

A CONTINUOUS ANALYSIS OF NEUMANN-NEUMANN METHODS: SCALABILITY AND NEW COARSE SPACES*

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Abstract. We present a new coarse space correction for the iterative Neumann-Neumann method. We describe the method for general elliptic partial differential equations, and perform the analysis for the case of the Poisson and screened Poisson equation (sometimes also called positive definite Helmholtz equation, or Helmholtz equation with the good sign). We prove that the new two-level Neumann-Neumann method converges after one iteration, both at the continuous and discrete level, which means the new coarse space is optimal in the sense of best possible, and it makes the two-level method a direct solver. In two and three space dimensions, the new coarse space is too high dimensional in practice, and we introduce a spectral approximation, which transforms a divergent iterative Neumann-Neumann method into a convergent one. We also identify what the optimized choice of coarse space functions is in the approximation. Our new coarse space thus also addresses convergence or robustness problems of the underlying domain decomposition iteration, similarly to the new coarse spaces GenEO, SHEM, and ACMS, which were designed to treat different convergence difficulties of the underlying domain decomposition method, namely the presence of high contrast media. Several numerical experiments are carried out to demonstrate the performance of this new coarse space correction, also including decompositions with cross points.

Key words. domain decomposition methods, elliptic problems, Neumann-Neumann methods

AMS subject classifications. 65N55, 65F08, 65F10, 65F50

1. Introduction. We design and analyze a new coarse space correction for domain decomposition solvers of algebraic equations arising from the discretization of second-order elliptic PDEs. Domain decomposition methods are of great interest because of their natural parallelism permitting the use of parallel architectures in order to approximate the solution of partial differential equations. Another important benefit of these methods is that they allow for a better treatment of complex geometries and they are well suited for heterogeneous problems, i.e., problems that have different physics in different parts of the domain, see e.g. [30, 33, 22]. We focus here on a class of non-overlapping domain decomposition methods known in the literature as the Neumann-Neumann methods (NNMs). As most domain decomposition methods, the classical one-level NNMs lack global communication, since the only mechanism for sharing data is through interfaces. A remedy for this is to introduce an additional coarse space correction with a relatively small cost compared to the size of the original problem, see the seminal early contributions [29, 9]. The design of efficient coarse spaces is at the heart of domain decomposition theory and it is in general not straightforward. For the case of algebraic NNM preconditioners, seminal contributions are e.g. [25, 26, 27, 10, 8]. However, the NNM was originally described in [11] as an iteration at the continuous level like the classical Schwarz method, but only for two subdomains. This is because unlike the Schwarz method [24], the convergence of the iterative NNM is not guaranteed for the case of many subdomains, even in the case of a one way decomposition without cross points, when the subdomain aspect

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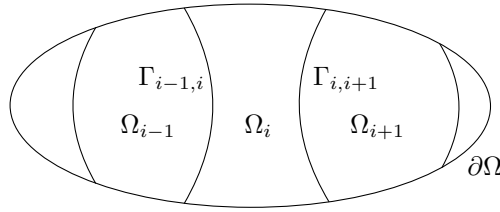


Fig. 1: One way decomposition of the domain Ω .

ratio is unfavorable, see [2, 1]. We are interested here in developing a coarse space correction that leads to a convergent iterative NNM.

An interesting discovery was made in [15], where the authors described a coarse space for the parallel Schwarz method that leads to convergence after one coarse correction step. This became known in the literature as an optimal coarse space; i.e. better convergence cannot be achieved (note optimal here is not to be understood in the sense of scalable). This approach allowed the authors in [14] to describe a new coarse space for optimized Schwarz methods, and led to the new Spectral Harmonically Enriched Multiscale (SHEM) coarse space [19, 18], which was compared to the new ACMS coarse space in [23]. Like GenEO [31], SHEM, and ACMS were developed to handle convergence problems of the underlying domain decomposition method in the presence of high contrast. Our new coarse space here is designed to handle convergence problems for subdomain aspect ratios which create convergence problems of the NNM. It is based on ideas in the short conference paper [3], and we present a complete formulation and detailed analysis here, including the cross point case. We also present coarse space corrections for problems for which the NNM cannot directly be applied due to the problem of floating subdomains.

The paper is organized as follows: in Section 2, we recall the continuous iterative NNM and give a convergence analysis for one-way decompositions in one, two, and three spatial dimensions. We explain in Section 3 how to construct the coarse space correction in one, two, and three spatial dimensions, and also for problems for which the NNM is not well defined due to floating subdomains, and we give corresponding convergence estimates. We show numerical results illustrating the performance of the new method in Section 4, and present the optimal coarse space in the cross point case at the discrete level in Section 5.

2. Analysis of the one-level NNM. Let Ω be a bounded domain in \mathbb{R}^d , $d = 1, 2$. We consider the partial differential equation

$$(2.1) \quad \begin{aligned} \mathcal{L}u &= f, \text{ in } \Omega, \\ u &= 0, \text{ on } \partial\Omega, \end{aligned}$$

where \mathcal{L} is a second order elliptic operator, and $f \in L^2(\Omega)$. We want to solve the problem defined in (2.1) using the NNM. Let $\Omega_1, \dots, \Omega_N$ be a one way non-overlapping decomposition of Ω as shown in Figure 1 (for decompositions with cross points see Section 5). Let n_i denote the unit outward normal on $\partial\Omega_i$. The iterative one-level NNM is then given by Algorithm 2.1.

