightarrow W$ :

However paths must go through boundary !

Eg:  $e = (w, b); e' = (w', b')$

$$\begin{aligned}\mathbb{P}(e, e' \in \mathbf{m}) &= \text{Pf} \begin{pmatrix} 0 & K^{-1}(w, b) & K^{-1}(w, w') & K^{-1}(w, b') \\ & 0 & K^{-1}(b, w') & K^{-1}(b, b') \\ & & 0 & K^{-1}(w', b') \\ & & & 0 \end{pmatrix} \\ &= \overbrace{K^{-1}(w, b)}^{\mathbb{P}(e \in \mathbf{m})} \overbrace{K^{-1}(w', b')}^{\mathbb{P}(e' \in \mathbf{m})} + K^{-1}(b, w')K^{-1}(w, b') \\ &\quad - K^{-1}(w, w')K^{-1}(b, b')\end{aligned}$$

so

$$\text{Cov}(1_{e \in \mathbf{m}}; 1_{e' \in \mathbf{m}}) = K^{-1}(b, w')K^{-1}(w, b') - K^{-1}(w, w')K^{-1}(b, b')$$

Leads to scaling limit eventually...!

