

Multigrid interpretation of a three-level Parareal Algorithm

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1 Introduction

Parallel-in-time methods, of which parareal [13] and multigrid reduction in time (MGRIT) [3] are well-known examples, are important tools for increasing parallelism beyond traditional spatially parallel methods, see [6, 14] and references therein. As a two-level method, parareal performs the fine but expensive integration independently (and in parallel) over many short time intervals, and it uses a cheap (but coarse) integrator to correct values across time subintervals sequentially. For linear ODE systems, parareal iterates are known to be equivalent to two-level MGRIT ones for a specific choice of initial guess, restriction/prolongation operators and relaxation scheme, cf. [10, 3, 9]. One can thus analyze parareal convergence in two ways: one can make hypotheses on Lipschitz constants and truncation errors, which is typical in the ODE community, cf. [13, 1, 8], or one can use spectral information of all-at-once matrices, as is common in the multigrid community, see [3, 5, 2, 15].

When parareal and MGRIT are used with many time subintervals, the coarse correction step becomes a computational bottleneck. To overcome this, one can parallelize the coarse solution by subdividing the coarse problem and using a coarser level to ensure global communication. For MGRIT, this leads to a multilevel variant [11]; for parareal, a three-level variant has been introduced and analyzed in [12]. In this paper, we show that there is a choice of restriction/prolongation operators and relaxation schemes such that the resulting MGRIT method is equivalent to three-level parareal when applied to linear problems. The existing MGRIT literature can thus add to our understanding of three-level parareal, beyond what is shown in [12].

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2 The three-level parareal algorithm

Suppose one wishes to solve the *linear* system of ODEs $u' = \Phi u + f(t)$ with initial conditions $u(0) = u_0$ on the interval $[0, T]$. To obtain the temporal grid for both parareal and MGRIT, we subdivide the interval hierarchically as follows:¹

- The interval $[0, T]$ is subdivided into p coarsest intervals $\mathcal{I}_i = [T_{i-1}, T_i]$, $i = 1, \dots, p$, each of length $\Delta T = T/p$;
- Each coarsest interval \mathcal{I}_i is subdivided into m subintervals $\mathcal{I}_{i,j} = [t_{i,j}, t_{i,j+1}]$, $j = 0, 1, \dots, m-1$, of length $\Delta t = \Delta T/m$;
- Each $\mathcal{I}_{i,j}$ is divided into intervals $[t_{i,j,k}, t_{i,j,k+1}]$ ($0 \leq k < n$) of length $\delta t = \Delta t/n$.

We can now define the following *propagators*, which take an initial value at the beginning of \mathcal{I}_i , $\mathcal{I}_{i,j}$ or $\mathcal{I}_{i,j,k}$ and return the solution at the end of the interval:²

- F_0 is the action of the fine integrator over one fine time step δt . For a linear problem, we have $F_0 u_{i-1} = \Phi_0 u_{i-1} + f_i$.
- $F = F_0^n$ is the action of the fine integrator over one intermediate time step $\Delta t = n\delta t$. For a linear problem, we have $F u_{i-n} = \Phi_0^n u_{i-n} + \sum_{k=0}^{n-1} \Phi_0^k f_{i-k}$.
- G is the action of the intermediate integrator over one intermediate time step Δt . For a linear problem, we have $G U_{i,j-1} = \Phi_1 U_{i,j-1} + \gamma_{i,j}$.
- H is the action of the coarse integrator over one coarse time step $\Delta T = m\Delta t$. For a linear problem, we have $H Y_{i-1} = \Phi_2 Y_{i-1} + \eta_i$.

The three-level parareal algorithm, as introduced in [12], iterates on the level-1 state variables $U_{i,j}$ and level-2 state variables Y_i as follows:

1. Initialization (with iteration indices appearing as superscripts):

$$\begin{aligned} Y_0^0 &= u_0, & Y_i^0 &= H Y_{i-1}^0 \\ U_{i,0}^0 &= Y_{i-1}^0, & U_{i,j}^0 &= G U_{i,j-1}^0 \quad (1 \leq j \leq m), \end{aligned}$$

2. Iteration: for $\nu = 0, 1, 2, \dots$,

$$U_{i,0}^{\nu+1} = Y_{i-1}^{\nu}, \quad U_{i,j}^{\nu+1} = F U_{i,j-1}^{\nu} + G U_{i,j-1}^{\nu+1} - G U_{i,j-1}^{\nu} \quad (1 \leq j \leq m), \quad (1)$$

$$Y_i^{\nu+1} = u_0, \quad Y_i^{\nu+1} = U_{i,m}^{\nu+1} + H Y_{i-1}^{\nu+1} - H Y_{i-1}^{\nu}. \quad (2)$$