2.1. One-dimensional analysis. We consider a one dimensional decomposition of Ω into N equally-sized subdomains Ω_i as shown in Figure 2, and study the

Algorithm 2.1 One-level NNM

1. Set $g_{i,j}^0$ to zero or any inexpensive initial guess on the interfaces $\partial\Omega_i \cap \partial\Omega_j$
2. Repeat until convergence
 - (a) Solve the Dirichlet followed by the Neumann problems

$$\begin{aligned} \mathcal{L}u_i^n &= f \text{ in } \Omega_i, & \mathcal{L}\psi_i^n &= 0 \text{ in } \Omega_i, \\ u_i^n &= g_{i,j}^n \text{ on } \partial\Omega_i \cap \partial\Omega_j, & \partial_{n_i}\psi_i^n &= (\partial_{n_i}u_i^n + \partial_{n_j}u_j^n)/2 \text{ on } \partial\Omega_i \cap \partial\Omega_j, \\ u_i^n &= 0 \text{ on } \partial\Omega_i \cap \partial\Omega. & \psi_i^n &= 0 \text{ on } \partial\Omega_i \cap \partial\Omega. \end{aligned}$$

- (b) Update the traces

$$g_{i,j}^{n+1} := g_{i,j}^{n+1} - \frac{1}{2} (\psi_i^n + \psi_j^n) \text{ on } \partial\Omega_i \cap \partial\Omega_j.$$

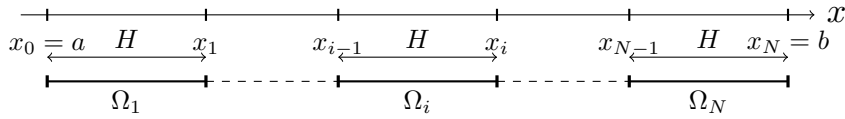


Fig. 2: One-dimensional geometry.

77 convergence of Algorithm 2.1 for the screened Laplace operator $\mathcal{L} = \eta^2 - \Delta$. By lin-
 78 earity of (2.1), it suffices to set $f = 0$ and study the convergence of $(g_{i,i+1}^n)_{1 \leq i \leq N-1}$ to
 79 zero. To simplify the notation, we set here $g_i := g_{i,i+1}$ on the interface $\Gamma_i := \Gamma_{i,i+1} =$
 80 $\partial\Omega_i \cap \partial\Omega_{i+1}$. We first prove a lemma that gives a recurrence relation between the
 81 iterates:

82 LEMMA 2.1. Let $\mathbf{g}^n = [g_1^n, g_2^n, \dots, g_{N-1}^n]^T \in \mathbb{R}^{N-1}$, then for $N \geq 3$, we have
 83 $\mathbf{g}^n = T\mathbf{g}^{n-1}$, where $T \in \mathbb{R}^{(N-1) \times (N-1)}$ is given by

$$84 \quad T := -\frac{1}{4 \sinh^2(\eta H)} \begin{bmatrix} 1 & \frac{1}{\cosh(\eta H)} & -1 & 0 & \cdots & \cdots & 0 \\ 0 & 2 & 0 & -1 & \ddots & & \vdots \\ -1 & 0 & 2 & 0 & -1 & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots & 0 \\ \vdots & \ddots & \ddots & \ddots & 2 & 0 & -1 \\ \vdots & & \ddots & -1 & 0 & 2 & 0 \\ 0 & \cdots & \cdots & 0 & -1 & \frac{1}{\cosh(\eta H)} & 1 \end{bmatrix}.$$

85 *Proof.* The subdomain solutions are given by

$$86 \quad u_i^n(x) = g_i^n \frac{\sinh(\eta(x - x_{i-1}))}{\sinh(\eta H)} + g_{i-1}^n \frac{\sinh(\eta(x_i - x))}{\sinh(\eta H)},$$

87 for $i = 1, \dots, N$, where we defined $g_0^n = g_N^n = 0$ for simplicity. Similarly we obtain

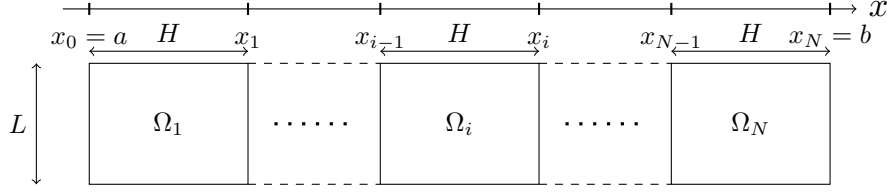


Fig. 3: Two-dimensional geometry.

88 for the local corrections

$$89 \quad \psi_i^n = \left(2g_i^n \frac{\cosh(\eta H)}{\sinh(\eta H)} - \frac{g_{i-1}^n}{\sinh(\eta H)} - \frac{g_{i+1}^n}{\sinh(\eta H)} \right) \frac{\cosh(\eta(x - x_{i-1}))}{2 \sinh(\eta H)} \\ + \left(2g_{i-1}^n \frac{\cosh(\eta H)}{\sinh(\eta H)} - \frac{g_{i-2}^n}{\sinh(\eta H)} - \frac{g_i^n}{\sinh(\eta H)} \right) \frac{\cosh(\eta(x_i - x))}{2 \sinh(\eta H)},$$

