Tautomerization in 2,7,12,17-Tetraphenylporphycene and 9-Amino-2,7,12,17-tetraphenylporphycene: Influence of Asymmetry on the Direction of the Transition Moment


Abstract: Femtosecond transient absorption anisotropy studies have been performed for two porphycenes of different symmetry. In 2,7,12,17-tetraphenylporphycene, the chemical identity of two trans forms implies a change in the S₀–S₁ transition-moment direction upon tautomerization. Exploiting this phenomenon, the rates of double hydrogen transfer in both the S₀ and S₁ states (1.4 × 10¹¹ s⁻¹ and 2.7 × 10¹¹ s⁻¹, respectively) have been determined by performing time-resolved anisotropy studies. In the asymmetric 9-amino-2,7,12,17-tetraphenylporphycene, tautomerization occurs with a similar rate in the ground state. In the S₁ state, the reaction is hindered in its vibrationally relaxed form, but the excitation spectra suggest that it may occur for an unrelaxed molecule. Unlike all porphycenes that have been studied so far, 9-amino-2,7,12,17-tetraphenylporphycene does not reveal significant changes in anisotropy owing to intramolecular double hydrogen transfer; rather, the transition-moment directions are similar in the two tautomeric forms. Analysis of the molecular orbitals allows for an explanation of the “locking” of the transition moments: it is due to a large splitting of the two HOMO orbitals, which retain the order of their energies in the two tautomers.

Keywords: hydrogen transfer · porphycenes · tautomerism · transition moments

Introduction

The spectroscopic- and photophysical properties of porphycenes, which are structural isomers of porphyrins, make these compounds very attractive for applications that involve the use of light.[31] Intense absorption in the red-light spectroscopic region and good yields of singlet-oxygen generation suggest that porphycenes may be efficient in photodynamic therapy and in photodiagnostics. This prediction has been positively tested for various derivatives, both in vitro and in vivo.[2] Another widely explored area is related to the mechanism of intramolecular tautomerism.[3–22] Being of fundamental importance for understanding the multidimensional character of tautomerization coordinates, mode-selectivity, tunneling, and the cooperativity between the motions of the two protons, such studies can also lead to applications of porphycenes, for example, as ultrafast switches or as information-storage media.[23,24]

Our previous investigations of intramolecular double hydrogen transfer in porphycenes and porphyrins have resulted in the development of techniques for the determination of tautomerization rates.[25–29] These procedures are based on the use of polarized spectroscopy, both in steady-state and time-resolved regimes. The former regime has allowed the obtaining of rates for the reaction that occurs in the lowest excited singlet state,[28] as well as the detection of tautomerization in single molecules.[30,31] The latter regime, especially femtosecond time-resolved absorption,[25–27] enabled simultaneous characterization of both the ground- and excited-state reaction kinetics. For a series of five differently alkyl-substituted porphycenes, the values of the ground-state rate constants vary in the range of less than 10¹⁰ s⁻¹ to more than 10¹¹ s⁻¹.[26] The S₁ values, which are several times lower in each compound, also span a similar range. These huge dif-
ferences are indicative of double proton tunneling, which is important, even at room temperature.

Herein, we use spectroscopy with polarized light to study tautomerism in aryl-substituted porphycenes, a class of compounds for which their tautomerism has not been investigated before. We compare 2,7,12,17-tetraphenylporphycene (TPPo) with 9-amino-2,7,12,17-tetraphenylporphycene (9-ATPPo, Scheme 1). Previous studies of 9-ATPPo revealed quite unique properties of this molecule. Dual fluorescence was observed and the energy gap between the two bands exceeded 1500 cm\(^{-1}\). The emission decays were biexponential, thereby clearly showing the presence of two different forms, which were assigned to trans tautomers. Thus, the amino substituent selectively and substantially lowers the electronic transition energy of only one tautomer; such an effect has not been observed for TPPo that is substituted at the 9-position with acetoxy- or nitro groups.