This method is shown in [12] to converge to the fine solution *in finitely many steps*, i.e., $U_{i,j}^{\nu} = F^{(i-1)m+j} u_0$ for $\nu \geq i(m+1)$, for any choice of G and H . Note that this is not a nested iteration, where one needs to iterate U or Y to sufficient accuracy before switching levels; instead, only one parareal step on $U_{i,j}$ is performed before it is used in (2), and one coarse parareal step (2) is performed before the Y_i are used as new initial values in (1).

¹ For ease of explanation, we assume that all subdivisions have equal length, although it is easy to see that similar results hold for non-uniform subdivisions.

² To lighten the notation, the time index is only indicated in the variable on which the propagators are applied, and not in the propagators themselves.

Algorithm 1 MGRIT($\ell, \tilde{\mathbf{g}}^{(\ell)}$) (in correction form, as defined in [3])

if ℓ is the coarsest level L **then**
 Solve coarse grid system $A_L \mathbf{u}^{(L)} = \tilde{\mathbf{g}}^{(L)}$
else
 Relax on $A_\ell \mathbf{u}^{(\ell)} = \tilde{\mathbf{g}}^{(\ell)}$ using F -relaxation
 Compute and restrict residual using injection: $\tilde{\mathbf{g}}^{(\ell+1)} = R_\ell^{\ell+1}(\tilde{\mathbf{g}}^{(\ell)} - A_\ell \mathbf{u}^{(\ell)})$
 Solve on the next level : MGRIT($\ell + 1, \tilde{\mathbf{g}}^{(\ell+1)}$)
 Correct: $\mathbf{u}^{(\ell)} \leftarrow \mathbf{u}^{(\ell)} + P_{\ell+1}^\ell \mathbf{u}^{(\ell+1)}$
end if

Algorithm 2 MGRIT-FAS($\ell, \mathbf{u}^{(\ell)}, \mathbf{g}^{(\ell)}$) (as defined in [4])

if ℓ is the coarsest level L **then**
 Solve coarse grid system $A_L(\mathbf{u}^{(L)}) = \mathbf{g}^{(L)}$
else
 Relax on $A_\ell(\mathbf{u}^{(\ell)}) = \mathbf{g}^{(\ell)}$ using F -relaxation to obtain $\mathbf{v}^{(\ell)}$
 Compute FAS right hand side: $\mathbf{g}^{(\ell+1)} = R_\ell^{\ell+1}(\mathbf{g}^{(\ell)} - A_\ell(\mathbf{v}^{(\ell)})) + A_{\ell+1}(R_\ell^{\ell+1} \mathbf{v}^{(\ell)})$
 Solve on the next level : MGRIT-FAS($\ell + 1, \mathbf{u}^{(\ell+1)}, \mathbf{g}^{(\ell+1)}$)
 Correct: $\mathbf{u}^{(\ell)} \leftarrow \mathbf{u}^{(\ell)} + P_{\ell+1}^\ell(\mathbf{u}^{(\ell+1)} - R_\ell^{\ell+1} \mathbf{v}^{(\ell)})$
end if

3 Equivalence with the MGRIT V-cycle

The initial value problems that are solved by the propagators can also be written as linear systems of the type $A_\ell \mathbf{u}^{(\ell)} = \mathbf{g}^{(\ell)}$, where

$$A_\ell = \begin{bmatrix} I & & & & \\ -\Phi_\ell & I & & & \\ & & \ddots & \ddots & \\ & & & & -\Phi_\ell & I \end{bmatrix}.$$

The index ℓ here indicates the level of coarseness of the temporal grid, with $\ell = 0$ being the finest grid, and $\ell = 2$ being the coarsest for a three-level method. Such systems can be solved using the MGRIT V-cycle with F-relaxation algorithm, which can be written in correction form [3] or as a full approximation scheme (FAS) [4], see Algorithms 1 and 2. Here, we consider the special case of $L = 2$, i.e., the three-level algorithm. For the purpose of writing the recurrence, we will index the fine grid (level-0) solution as $u_{i,j,k} \approx u(t_{i,j,k})$. The level-1 vectors will be double indexed as $u_{i,j} \approx u(t_{i,j})$, and level-2 vectors are singly indexed as $u_i \approx u(T_{i-1})$. If injection is used for $P_{\ell+1}^\ell$ and $R_\ell^{\ell+1} = (P_{\ell+1}^\ell)^T$ in Algorithm 2, then one V-cycle of MGRIT-FAS with F-relaxation for solving $A_0 \mathbf{u} = f$ updates the iterate $u_{i,j,k}$ as follows:

1. Relax on level 0:

$$v_{i,j,k} = \begin{cases} f_{i,j,k} + \Phi_0 v_{i,j,k-1}, & 1 \leq k \leq n-1, \forall i, j, \\ u_{i,j,0}, & k = 0. \end{cases}$$

2. Compute FAS right-hand side for level 1:

$$g_{i,j} = \begin{cases} f_{i,j,0} + \Phi_0 v_{i,j-1,n-1} - \Phi_1 u_{i,j-1,0}, & 1 \leq j \leq m-1, \\ f_{i,0,0} + \Phi_0 v_{i-1,m-1,n-1} - \Phi_1 u_{i-1,m-1,0}, & j = 0. \end{cases}$$

3. Relax on level 1 using initial guess $(\mathbf{u}^{(1)})_{i,j} = v_{i,j,0} = u_{i,j,0}$:

$$v_{i,j} = \begin{cases} f_{i,j,0} + \Phi_0 v_{i,j-1,n-1} + \Phi_1 (v_{i,j-1} - u_{i,j-1,0}), & 1 \leq j \leq m-1, \\ u_{i,0,0}, & j = 0. \end{cases}$$

4. Compute FAS right-hand side for level 2:

$$g_i = f_{i,0,0} + \Phi_0 v_{i-1,m-1,n-1} + \Phi_1 (v_{i-1,m-1} - u_{i-1,m-1,0}) - \Phi_2 u_{i-1,0,0}.$$

5. Solve the level-2 system:

$$u_i^{\text{new}} = f_{i,0,0} + \Phi_0 v_{i-1,m-1,n-1} + \Phi_1 (v_{i-1,m-1} - u_{i-1,m-1,0}) + \Phi_2 (u_{i-1}^{\text{new}} - u_{i-1,0,0}).$$

6. Correct on level 1 and then on level 0, using injection for both levels:

$$u_{i,j,k}^{\text{new}} = \begin{cases} f_{i,j,k} + \Phi_0 v_{i,j,k-1}, & 1 \leq k \leq n-1, \forall i, j, \\ f_{i,j,0} + \Phi_0 v_{i,j-1,n-1} + \Phi_1 (v_{i,j-1} - u_{i,j-1,0}), & k = 0, 1 \leq j \leq m-1, \forall i, \\ f_{i,0,0} + \Phi_0 v_{i-1,m-1,n-1} + \Phi_1 (v_{i-1,m-1} - u_{i-1,m-1,0}) + \Phi_2 (u_{i-1}^{\text{new}} - u_{i-1,0,0}), & j = k = 0, \forall i. \end{cases}$$

We can now prove the following equivalence theorem.

Theorem 1 For the linear problem $u' = \Phi u + f(t)$, assume that $u_{i,j,k}^0$ satisfies

$$u_{1,0,0}^0 = u_0, \quad u_{i,0,0}^0 = H u_{i-1,0,0}^0 \quad \forall i \geq 1, \quad u_{i,j,0}^0 = G u_{i,j-1,0}^0 \quad \forall j = 1, \dots, m-1.$$

Then for all $\nu \geq 0$, the three-level MGRIT-FAS V-cycle with F-relaxation and with injection as the prolongation operator is equivalent to three-level parareal via

$$u_{i,j,k}^{\nu+1} = \begin{cases} F_0^k U_{i,j}^\nu, & 1 \leq k \leq n-1, \forall i, j, \\ U_{i,j}^{\nu+1}, & k = 0, 1 \leq j \leq m-1, \forall i \geq 1, \\ Y_{i-1}^{\nu+1}, & j = k = 0, \forall i \geq 1. \end{cases}$$

Proof From the initialization conditions, we have for $\nu = 0$ that $u_{i,j,0}^\nu = U_{i,j}^\nu$ for $1 \leq j \leq m-1$, and $u_{i,0,0}^\nu = Y_{i-1}^\nu$ for all i . We will prove by induction that these two equalities also hold for $\nu \geq 1$. To do so, we rewrite $u_{i,j,k}^{\text{new}}$ in terms of the propagators F_0, F, G and H . The update formula at step 6 leads us to consider three cases:

Case 1 ($k \neq 0$). Step 1 at iteration ν reads

$$u_{i,j,k}^{\text{new}} = v_{i,j,k} = F_0 v_{i,j,k-1} = \dots = F_0^k v_{i,j,0} = F_0^k U_{i,j}^\nu.$$

Case 2 ($k = 0, j \neq 0$). This case is given by step 3, where

$$u_{i,j,0}^{\text{new}} = v_{i,j} = F_0 v_{i,j-1,n-1} + \Phi_1(v_{i,j-1} - U_{i,j-1}^Y) = F_0^n U_{i,j-1}^Y + G v_{i,j-1} - G U_{i,j-1}^Y.$$

Here, we have replaced the difference of Φ_1 by a difference of G , because G is affine. Thus, we have $v_{i,j} = U_{i,j}^{Y^{v+1}}$ for $1 \leq j \leq m-1$, since both quantities are initialized the same way (we have $v_{i,0} = u_{i,0,0} = Y_{i-1}^Y = U_{i,0}^{Y^{v+1}}$) and satisfy the same recurrence.

Case 3 ($j = k = 0$). Here we have $u_{i,0,0}^{\text{new}} = u_i^{\text{new}}$, so step 5 gives, for $i \geq 2$,

$$\begin{aligned} u_i^{\text{new}} &= f_{i,0,0} + \Phi_0 v_{i-1,m-1,n-1} + \Phi_1(v_{i-1,m-1} - u_{i-1,m-1,0}) + \Phi_2(u_{i-1}^{\text{new}} - u_{i-1,0,0}) \\ &= F_0 v_{i-1,m-1,n-1} + \Phi_1(v_{i-1,m-1} - U_{i-1,m-1}^Y) + \Phi_2(u_{i-1}^{\text{new}} - u_{i-1,0,0}) \\ &= F_0^n U_{i-1,m-1}^Y + G v_{i-1,m-1} - G U_{i-1,m-1}^Y + H u_{i-1}^{\text{new}} - H u_{i-1,0,0} \\ &= U_{i-1,m}^{Y^{v+1}} + H u_{i-1}^{\text{new}} - H Y_{i-2}^Y. \end{aligned}$$

For $i = 1$, we have $u_1^{\text{new}} = u_0 = Y_0^{Y^{v+1}}$; thus, u_i^{new} and $Y_{i-1}^{Y^{v+1}}$ satisfy the same recurrence with the same initial condition. This leads to $u_{i,0,0}^{\text{new}} = Y_{i-1}^{Y^{v+1}}$ for all i , as claimed. \square

We can now use the FAS formulation to deduce the equivalence in classical (correction) form. We define the following operators:

$$E_\ell = I - P_{\ell+1}^\ell R_\ell^{\ell+1}, \quad M_\ell = \text{diag}((A_\ell)_{11}, (A_\ell)_{22}, \dots),$$

where $(A_\ell)_{ii}$ are diagonal blocks of A_ℓ corresponding to the i th subinterval, *starting* with the coarse point and including all the fine points until (but excluding) the next coarse point. In other words, M_ℓ is the block Jacobi smoother for level ℓ , and E_ℓ blanks out the coarse points and retains the fine points when applied to a vector of values at level ℓ . Similar operators were defined in [10], where the authors proved the equivalence between two-level parareal and a geometric multigrid method with block Jacobi smoothing and aggressive coarsening in the FAS setting; however, the blocks in [10] are defined differently, with the coarse points appearing at the end of the block rather than the beginning. We write the change in the solution at step 6 as

$$u_{i,j,k}^{\text{new}} - u_{i,j,k} = \begin{cases} v_{i,j,k} - u_{i,j,k} =: (\Delta \mathbf{u}^{(0)})_{i,j,k}, & k \neq 0, \\ v_{i,j} - u_{i,j,0} =: (\Delta \mathbf{u}^{(1)})_{i,j}, & k = 0, j \neq 0, \\ u_i^{\text{new}} - u_{i,0,0} =: (\Delta \mathbf{u}^{(2)})_i, & j = k = 0. \end{cases}$$