90 for $i = 2, \dots, N - 2$, and

$$91 \quad \psi_1^n = \left(2g_1^n \frac{\cosh(\eta H)}{\sinh(\eta H)} - \frac{g_2^n}{\sinh(\eta H)} \right) \frac{\sinh(\eta(x - x_0))}{2 \cosh(\eta H)}, \\ \psi_N^n = \left(2g_{N-1}^n \frac{\cosh(\eta H)}{\sinh(\eta H)} - \frac{g_{N-2}^n}{\sinh(\eta H)} \right) \frac{\sinh(\eta(x_N - x))}{2 \cosh(\eta H)}.$$

92 Hence, we get the stated relation. \square

93 **THEOREM 2.2.** *If the width of the subdomains satisfies $H > \frac{\ln(1+\sqrt{2})}{\eta}$, then the*
 94 *one-level NNM for the screened Laplace problem $(\eta^2 - \Delta)u = 0$ in 1D given by Algo-*
 95 *rithm 2.1 is convergent, and satisfies the convergence estimate*

$$96 \quad (2.2) \quad \max_{1 \leq i \leq N-1} |g_i^n| \leq \frac{1}{\sinh^{2n}(\eta H)} \max_{1 \leq i \leq N-1} |g_i^0|.$$

97 *Proof.* By Lemma 2.1, we have that

$$\|T\|_\infty = \max \left\{ \frac{1}{\sinh^2(\eta H)}, \frac{1}{2 \sinh^2(\eta H)} + \frac{1}{4 \sinh^2(\eta H) \cosh(\eta H)} \right\} \leq \frac{1}{\sinh^2(\eta H)}.$$

98 \square

99 **2.2. Two-dimensional analysis.** We suppose that the decomposition of the
 100 domain Ω is as in Figure 3 and analyze the convergence of Algorithm 2.1 again for
 101 $\mathcal{L} = \eta^2 - \Delta$. Since the subdomains are rectangular, the iterates can be expanded in
 102 a sine series, i.e.

$$103 \quad (2.3) \quad u_i^n(x, y) = \sum_{m=1}^{\infty} \hat{u}_i^n(x, m) \sin\left(\frac{m\pi}{L} y\right), \quad \psi_i^n(x, y) = \sum_{m=1}^{\infty} \hat{\psi}_i^n(x, m) \sin\left(\frac{m\pi}{L} y\right),$$

104 where \hat{u}_i^n and $\hat{\psi}_i^n$ are the Fourier coefficients of u_i^n and ψ_i^n respectively. As in the
 105 one-dimensional case, we obtain a lemma giving the recurrence relation of the Fourier
 106 coefficients:

107 LEMMA 2.3. Let $\hat{\mathbf{g}}^n(m) = [\hat{g}_1^n(m), \hat{g}_2^n(m), \dots, \hat{g}_{N-1}^n(m)]^T \in \mathbb{R}^{N-1}$, then for $N \geq$
 108 3 we have $\hat{\mathbf{g}}^n(m) = T_m \hat{\mathbf{g}}^{n-1}(m)$, where $T_m \in \mathbb{R}^{(N-1) \times (N-1)}$ is given by

$$109 \quad T_m := -\frac{1}{4 \sinh^2(k_m H)} \begin{bmatrix} 1 & \frac{1}{\cosh(k_m H)} & -1 & 0 & \cdots & \cdots & 0 \\ 0 & 2 & 0 & -1 & \ddots & & \vdots \\ -1 & 0 & 2 & 0 & -1 & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots & 0 \\ \vdots & \ddots & \ddots & \ddots & 2 & 0 & -1 \\ \vdots & & \ddots & -1 & 0 & 2 & 0 \\ 0 & \cdots & \cdots & 0 & -1 & \frac{1}{\cosh(k_m H)} & 1 \end{bmatrix}$$

110 where $k_m := \sqrt{\eta^2 + \frac{m^2 \pi^2}{L^2}}$.

111 *Proof.* For each $m \geq 1$ and $i = 2, \dots, N-1$, $u_i^n(x, m)$ and $\psi_i^n(x, m)$ satisfy

$$112 \quad \begin{aligned} k_m^2 \hat{u}_i^n - \partial_{xx} \hat{u}_i^n &= 0, & k_m^2 \hat{\psi}_i^n - \partial_{yy} \hat{\psi}_i^n &= 0, \\ \hat{u}_i^n(x_{i-1}, m) &= \hat{g}_{j-1}^n(m), & \hat{\psi}_i^n(x_{i-1}, m) &= (\partial_x \hat{u}_i^n(x_{j-1}, m) - \partial_x \hat{u}_{i-1}^n(x_{j-1}, m))/2, \\ \hat{u}_i^n(x_i, m) &= \hat{g}_i^n(m), & \hat{\psi}_i^n(x_i, m) &= (\partial_x \hat{u}_i^n(x_i, m) - \partial_x \hat{u}_{i+1}^n(x_i, m))/2. \end{aligned}$$

113 The solution of the Dirichlet problems on the interior subdomains are thus

$$114 \quad \hat{u}_i^n(x, m) = \hat{g}_i^n(m) \frac{\sinh(k_m(x - x_{i-1}))}{\sinh(k_m H)} + \hat{g}_{i-1}^n(m) \frac{\sinh(k_m(x_i - x))}{\sinh(k_m H)},$$