The most-intriguing result of this previous work was the suggestion, based on calculations, that trans–trans tautomerization in 9-ATPPo does not lead to a change in the \(S_0–S_1\) and \(S_0–S_2\) transition-moment (TM) directions. Such an effect, which has never been observed previously for porphycenes, prompted us to perform detailed investigations of this molecule and of its symmetrical counterpart, TPPo. We demonstrate that, whilst the TM directions do rotate as a result of tautomerization in TPPo, this is not the case for 9-ATPPo. Strictly speaking, a small change in the TM direction is observed by transient absorption spectroscopy, thus demonstrating that tautomerization in 9-ATPPo occurs with a considerable rate. We discuss the reasons for the presence (or lack of) TM changes in asymmetrical porphycenes on the basis of orbital-energy patterns and its changes during double-hydrogen-transfer reactions.

**Results and Discussion**

**Structures of low-energy tautomers**: Figure 1 presents the absorption spectra of both porphycenes. A dramatic change in the absorption pattern for 9-ATPPo has been explained previously by the presence of two trans tautomeric forms (Figure 2) that differ in their \(S_0–S_1\) transition energies. Other possible tautomeric species, four cis-type tautomers, for which the inner protons reside on the adjacent pyrrole rings, were computed at significantly higher energies. The presence of two forms of 9-ATPPo is clearly manifested in its emission spectrum: dual fluorescence is observed with a separation of 84 nm (about 1500 cm\(^{-1}\)) between the maxima. TD-DFT calculations of the absorption...
spectrum predict a separation of 900 cm\(^{-1}\), mostly owing to the difference between the \(S_1\) energies. In the ground state, both forms are calculated to be almost isoenergetic, with a slightly lower predicted energy for the \textit{trans}-1 species (0.10 and 0.14 kcal mol\(^{-1}\) without- and with the zero-point-energy included, respectively). The computed near-degeneracy of both forms is confirmed experimentally by recording the absorption spectrum as a function of temperature (Figure 3).

Besides a narrowing of the bands at lower temperatures, no features are observed that would suggest a significant change in the population ratio of the two forms. Analogous behavior has been reported previously for 9-acetoxy-2,7,12,17-tetra-\(n\)-propylporphycene (9-AcTPrPo); in that case, this behavior was interpreted as evidence for the ground-state tunneling of two hydrogen atoms between two non-equivalent forms.[38]

Figure 2 also shows the calculated \(S_0\text{-}S_1\) and \(S_0\text{-}S_2\) TM directions for the symmetrical- and asymmetrical porphycene derivatives. The symmetrical double-minimum character of the potential-energy surface in TPPo dictates that each TM changes its direction upon tautomerization: It is rotated by an angle that is twice as large as that at which it forms with the vertical (or horizontal) axis. The situation is completely different in 9-ATTPo: Now, the calculated TM directions are quite similar in the two \textit{trans} forms; the TMs are rotated by 12° and 18° for the \(S_0\text{-}S_1\) and \(S_0\text{-}S_2\) transitions, respectively. Such behavior, which was confirmed experimentally (see below), is contrary to the case of 9-AcTPrPo, another asymmetrical derivative for which the TMs have been shown to rotate upon \textit{trans\textendash}-\textit{trans} conversion, in a similar fashion to that of its symmetrical counterpart, 2,7,12,17-tetra-\(n\)-propylporphycene (TPrPo).[25,38] We will explain this difference by analysis of the molecular orbitals that are involved in the transitions.

**Tautomerization in TPPo:** Figure 4 shows the temporal changes in the transient absorption (TA) anisotropy at two different wavelengths, which correspond to bleaching in the regions of the \(S_0\text{-}S_1\) and \(S_0\text{-}S_2\) electronic transitions (cf. Figure 1). The excitation wavelength corresponded to the \(S_0\text{-}S_1\) transition. The two curves show similar temporal characteristics but opposite trends. The anisotropy probed in the excitation region decays from an initial value of 0.4 to about 0.1, whereas the kinetic curve recorded in the \(S_0\text{-}S_2\) region rises monotonously, from a value of −0.2 to about 0.1. These results are in perfect agreement with the predictions based on calculations, which yield a value of 89° for the angle that is formed by the \(S_0\text{-}S_1\) and \(S_0\text{-}S_2\) transition moments (Figure 2). Notably, though, the absolute positions of the TM directions are probably calculated with an error, because the final value of the anisotropy (0.1), corresponding to the plateau region in Figure 4, suggests that the TM directions are tilted by 45° from both the vertical- and horizontal directions in Figure 2. We neglect herein the possible influence of rotational depolarization because it takes much longer than 10 ps, which is the time span of the experiment shown in Figure 4.