To compute $\Delta \mathbf{u}^{(0)}$, note that $v_{i,j,k} - u_{i,j,k} = 0$ when $k = 0$; for $k \neq 0$, we have

$$\begin{aligned} (\Delta \mathbf{u}^{(0)})_{i,j,k} &= v_{i,j,k} - u_{i,j,k} = f_{i,j,k} + \Phi_0(v_{i,j,k-1} - u_{i,j,k-1}) + \Phi_0 u_{i,j,k-1} - u_{i,j,k} \\ &= (\mathbf{f} - A_0 \mathbf{u})_{i,j,k} + \Phi_0 (\Delta \mathbf{u}^{(0)})_{i,j,k-1}. \end{aligned}$$

If we move $\Phi_0 (\Delta \mathbf{u}^{(0)})_{i,j,k-1}$ to the left and recall the definition of M_0 , we get

$$M_0 \Delta \mathbf{u}^{(0)} = E_0 (\mathbf{f} - A_0 \mathbf{u}) \implies \Delta \mathbf{u}^{(0)} = M_0^{-1} E_0 \tilde{\mathbf{g}}^{(0)},$$

where $\tilde{\mathbf{g}}^{(0)} = \mathbf{f} - A_0\mathbf{u}$ is the initial residual. This is almost the same as in [10], except the residual is blanked before the smoothing, instead of after. Next, we calculate

$$(\Delta\mathbf{u}^{(1)})_{i,j} = v_{i,j} - u_{i,j,0} = \begin{cases} 0, & j = 0, \\ g_{i,j} + \Phi_1(v_{i,j-1} - u_{i,j-1,0}) + \Phi_1 u_{i,j-1,0} - u_{i,j,0}, & j \neq 0, \end{cases}$$

which implies

$$M_1\Delta\mathbf{u}^{(1)} = E_1(\mathbf{g}^{(1)} - A_1R_0^1\mathbf{u}) = E_1R_0^1(\mathbf{f}^{(0)} - A_0(\mathbf{u} + \Delta\mathbf{u}^{(0)})).$$

Thus, $\Delta\mathbf{u}^{(1)} = M_1^{-1}E_1\tilde{\mathbf{g}}^{(1)}$, where $\tilde{\mathbf{g}}^{(1)} = R_0^1(\tilde{\mathbf{g}}^{(0)} - A_0\Delta\mathbf{u}^{(0)})$. Finally, we have

$$(\Delta\mathbf{u}^{(2)})_i = u_i^{\text{new}} - u_{i,0,0} = g_i + \Phi_2(u_{i-1}^{\text{new}} - u_{i-1,0,0}) + \Phi_2 u_{i-1,0,0} - u_{i,0,0},$$

which leads to

$$A_2\Delta\mathbf{u}^{(2)} = \mathbf{g}^{(2)} - A_2R_0^2\mathbf{u} = R_1^2(\mathbf{g}^{(1)} - A_1(R_0^1\mathbf{u} + \Delta\mathbf{u}^{(1)})) = R_1^2(\tilde{\mathbf{g}}^{(1)} - A_1\Delta\mathbf{u}^{(1)}).$$

We conclude, by replacing $\Delta\mathbf{u}^{(1)}$ with $M_1^{-1}E_1\tilde{\mathbf{g}}^{(1)}$ in the last step, that

$$\begin{aligned} \mathbf{u}^{\text{new}} - \mathbf{u} &= \Delta\mathbf{u}^{(0)} + P_1^0\Delta\mathbf{u}^{(1)} + P_2^0\Delta\mathbf{u}^{(2)} \\ &= \Delta\mathbf{u}^{(0)} + P_1^0(\Delta\mathbf{u}^{(1)} + P_2^1A_2^{-1}R_1^2(\tilde{\mathbf{g}}^{(1)} - A_1\Delta\mathbf{u}^{(1)})) \\ &= \Delta\mathbf{u}^{(0)} + P_1^0((I - P_2^1A_2^{-1}R_1^2A_1)M_1^{-1}E_1 + P_2^1A_2^{-1}R_1^2)\tilde{\mathbf{g}}^{(1)} \end{aligned}$$

Defining $T = (I - P_2^1A_2^{-1}R_1^2A_1)M_1^{-1}E_1 + P_2^1A_2^{-1}R_1^2$, we continue to calculate

$$\begin{aligned} \mathbf{u}^{\text{new}} - \mathbf{u} &= \Delta\mathbf{u}^{(0)} + P_1^0TR_0^1(\tilde{\mathbf{g}}^{(0)} - A_0\Delta\mathbf{u}^{(0)}) \\ &= (P_1^0TR_0^1 + (I - P_1^0TR_0^1A_0)M_0^{-1}E_0)(\mathbf{f} - A\mathbf{u}) =: \mathcal{P}(\mathbf{f} - A\mathbf{u}). \end{aligned}$$