115 for $i = 1, \dots, N$, where we defined $\hat{g}_0^n = \hat{g}_N^n = 0$ to include the solution to boundary
 116 subdomains. Similarly for the Neumann problems on the interior subdomains, we
 117 obtain

$$118 \quad \begin{aligned} \hat{\psi}_i^n(x, m) &= \left(2 \hat{g}_i^n(m) \frac{\cosh(k_m H)}{\sinh(k_m H)} - \frac{\hat{g}_{i-1}^n(m)}{\sinh(k_m H)} - \frac{\hat{g}_{i+1}^n(m)}{\sinh(k_m H)} \right) \frac{\cosh(k_m(x - x_{j-1}))}{2 \sinh(k_m H)} \\ &+ \left(2 \hat{g}_{i-1}^n(m) \frac{\cosh(k_m H)}{\sinh(k_m H)} - \frac{\hat{g}_{i-2}^n(m)}{\sinh(k_m H)} - \frac{\hat{g}_i^n(m)}{\sinh(k_m H)} \right) \frac{\cosh(k_m(x_i - x))}{2 \sinh(k_m H)}, \end{aligned}$$

119 and for the first and last subdomains, we find

$$120 \quad \begin{aligned} \hat{\psi}_1^n(x, m) &= \left(2 \hat{g}_1^n(m) \frac{\cosh(k_m H)}{\sinh(k_m H)} - \frac{\hat{g}_2^n(m)}{\sinh(k_m H)} \right) \frac{\sinh(k_m(x - x_0))}{2 \cosh(k_m H)}, \\ \hat{\psi}_N^n(x, m) &= \left(2 \hat{g}_{N-1}^n(m) \frac{\cosh(k_m H)}{\sinh(k_m H)} - \frac{\hat{g}_{N-2}^n(m)}{\sinh(k_m H)} \right) \frac{\sinh(k_m(x_N - x))}{2 \cosh(k_m H)}, \end{aligned}$$

121 from which we obtain the stated formula. \square

122 THEOREM 2.4. If the width of the subdomains satisfies $H > \frac{\ln(1+\sqrt{2})}{k_1}$, $k_1 =$
 123 $\sqrt{\eta^2 + \frac{\pi^2}{L^2}}$ with L the height of the subdomains, then the one-level NNM for the
 124 screened Laplace problem $(\eta^2 - \Delta)u = 0$ in 2D given by Algorithm 2.1 is conver-
 125 gent and satisfies the L^2 convergence estimate

$$126 \quad (2.4) \quad \max_{1 \leq i \leq N-1} \|g_i^n\|_2 \leq \frac{1}{\sinh^{2n}(k_1 H)} \max_{1 \leq i \leq N-1} \|g_i^0\|_2.$$

127 *Proof.* We define the sequences $\Lambda_i^n = \{\hat{g}_i^n(m)\}_{m \geq 1}$. By Lemma 2.3, we have

$$128 \quad \hat{g}_i^{n+1}(m) = \frac{1}{\sinh(k_m H)^2} \left(\frac{1}{4} \hat{g}_{i-2}^n(m) - \frac{1}{2} \hat{g}_i^n(m) + \frac{1}{4} \hat{g}_{i+2}^n(m) \right)$$

129 for each $m \geq 1$. Using then Parseval's identity $\|\Lambda_i^n\|_2^2 = \frac{L}{2} \|g_i^n\|_2^2$, we have for $i =$
130 $3, \dots, N-3$

$$\begin{aligned} \|\Lambda_i^{n+1}\|_2 &\leq \frac{1}{\sinh^2(k_1 H)} \left(\frac{1}{4} \|\Lambda_{i-2}^n + \frac{1}{2} \Lambda_i^n + \frac{1}{4} \Lambda_{i+2}^n\|_2 \right), \\ 131 \quad &\leq \frac{1}{\sinh^2(k_1 H)} \left(\frac{1}{4} \|\Lambda_{i-2}^n\|_2 + \frac{1}{2} \|\Lambda_i^n\|_2 + \frac{1}{4} \|\Lambda_{i+2}^n\|_2 \right), \\ &\leq \frac{1}{\sinh^2(k_1 H)} \max_{1 \leq i \leq N-1} \|\Lambda_i^n\|_2, \end{aligned}$$

132 where we used the triangle inequality, and the monotonicity of $m \mapsto 1/\sinh^2(k_m H)$.
133 Similarly one can show that the same bound also holds for the remaining subdomains
134 $i = 1, 2, N-2, N-1$, and hence we get the stated result. \square

135 From our analysis, we see that the one-level iterative NNMs only converge provided
136 the subdomain width H is large enough (compared to the height L in 2D). If this is
137 not the case, the iterative method diverges, but could still be used as a preconditioner
138 for a Krylov method, see the numerical experiments in Section 4. Note also that a
139 large η screening parameter helps convergence.

140 **2.3. Three-dimensional analysis.** In this part, we suppose that the domain
141 $\Omega = [a, b] \times [0, L] \times [0, L']$ is decomposed into N non-overlapping subdomains $\Omega_i =$
142 $[x_{i-1}, x_i] \times [0, L] \times [0, L']$, where $i = 1, \dots, N$. Since the subdomains represent three-
143 dimensional bricks, we can expand the iterates u_i^n and ψ_i^n in a double sine series,
144 i.e.,

$$\begin{aligned} 145 \quad (2.5) \quad u_i^n(x, y, z) &= \sum_{m, m' \geq 1}^{\infty} \hat{u}_i^n(x, m, m') \sin\left(\frac{m\pi}{L} y\right) \sin\left(\frac{m'\pi}{L'} z\right), \\ \psi_i^n(x, y, z) &= \sum_{m, m' \geq 1}^{\infty} \hat{\psi}_i^n(x, m, m') \sin\left(\frac{m\pi}{L} y\right) \sin\left(\frac{m'\pi}{L'} z\right), \end{aligned}$$