Assuming that the only origin of the change in anisotropy is the rotation of the TM direction, caused by reversible \textit{trans\textendash}-\textit{trans} double hydrogen transfer, one can determine the rates of tautomerization in both the \(S_0\) and \(S_1\) states (\(k^0\text{PT}\) and \(k^{1\text{PT}}\), respectively), according to Equation (1 ), where \(r(t)\) is the measured anisotropy of the TA signal: \(r(t) = (\Delta OD(t)_{\text{par}} - \Delta OD(t)_{\text{perp}})/\Delta OD(t)_{\text{par}} + 2\Delta OD(t)_{\text{perp}}\); the subscripts indicate parallel or perpendicular polarization of the probe beam with respect to that of the pump.

\[
    r(t) = f[(r_0^0 - r_0^1)e^{-2k^0_{\text{PT}}t} + r_0^1] + (1 - f)[(r_0^1 - r_\infty^1)e^{-2k^{1\text{PT}}t} + r_\infty^1] \\
    \quad (1 - f)[(r_0^1 - r_\infty^1)e^{-2k^{1\text{PT}}t} + r_\infty^1]
\]

The contributions to anisotropy that originate from the ground- and excited states are labeled as \(f\) and \(1 - f\), respectively; \(r_0^1\) and \(r_\infty^1\) are the asymptotic values owing to the \(S_0\) and \(S_1\) states at an infinitely long delay. The anisotropies at the moment of excitation are denoted as \(r_0^1\) and \(r_\infty^1\).

By using this procedure, the rate constants for tautomerization have been determined for TPPo. In the ground state, the reaction occurs within 0.7(±0.4) ps, whilst it is about 5-fold slower in the \(S_0\) state (3.7(±1.3) ps). Within experimental error, these values are the same as those for TPrPo (0.8-
A C H T U N G T R E N N U N G

Similar values of the rate constants suggest that the strengths of the intramolecular hydrogen bonds are not much different between the three compounds. This result is confirmed by similar values of the $^1H$ NMR shifts: $\delta = 3.5$, 3.04, and 3.58 ppm for TPPo,\[34]\ TPrPo,\[39] and TiBPo,\[26] respectively. Another measure of the hydrogen-bond strength is the distance between the H-bonded nitrogen atoms. The X-ray data give values of 2.615 Å for TPrPo\[39] and 2.63 Å for TiBPo;\[26] no X-ray data are available for TPPo. The calculations predict a N–N distance of 2.646 Å, which is slightly shorter than the value of 2.655 Å that was obtained by using the same model and basis set for the parent, unsubstituted porphycene (for which the X-ray data yield 2.63 Å\[40]). These data are in agreement with the observation that the rate of tautomerization in the parent porphycene is about two-times slower than in the 2,7,12,17-substituted derivatives. These results also show that the alkyl- and aryl substituents influence the geometry of the inner cavity, hydrogen-bond strength, and the rate of tautomerization to the same degree.

**Transient absorption studies of 9-ATPPO:** Time-resolved absorption spectra (Figure 5) were recorded at different excitation wavelengths: 400, 660, and 700 nm. These latter two wavelengths correspond to the $S_0$–$S_1$ transition in the trans-2 and trans-1 forms, respectively, whilst both tautomers are supposedly excited at 400 nm (see Figure 1).

All of these spectra have several common features: bleaching bands are observed, which correspond to the Soret- and Q transitions, along with an excited-state absorption (ESA) peak at around 520 nm, which overlaps with the bleaching. This overlap is so strong that the shortest-wave absorption maximum of the Q bands (590 nm) only appears as a dip in the ESA band. The temporal evolution of the spectra at these three excitation wavelengths is very similar: within the first 50 ps, the absorption changes very weakly below 590 nm and increases slightly in the range 610–640 nm with a characteristic time in the range 5–12 ps, depending on the excitation wavelength (decay-associated difference spectra that result from a global analysis of the TA data; see the Supporting Information, Figure S1). More-significant changes can only be seen in the bleaching band: the negative band at around 715 nm decays by approximately 30% with a decay time of 0.3–0.7 ps. The characteristic times of both processes increase with the excitation energy; therefore we ascribe them to intramolecular vibrational redistribution and energy-relaxation.