We conclude that the error propagator reads

$$S = I - \mathcal{P}A_0 = (I - P_1^0TR_0^1A_0)(I - M_0^{-1}E_0A_0),$$

where the operator T satisfies $I - TA_1 = (I - P_2^1A_2^{-1}R_1^2A_1)(I - M_1^{-1}E_1A_1)$. Note that the preconditioners \mathcal{P} and T can also be written as

$$\mathcal{P} = M_0^{-1}E_0 + P_1^0TR_0^1(I - A_0M_0^{-1}E_0), \quad T = M_1^{-1}E_1 + P_2^1A_2^{-1}R_1^2(I - A_1M_1^{-1}E_1).$$

We can hence interpret the action of the preconditioner \mathcal{P} as follows:

1. $M_0^{-1}E_0$: Take the fine residual, blank out the coarse points and apply block Jacobi.
2. $I - A_0M_0^{-1}E_0$: Update the residual after relaxation.
3. $P_1^0TR_0^1$: Restrict the new residual, recursively solve the coarse problem, then update the coarse points by injection.

Since T acts the same way but at a coarser level, the action of \mathcal{P} corresponds to exactly one MGRIT V-cycle with F -relaxation, written in correction form.

Remark If one replaces injection with injection plus F-relaxation (like in standard MGRIT), then the equivalent parareal formulation at the ν th iteration would be

$$\begin{aligned} U_{i,0}^{\nu+1/2} &= Y_{i-1}^\nu, & U_{i,j}^{\nu+1/2} &= GU_{i,j-1}^{\nu+1/2} + FU_{i,j-1}^\nu - GU_{i,j-1}^\nu & (1 \leq j \leq m), \\ Y_0^{\nu+1} &= u_0, & Y_i^{\nu+1} &= U_{im}^{\nu+1/2} + HY_{i-1}^{\nu+1} - HY_{i-1}^\nu, \\ U_{i,0}^{\nu+1} &= Y_{i-1}^{\nu+1}, & U_{i,j}^{\nu+1} &= GU_{i,j-1}^{\nu+1} + FU_{i,j-1}^\nu - GU_{i,j-1}^\nu & (1 \leq j \leq m). \end{aligned}$$

Note that the term $FU_{i,j-1}^\nu - GU_{i,j-1}^\nu$ is used twice, but it only needs to be computed once using a fine propagation. The intermediate propagation G , however, needs to be computed twice, since it is applied once to $U_{i,j-1}^{\nu+1/2}$, and another time to $U_{i,j-1}^{\nu+1}$.

4 Numerical example

We present the numerical example in [7], where the advection-diffusion equation $u_t = u_x + \kappa u_{xx}$ with periodic boundary conditions $u(0, t) = u(2, t)$, $u_x(0, t) = u_x(2, t)$ is solved on $t \in (0, 4)$, with $\kappa = 1/1024$ (advection-dominated case) and $u(x, 0) = e^{-20(x-1)^2}$. We discretize the problem using second order finite difference in space and backward Euler in time, with $\Delta x = 1/20$ and $\delta t = 1/1280$. For two-level parareal, the coarse propagator is backward Euler with $\Delta T = 1/2$ (8 coarse steps with 640 fine steps per coarse step). For three-level parareal, we use an intermediate level with $\Delta t = 1/128$ (10 fine steps per intermediate step), while keeping $\Delta T = 1/2$ for the coarsest level (i.e., 64 intermediate steps per coarse step). In Figure 1, we compare two-level and three-level parareal, both with and without post-smoothing. We compare both the iteration count and the *idealized* running time, as measured by the number of *non-concurrent* backward Euler steps taken at all levels; this cost is normalized by that of sequential time-stepping, so that a cost of 1 means the same cost as sequential time-stepping without parallelization. We see that two-level parareal converges to the exact solution in 8 iterations, whereas the three-level variants take many more iterations. However, the three-level iterations are much more parallel and take less time to run than a two-level iteration. In particular, both three-level versions converge with cost much lower than 1; such speedup is not possible for two-level parareal. Finally, although post-smoothing reduces the number of three-level iterations, the higher cost per iteration (two intermediate propagations rather than one) makes it slower than no post-smoothing once the normalized cost is considered.