146 where \hat{u}_i^n and $\hat{\psi}_i^n$ are the Fourier coefficients of u_i^n and ψ_i^n . Following the analysis in
147 the one- and two-dimensional case, we obtain:

148 **LEMMA 2.5.** *Let $\hat{\mathbf{g}}^n(m, m') = [\hat{g}_1^n(m, m'), \hat{g}_2^n(m, m'), \dots, \hat{g}_{N-1}^n(m, m')]^T \in \mathbb{R}^{N-1}$,*
149 *then for $N \geq 3$ we have $\hat{\mathbf{g}}^n(m, m') = T_{m, m'} \hat{\mathbf{g}}^{n-1}(m, m')$, where $T_{m, m'} \in \mathbb{R}^{(N-1) \times (N-1)}$*

150 is given by

$$151 \quad T_{m,m'} := -\frac{1}{4 \sinh^2(k_{m,m'}H)} \begin{bmatrix} 1 & \frac{1}{\cosh(k_{m,m'}H)} & -1 & 0 & \cdots & \cdots & 0 \\ 0 & 2 & 0 & -1 & \ddots & & \vdots \\ -1 & 0 & 2 & 0 & -1 & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots & 0 \\ \vdots & \ddots & \ddots & \ddots & 2 & 0 & -1 \\ \vdots & & & & & & \\ \vdots & & & & & & \\ 0 & \cdots & \cdots & 0 & -1 & \frac{1}{\cosh(k_{m,m'}H)} & 1 \end{bmatrix},$$

$$152 \quad \text{with } k_{m,m'} := \sqrt{\eta^2 + \frac{m^2\pi^2}{L^2} + \frac{m'^2\pi^2}{L'^2}}.$$

153 *Proof.* The proof of Lemma 2.5 is similar to the one of Lemma 2.3. In fact, it
154 suffices to observe that the Fourier coefficients $\hat{u}_i^n(x, m, m')$ and $\hat{\psi}_i^n(x, m, m')$ satisfy

$$155 \quad \begin{aligned} k_{m,m'}^2 \hat{u}_i^n - \partial_x \hat{u}_i^n &= 0, & k_{m,m'}^2 \hat{\psi}_i^n - \partial_{yy} \hat{\psi}_i^n &= 0, \\ \hat{u}_i^n(x_{i-1}, m, m') &= \hat{g}_{j-1}^n(m, m'), & \hat{\psi}_i^n(x_{i-1}, m, m') &= (\partial_x \hat{u}_i^n(x_{j-1}, m, m') - \partial_x \hat{u}_{i-1}^n(x_{j-1}, m, m'))/2, \\ \hat{u}_i^n(x_i, m, m') &= \hat{g}_i^n(m, m'), & \hat{\psi}_i^n(x_i, m, m') &= (\partial_x \hat{u}_i^n(x_i, m, m') - \partial_x \hat{u}_{i+1}^n(x_i, m, m'))/2, \end{aligned}$$

156 which are similar to the ones in Lemma 2.3, except that $k_{m,m'}$ includes the frequen-
157 cies of the Fourier expansion of both the y - and z -direction. Using then the same
158 arguments of Lemma 2.3, we obtain the desired result. \square

159 **THEOREM 2.6.** *If the width of the subdomains in 3D satisfies $H > \frac{\ln(1+\sqrt{2})}{k_{1,1}}$,*
160 $k_{1,1} = \sqrt{\eta^2 + \frac{\pi^2}{L^2} + \frac{\pi^2}{(L')^2}}$ *with L, L' the dimensions of the subdomains in y - and*
161 z -*direction, then the one-level NNM for the screened Laplace problem $(\eta^2 - \Delta)u = 0$*
162 *in 3D given by Algorithm 2.1 is convergent and satisfies the L^2 convergence estimate*

$$163 \quad (2.6) \quad \max_{1 \leq i \leq N-1} \|g_i^n\|_2 \leq \frac{1}{\sinh^{2n}(k_{1,1}H)} \max_{1 \leq i \leq N-1} \|g_i^0\|_2.$$

164 *Proof.* Similar to the proof of Theorem 2.4, it suffices to define the sequence
165 $\Lambda_i^n = \{|\hat{g}_i^n(m, m')|\}_{m, m' \geq 1}$. Then, using Lemma 2.5, we obtain

$$166 \quad \hat{g}_i^{n+1}(m, m') = \frac{1}{\sinh(k_{m,m'}H)^2} \left(\frac{1}{4} \hat{g}_{i-2}^n(m, m') - \frac{1}{2} \hat{g}_i^n(m, m') + \frac{1}{4} \hat{g}_{i+2}^n(m, m') \right)$$

167 for each $m, m' \geq 1$. Noting that the Parseval identity still holds in 3D, and is given
168 by $\|\Lambda_i^n\|_2^2 = \frac{LL'}{4} \|g_i^n\|_2^2$, we obtain by the triangle inequality that

$$\begin{aligned} \|\Lambda_i^{n+1}\|_2 &\leq \frac{1}{\sinh^2(k_{1,1}H)} \left(\frac{1}{4} \|\Lambda_{i-2}^n\|_2 + \frac{1}{2} \|\Lambda_i^n\|_2 + \frac{1}{4} \|\Lambda_{i+2}^n\|_2 \right), \\ 169 \quad &\leq \frac{1}{\sinh^2(k_{1,1}H)} \left(\frac{1}{4} \|\Lambda_{i-2}^n\|_2 + \frac{1}{2} \|\Lambda_i^n\|_2 + \frac{1}{4} \|\Lambda_{i+2}^n\|_2 \right), \\ &\leq \frac{1}{\sinh^2(k_{1,1}H)} \max_{1 \leq i \leq N-1} \|\Lambda_i^n\|_2. \end{aligned}$$