Over a longer timescale, from 50 to 500 ps, the transient signals decay proportionally at all wavelengths. This slow evolution most-likely exclusively results from depopulation of the excited state and repopulation of the ground state.

Figure 6 shows the kinetics of anisotropy at four selected wavelengths for each of three different excitation energies. Contrary to other porphycenes, the anisotropy does not change significantly over a short timescale (for longer delays of up to 500 ps, decays of anisotropy that correspond to very slow rotational diffusion can be detected). Very small changes in anisotropy within the first 10 ps can only be seen for selected combinations of excitation-/probe wavelengths.
The recorded data indicate the absence of a fast process that would lead to a significant change in the direction of the transition dipole moment. Therefore, either the tautomerization is not accompanied by a significant change in the transition dipole moments or the tautomerization does not occur on a timescale shorter than hundreds of picoseconds. This latter hypothesis seems to be very unlikely, because the geometry of the inner cavity is very similar to that of alkyl-substituted porphycenes, for which fast tautomerization has been confirmed. On the other hand, the experimental results are consistent with theoretical predictions that, upon tautomerization in 9-ATPPo, the TMs are rotated by a very small angle. Therefore, we focused more closely on the anisotropy behavior to check whether such minor changes can be reliably analyzed.

Figure 7 shows the anisotropy for two different delays from the moment of excitation. The anisotropy remains exactly the same in the region of strong excited-state absorption (450–550 nm), which confirms the absence of tautomerization in the S1 state on the picosecond timescale. However, the anisotropy clearly changes in the region of the bleaching of the Q-absorption bands (600–700 nm). The anisotropy decay in this range can be reproduced by a double exponential function with decay times of τ1 = 550(±50) fs and, with a much smaller contribution, τ2 = 8(±3) ps. This former value is close to the times of ground-state tautomerization for TPPo, TPrPo, and TtBPo. The slower decay may correspond to the vibrational cooling and thermal equilibration of the initially excited S1 state of trans-1 (for the parent porphycene, a characteristic time of 6.2(±0.5) ps has been reported to account for this process).

Similar values of the anisotropy decay time have been found for other excitation-emission-wavelength combinations. For instance, probing at 740 nm with excitation at 700 nm yielded a time of 680(±50) fs. Importantly, it is practically impossible to find a wavelength that exhibits a pure bleaching signal, owing to its strong overlap with the excited-state absorption or stimulated emission. Still, we conclude that these data strongly suggest that the sub-picosecond anisotropy decay is due to ground-state tautomerization, which occurs in 9-ATPPo at a similar rate as in other similarly substituted symmetrical porphycene derivatives.

An even-more-challenging problem is the estimation of the rate of tautomerization for the S1 excited state. If the more-stable trans-1 form is excited, no tautomerization is to be expected, owing to the large energy difference between the S1 states of trans-1 and trans-2, with the latter form located at much-higher energy. On the other hand, if the trans-2 form is excited, tautomerization in the S1 state should be a downhill process. Such a case has been investigated previously for 9-acetoxy-2,7,12,17-tetrapropylporphycene (9-AcTPrPo). This downhill process was determined to occur within 2.6 ps at 293 K. Such a fast rate of the irreversible excited-state reaction could explain the presence of just one fluorescence band, whereas two tautomeric forms could be readily detected in the absorption spectrum. For 9-ATPPo, the situation is different because dual emission is observed. Most importantly, the lifetime of fluorescence from trans-2, above 1 ns, puts an upper limit of 10^9 s^-1 on the rate of tautomerization from the relaxed S1 state. In principle, one should not exclude a higher rate from a vibrationally unrelaxed S1 state, or from higher electronic states. Still, one should expect fast anisotropy changes, as well as a decrease in the emission intensity of the trans-2 form for higher-energy excitations. Such effects were not clearly detected; in particular, slight anisotropy decays, which were observed after excitation at 400 nm, could not be unequivocally interpreted, owing to the presence of internal-conversion processes that also lead to anisotropy changes. In addition, the selective excitation of the trans-2 form is not possible because its absorption overlaps with that of the trans-1 form.