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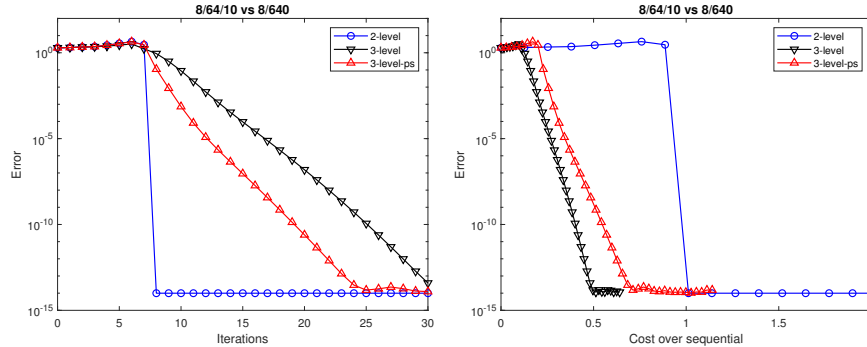


Fig. 1 Left: Iteration count for two-level parareal, and three-level parareal, with and without post-smoothing. Right: Computational cost of the three methods, as measured by the number of backward Euler steps taken, normalized by the cost of sequential time-stepping.

References

1. G. Bal. On the convergence and the stability of the parareal algorithm to solve partial differential equations. In *Domain Decomposition Methods in Science and Engineering XV*, pages 425–432. Springer, 2005.
2. V. Dobrev, T. Kolev, N. A. Petersson, and J. B. Schroder. Two-level convergence theory for multigrid reduction in time (MGRIT). *SIAM J. Sci. Comput.*, 39(5):S501–S527, 2017.
3. R. D. Falgout, S. Friedhoff, T. V. Kolev, S. P. MacLachlan, and J. B. Schroder. Parallel time integration with multigrid. *SIAM J. Sci. Comput.*, 36(6):C635–C661, 2014.
4. R. D. Falgout, A. Katz, T. V. Kolev, J. B. Schroder, A. Wissink, and U. M. Yang. Parallel time integration with multigrid reduction for a compressible fluid dynamics application. Technical Report LLNL-JRNL-663416, Lawrence Livermore National Laboratory, 2015.
5. S. Friedhoff and S. MacLachlan. A generalized predictive analysis tool for multigrid methods. *Numer. Linear Algebra Appl.*, 22(4):618–647, 2015.
6. M. J. Gander. 50 years of time parallel time integration. In *Multiple Shooting and Time Domain Decomposition Methods*, pages 69–113. Springer, 2015.
7. M. J. Gander. Five decades of time parallel time integration, and a note on the degradation of the performance of the Parareal algorithm as a function of the Reynolds number. *Oberwolfach Report*, 2017.
8. M. J. Gander and E. Hairer. Nonlinear convergence analysis for the parareal algorithm. In *Domain Decomposition Methods in Science and Engineering XVII*, pages 45–56. 2008.
9. M. J. Gander, F. Kwok, and H. Zhang. Multigrid interpretations of the parareal algorithm leading to an overlapping variant and MGRIT. *Comput. Visualization Sci.*, 19(3):59–74, 2018.
10. M. J. Gander and S. Vandewalle. Analysis of the parareal time-parallel time-integration method. *SIAM J. Sci. Comput.*, 29(2):556–578, 2007.
11. A. Henthaler, B. S. Southworth, D. Nordsletten, O. Röhrle, R. D. Falgout, and J. B. Schroder. Multilevel convergence analysis of multigrid-reduction-in-time. *SIAM J. Sci. Comput.*, 42(2):A771–A796, 2020.
12. F. Kwok. Analysis of a three-level variant of parareal. In preparation, 2022.
13. J.-L. Lions, Y. Maday, and G. Turinici. Résolution d’EDP par un schéma en temps «pararéel». *C. R. Acad. Sci., Ser. I: Math.*, 332(7):661–668, 2001.
14. B. W. Ong and J. B. Schroder. Applications of time parallelization. *Comput. Visualization Sci.*, 23(1):1–15, 2020.
15. B. S. Southworth. Necessary conditions and tight two-level convergence bounds for parareal and multigrid reduction in time. *SIAM J. Matrix Anal. Appl.*, 40(2):564–608, 2019.