170 The same bound holds for the subdomains $i = 1, 2, N-2, N-1$. This shows that
171 Algorithm 2.1 converges in 3D if the condition $k_{1,1}H > \ln(1 + \sqrt{2})$ is satisfied. \square

172 **3. Analysis of the two-level NNM.** In this section, we explain how to add
 173 a coarse space correction to Algorithm 2.1. While in the classical approach a coarse
 174 space is added to make the method scalable, here we ask the coarse space to do more,
 175 namely make the method convergent for all subdomain width H given a subdomain
 176 height L . This is similar to the new coarse spaces developed over the last decade
 177 like GenEO, ACMS, and SHEM, which were also designed to deal with convergence
 178 or robustness problems of the underlying domain decomposition methods. The key
 179 idea for the new coarse correction comes from the observation that at convergence of
 180 the NNM, the intermediate steps $u_i^{n+\frac{1}{2}} := u_i^n - \psi_i^n$ are continuous in the Dirichlet
 181 trace, but during the iteration before convergence they are not, which means that
 182 an efficient correction should be able to reduce the jump between neighboring $u_i^{n+\frac{1}{2}}$
 183 as much as possible in order to improve convergence. A good coarse space should
 184 thus contain enough harmonic functions¹ that are discontinuous across subdomain
 185 boundaries. Thus a complete coarse space X can be defined as

$$186 \quad (3.1) \quad X := \{v \in L^2(\Omega) : v_i := v|_{\Omega_i} \in H^1(\Omega_i), \mathcal{L}v_i = 0 \text{ in } \Omega_i, v_i = 0 \text{ on } \partial\Omega_i \cap \partial\Omega\}.$$

187 Moreover, the size of the coarse space X defined in (3.1) can be reduced. In fact, since
 188 the iterates $u_i^{n+\frac{1}{2}}$ are already continuous in the normal derivative, it is sufficient that
 189 the optimal coarse space is also continuous in the normal derivative across subdomain
 190 interfaces. The optimal coarse space is thus given taking all these functions,

$$191 \quad (3.2) \quad X_d := \left\{ v \in X : v_i := v|_{\Omega_i}, \frac{\partial v_i}{\partial n_i} + \frac{\partial v_j}{\partial n_j} = 0 \text{ on } \partial\Omega_i \cap \partial\Omega_j \right\}.$$

192 Now, since the coarse correction must reduce the Dirichlet jump of the iterates $u_i^{n+\frac{1}{2}}$,
 193 we choose the correction $U^n \in X_d$ such that it satisfies

$$194 \quad (3.3) \quad U^n := \arg \min_{v \in X_d \subset X} q(u^{n+\frac{1}{2}} + v),$$

195 where \tilde{X}_d is a finite-dimensional subspace of the optimal coarse space X_d , and the
 196 quadratic functional q is defined by

$$197 \quad (3.4) \quad \begin{aligned} q : X_d &\mapsto \mathbb{R}_+ \\ u &\mapsto \sum_{\partial\Omega_i \cap \partial\Omega_j \neq \emptyset} \int_{\partial\Omega_i \cap \partial\Omega_j} |u_i - u_j|^2 \, ds. \end{aligned}$$

198 This leads at the continuous level to the new two-level NNM given in Algorithm 3.1.

199 **3.1. One-dimensional analysis.** We now analyze the convergence of the two-
 200 level NNM for the screened Laplace operator $\mathcal{L} = \eta^2 - \Delta$ and the Laplace operator
 201 $\mathcal{L} = -\Delta$ in one dimension. The Laplace operator needs a separate analysis, since the
 202 Neumann correction problems are then not well posed for interior subdomains.

203 **3.1.1. Screened Laplace operator $\mathcal{L} = \eta^2 - \Delta$.** Since the optimal coarse
 204 space X_d in 1D is finite dimensional, it is a practical choice for \tilde{X}_d and makes the
 205 Algorithm 3.1 for the problem $\mathcal{L} = \eta^2 - \Delta$ optimal in the sense that it becomes a
 206 direct solver:

¹Here harmonic means satisfying the homogeneous equation $\mathcal{L}u = 0$.

Algorithm 3.1 Two-level NNM

1. Set $g_{i,j}^0$ to zero or any inexpensive initial guess on the interfaces $\partial\Omega_i \cap \partial\Omega_j$
2. Repeat until convergence
 - (a) Solve the Dirichlet followed by the Neumann problems

$$\begin{aligned} \mathcal{L}u_i^n &= f \text{ in } \Omega_i, & \mathcal{L}\psi_i^n &= 0 \text{ in } \Omega_i, \\ u_i^n &= g_{i,j}^n \text{ on } \partial\Omega_i \cap \partial\Omega_j, & \partial_{n_i}\psi_i^n &= (\partial_{n_i}u_i^n + \partial_{n_j}u_j^n)/2 \text{ on } \partial\Omega_i \cap \partial\Omega_j, \\ u_i^n &= 0 \text{ on } \partial\Omega_i \cap \partial\Omega. & \psi_i^n &= 0 \text{ on } \partial\Omega_i \cap \partial\Omega. \end{aligned}$$

- (b) Set

$$u_i^{n+\frac{1}{2}} := u_i^n - \psi_i^n \text{ in } \Omega_i.$$

- (c) Set

$$\tilde{u}^{n+1} := u^{n+\frac{1}{2}} + U^n$$

where U^n is as in (3.3).