Some indications regarding the extent- and direction of excited-state tautomerization can be extracted from the fluorescence excitation spectra. Figure 8 shows the excitation spectra of the lower- and higher-energy emissions (F1 and F2, respectively) for a solution of 9-ATPPo in n-propanol.
These spectra are very different, which immediately excludes the presence of an excited-state equilibrium between the two tautomers. The same conclusion has been drawn previously, basing on different F1 and F2 fluorescence-decay times and on different excitation spectra in toluene and 2-methyltetrahydrofuran.\(^{[22]}\) Interestingly, whereas the excitation- and absorption spectra are very different for F2, they strongly resemble each other for F1.

Assuming the possibility of excited-state conversion between the two tautomers, the intensity of F1 fluorescence can be expressed as \(I(F_1) = C_{ref} \phi_1(f_1 + 1)\). In this expression, \(C_{ref}\) is a constant that is characteristic of the detection system, \(I_{ex}\) denotes the intensity of excitation, \(f_i\) is the fraction of initially excited molecules of tautomer 1 that end up in the emitting state in the same tautomeric form, \(\sigma_i\) and \(\sigma_2\) are the absorption cross-sections for tautomers 1 and 2, respectively, and \(\phi_{i1}\) is the efficiency of 2→1 excited-state conversion. Finally, \(\phi_{i0}\) is the quantum yield of fluorescence from the vibrationally relaxed S1 level of tautomer 1. One should note that, except for \(\phi_{i0}\), all of the other terms are wavelength-dependent. For the intensity of F2, one obtains \(I(F_2) = C_{ref} \phi_2(f_1 + 1)\). These expressions can be simplified further. For instance, if excited-state tautomerization is the only process that is responsible for not all of the initially excited tautomer 1(2) species yielding F1(F2) fluorescence, then \(f_1 = 1 - \phi_{21}\) and \(f_2 = 1 - \phi_{21}\). Relevant for 9-ATTPo seems to be the case of the “downhill” 2→1 reaction. In such a case, \(\phi_{21} \approx 1\) and the excitation spectrum of F1 should resemble the absorption spectrum. Indeed, this result is observed. However, the detection of two different nanosecond lifetime for F1 and F2 indicates that \(\phi_{21}\) can only be large for molecules of tautomer 2 that are excited high above the S1 origin, but, after relaxation, \(\phi_{21}\) becomes small. Moreover, no significant changes were observed for different delay times and excitation wavelengths in the TA spectra. Evidently, further detailed studies seem to be necessary to prove the dependence of the intensities of the F1 and F2 emissions on the excitation wavelength and on the temperature. Combined with wavelength-dependent emission kinetics, such investigations could provide insight into the tautomerization mechanisms in the different excited electronic states of porphycene. In this context, we recall that at least four different electronic transitions contribute to the absorption in the region above 350 nm.

The origin of “locking” of the transition-moment directions in 9-ATTPo: Both experiment and theory demonstrate that tautomerization in 9-ATTPo does not lead to a significant change in the direction of the transition moment. In this respect, this molecule is different from all other porphycenes that have been investigated to date, including 9-AcTPrPo, which is a derivative with asymmetric potential for tautomerization. To explain the unique behavior of 9-ATTPo, we analyzed the order of the energies of the molecular orbitals that were responsible for the S1→S2 and S0→S2 transitions (Figure 9). It is well-known that porphycenes belong to the class of so-called “negative hard” chromophores, for which the orbital-energy splitting between the two highest-occupied molecular orbitals is small compared to the splitting of the two lowest-unoccupied orbitals (ΔHOMO ≪ ALUMO).\(^{[40]}\) A consequence of such orbital-energy ordering is that two lowest-energy transitions are dominated by configurations that involve the promotion of an electron into the same LUMO orbital. In the parent porphycene and in other symmetric derivatives, the largest contribution comes from HOMO→LUMO excitation for the S0→S1 transition, whereas in the S0→S2 transition, the HOMO→LUMO configuration dominates. The HOMO and HOMO→1 orbitals lie very close together in energy. Looking at their shapes (Figure 9), one realizes that tautomerization converts the HOMO orbital into the HOMO→1 orbital and vice versa. Naturally, the orbital configuration and state energies remain the same, owing to symmetry, but the transition moments are rotated because the spatial distribution of the electrons in the molecule changes as a result of trans→trans conversion.