- (d) Update the traces

$$g_{ij}^{n+1} := \frac{1}{2} (\tilde{u}_i^{n+1} + \tilde{u}_j^{n+1}) \text{ on } \partial\Omega_i \cap \partial\Omega_j.$$

207 **THEOREM 3.1.** *The optimal coarse space X_d is finite dimensional in 1D, and with*
 208 $\tilde{X}_d := X_d$, *Algorithm 3.1 in 1D for $\mathcal{L} = \eta^2 - \Delta$ converges after one iteration.*

209 *Proof.* Since the Dirichlet traces at the interfaces are just numbers, the optimal
 210 coarse space X_d is finite dimensional, and thus $\tilde{X}_d := X_d$ is a practical choice. Let
 211 u_i^0 be the solutions after solving with the initial guess. The function defined by
 212 $u_i^{\frac{1}{2}} := u_i^0 - \psi_i^0$ is in X_d . Now in order to compute the correction U^0 in (3.3), we proceed
 213 as follows: since the vector space X_d is of finite dimension in 1D, ($\dim(X_d) = N - 1$),
 214 it admits a finite basis $\phi_1, \phi_2, \dots, \phi_{N-1}$. We construct the coarse basis by choosing
 215 ϕ_i on each interface $x = x_i$ such that

$$216 \quad \phi_i := \begin{cases} -\frac{\cosh(\eta(x-x_{i-1}))}{\sinh(\eta H)} & \text{if } x \in (x_{i-1}, x_i), \\ \frac{\cosh(\eta(x_{i+1}-x))}{\sinh(\eta H)} & \text{if } x \in (x_i, x_{i+1}), \end{cases}$$

217 for $i = 2, \dots, N - 2$, and

$$218 \quad \phi_1 := \begin{cases} -\frac{\sinh(\eta(x-x_0))}{\cosh(\eta H)} & \text{if } x \in (x_0, x_1), \\ \frac{\cosh(\eta(x_2-x))}{\sinh(\eta H)} & \text{if } x \in (x_1, x_2), \end{cases}$$

219

$$220 \quad \phi_{N-1} := \begin{cases} \frac{\cosh(\eta(x-x_{N-2}))}{\sinh(\eta H)} & \text{if } x \in (x_{N-2}, x_{N-1}), \\ -\frac{\sinh(\eta(x_N-x))}{\cosh(\eta H)} & \text{if } x \in (x_{N-1}, x_N). \end{cases}$$

221 Since $U^0 \in X_d$, we have that $U^0 = \sum_{i=1}^{N-1} \alpha_i \phi_i$ for some coefficients α_i . Define $u_i^+ :=$

222 $u_{i+1}^{\frac{1}{2}}(x_i)$ and $u_i^- := u_i^{\frac{1}{2}}(x_i)$. The minimization problem in (3.3) reduces to finding

- 247 LEMMA 3.2 (Properties of the matrix L).
 248 1. L is surjective i.e. $(\text{rank}(L) = N - 2)$.
 249 2. $\ker(L) = \mathbb{R} [1, 1, \dots, 1]^T$.
 250 3. Its pseudoinverse $L^\dagger \in \mathbb{R}^{N-1 \times N-2}$ is given by

$$251 \quad L^\dagger = \frac{1}{N-1} \begin{bmatrix} N-2 & N-3 & \cdots & 1 \\ -1 & N-3 & \cdots & 1 \\ -1 & -2 & \cdots & 1 \\ \vdots & \vdots & \ddots & \vdots \\ -1 & -2 & \cdots & -(N-2) \end{bmatrix}.$$

252 *Proof.* The two first properties are straightforward. In order to compute the
 253 explicit formula of L^\dagger , we proceed as follows: The action of the pseudo-inverse of L
 254 on $\mathbf{x} = [x_1, \dots, x_{N-2}]^T$ is given by $L^\dagger \mathbf{x} = \mathbf{v}_1 + \mathbf{v}_2$ where $L\mathbf{v}_1 = \mathbf{x}$, $L\mathbf{v}_2 = 0$, and
 255 $\|L^\dagger \mathbf{x}\|_2$ is minimal. Furthermore, since $\text{rank } L = N - 2$, we know by the rank-nullity
 256 theorem that the null space of L is spanned by $[1, 1, \dots, 1, 1]^T$. Hence there exists
 257 $\alpha \in \mathbb{R}$ such that $\mathbf{v}_2 = \alpha[1, 1, \dots, 1, 1]^T$. Moreover, it is straightforward to verify that
 258 the vector

$$259 \quad \mathbf{v}_1 = \begin{bmatrix} x_1 + \dots + x_{N-2} \\ x_2 + \dots + x_{N-2} \\ \vdots \\ x_{N-2} \\ 0 \end{bmatrix}$$

260 is solution of $L\mathbf{v}_1 = \mathbf{x}$. Therefore, we have

$$261 \quad L^\dagger \mathbf{x} = \begin{bmatrix} x_1 + \dots + x_{N-2} \\ x_2 + \dots + x_{N-2} \\ \vdots \\ x_{N-2} \\ 0 \end{bmatrix} + \alpha \begin{bmatrix} 1 \\ 1 \\ \vdots \\ 1 \\ 1 \end{bmatrix} \in \mathbb{R}^{N-1},$$