When the symmetry of the double minimum is slightly perturbed, as in 9-AcTPrPo, neither the orbitals nor their energies remain the same after tautomerization. Upon passing from the trans-1 to trans-2 forms, the energy of the HOMO increases slightly, whereas that of the HOMO→1 is stabilized. The energy shifts are not large and the orbitals retain their shapes. However, the dominant configurations are now different in these two tautomers. For the trans-2 form, the dominant configurations for the two lowest singlet states are the same as in the symmetrical derivatives, but they are reversed in trans-1: for this tautomer, the S0→S1 transition is dominated by the HOMO→LUMO configuration, whereas the S0→S2 transition is dominated by the HOMO→1→LUMO configuration. As a result, the transition moments for both states are rotated upon hydrogen transfer.

For 9-ATTPo, the situation is different. The introduction of a strongly electron-donating group leads to an increase in the energies of both the HOMO→1 and HOMO orbitals,
but one of them is destabilized to a much-larger degree. This effect results in much-larger splitting between the two highest-occupied orbitals: $\Delta \text{HOMO} = 0.41$ and 0.48 eV for the $\text{trans}-1$ and $\text{trans}-2$ forms, respectively; the corresponding values for 9-AcTPPrPo are 0.03 and 0.12 eV. Such an energy pattern in 9-ATPPo ensures that no reversal in the dominant configuration upon tautomerization is to be expected. Indeed, for both tautomeric forms, the same configurations are dominant: HOMO$-1$—LUMO for $S_0$—$S_1$ and HOMO$-1$—LUMO for $S_0$—$S_2$. Because the orbitals are similar in both forms, no significant rotation of the TMs occurs. The calculated angles of 12° and 18° reflect slight changes in the orbital shape and, possibly, somewhat different contributions from other configurations.

**Conclusion**

The use of femtosecond transient absorption spectroscopy with polarized light resulted in the determination of the rates of tautomerization in two tetraphenyl-substituted derivatives. In the parent symmetrical TPPo, tautomerization occurs at a similar rate to that in alkylated porphycenes that are substituted at the same positions. This result is true for both the ground- and lowest-excited singlet states. The reaction is about five-fold slower in the $S_1$ state, most-probably owing to the expansion of the inner cavity in the excited chromophore, which leads to weakening of two intramolecular hydrogen bonds.

The asymmetrical derivative 9-ATPPo exhibits a similar tautomerization rate in the $S_0$ state, where the reaction occurs in less than a picosecond, but the reaction is orders of magnitude slower in the $S_1$ state, even for a thermodynamically favored $\text{trans}-2$$\rightarrow$$\text{trans}-1$ reaction. This result is very different from the situation that is encountered in 9-AcTPPrPo, where the $\text{trans}-2$$\rightarrow$$\text{trans}-1$ conversion requires 2.6 ps. Evidently, the barrier to tautomerization is much smaller in the latter compound, even though the energy difference between the two tautomers should be larger in 9-ATPPo. The size of 9-ATPPo is rather prohibitive for excitation- and probe-pulse widths that are generated by a non-linear optical parametric amplifier (NOPA) or frequency-doubled fundamental pulses, respectively. An interference band-pass filter with a center wavelength of 650 or 700 nm was installed in the white-light-seeding beam in the NOPA to obtain a narrower excitation spectrum (FWHM ≈ 20 nm) than that generated with the standard configuration. In each case, the energy of the pump pulses at the sample position was in the range 1–2 μJ.

Absorption changes were probed with white-light supercontinuum pulses that were generated in a 3 mm-thick CaF$_2$ plate. The useful spectroscopic range of the supercontinuum extended from 360 to 780 nm. Spectra for the parallel- and perpendicular polarization of the excitation- and probe pulses were recorded and the magic-angle spectra and anisotropy data were subsequently calculated. After acquisition, all of the data were corrected for the chirp of the probe pulses. The temporal resolution of the setup was approximately 300 fs.

Stationary electronic absorption spectra were measured on a Shimadzu UV 3100 spectrophotometer. Fluorescence excitation spectra were obtained on a Fluorolog-3 (Horiba Scientific) or on an Edinburgh FS 900 CDT spectrofluorimeter. Variable-temperature emission measurements were performed on a Jasny spectrofluorimeter.[37]

The calculations of the relative ground state and the transition energies were performed by using Gaussian 09, Revision B.01 program package. Density functional theory (DFT) and time-dependent DFT (1D-DFT) calculations were performed with the B3LYP functional and the 6-31G(d,p) basis set.

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