262 where α is such that $\|L^\dagger \mathbf{x}\|_2$ is minimal. Hence,

$$263 \quad \alpha = -\frac{1}{N-1} [1 \quad 1 \quad \cdots \quad 1 \quad 1] \begin{bmatrix} x_1 + \dots + x_{N-2} \\ x_2 + \dots + x_{N-2} \\ \vdots \\ x_{N-2} \\ 0 \end{bmatrix},$$

264 and we deduce that

$$\begin{aligned}
L^\dagger \mathbf{x} &= \begin{bmatrix} 1 & 1 & \dots & 1 & 1 \\ 0 & 1 & \dots & 1 & 1 \\ \vdots & \ddots & \ddots & 0 & 1 \\ 0 & \dots & \dots & 0 & 0 \end{bmatrix} \mathbf{x} - \frac{1}{N-1} \begin{bmatrix} 1 & 1 & \dots & 1 & 1 \\ 1 & 1 & \dots & 1 & 1 \\ \vdots & \ddots & \ddots & 1 & 1 \\ 1 & \dots & \dots & 1 & 1 \end{bmatrix} \begin{bmatrix} 1 & 1 & \dots & 1 & 1 \\ 0 & 1 & \dots & 1 & 1 \\ \vdots & \ddots & \ddots & 0 & 1 \\ 0 & \dots & \dots & 0 & 0 \end{bmatrix} \mathbf{x}, \\
&= \frac{1}{N-1} \begin{bmatrix} N-2 & -1 & \dots & -1 & -1 \\ -1 & N-2 & \dots & -1 & -1 \\ \vdots & \ddots & \ddots & N-2 & -1 \\ -1 & \dots & \dots & -1 & N-2 \end{bmatrix} \begin{bmatrix} 1 & 1 & \dots & 1 & 1 \\ 0 & 1 & \dots & 1 & 1 \\ \vdots & \ddots & \ddots & 0 & 1 \\ 0 & \dots & \dots & 0 & 0 \end{bmatrix} \mathbf{x}, \\
&= \frac{1}{N-1} \begin{bmatrix} N-2 & N-3 & \dots & 1 \\ -1 & N-3 & \dots & 1 \\ -1 & -2 & \dots & 1 \\ \vdots & \vdots & \ddots & \vdots \\ -1 & -2 & \dots & -(N-2) \end{bmatrix} \mathbf{x},
\end{aligned}$$

265

266 which concludes the proof. \square

267 The addition of the auxiliary variables \hat{b}_i allows ψ_i^n to be defined but up to a
268 constant. Moreover, the iterates $u_i^{n+\frac{1}{2}}$ have continuous normal derivatives on the
269 interfaces, but discontinuous Dirichlet traces. Hence, the constants on the interior of
270 the subdomains should be chosen such that they minimize the Dirichlet jump on the
271 interfaces. Therefore, we choose $\hat{\mathbf{c}} = [0, \hat{c}_2, \dots, \hat{c}_{N-1}, 0]^T$ such that

$$(3.11) \quad \hat{\mathbf{c}} := \arg \min_{\mathbf{c}} \sum_{i=1}^{N-1} \left| (u_{i+1}^{n+\frac{1}{2}} + c_{i+1}) - (u_i^{n+\frac{1}{2}} + c_i) \right|^2.$$

272

273 The NNM for $\mathcal{L} = -\Delta$ in 1D is then given in Algorithm 3.2.

Algorithm 3.2 NNM for $\mathcal{L} = -\Delta$ in 1D

1. Set g_i^0 to zero or any inexpensive initial guess at the interfaces Γ_i
2. Repeat until convergence
 - (a) Solve the Dirichlet problems

$$\begin{aligned} \eta u_i^n - \partial_{xx} u_i^n &= f, \quad \text{in } \Omega_i, \\ u_i^n(x_{i-1}) &= g_{i-1}^n, \quad u_i^n(x_i) = g_i^n. \end{aligned}$$

- (b) Solve the Neumann problems

$$\begin{aligned} -\partial_{xx} \psi_i^n &= 0 \quad \text{in } \Omega_i, \\ \partial_x \psi_i^n(x_{i-1}) &= -a_{i-1} + \hat{b}_{i-1}, \quad \partial_x \psi_i^n(x_i) = a_i + \hat{b}_i, \quad \int_{\Omega_i} \psi_i^n = 0, \end{aligned}$$

where the \hat{b}_i are defined in (3.9).

- (c) Set

$$u_i^{n+1/2} := u_i^n - \psi_i^n.$$

- (d) Set

$$\tilde{u}_i^{n+\frac{1}{2}} := u_i^{n+\frac{1}{2}} + \hat{c}_i,$$

where the \hat{c}_i are defined in (3.11).

- (e) Set

$$g_i^{n+1} := (\tilde{u}_{i+1}^{n+\frac{1}{2}}(x_i) + \tilde{u}_i^{n+\frac{1}{2}}(x_i))/2.$$

274 LEMMA 3.3. Let $\mathbf{g}^n = [g_1^n, g_2^n, \dots, g_{N-1}^n]^T \in \mathbb{R}^{N-1}$, then for $N \geq 2$, we have
 275 $\mathbf{g}^n = \tilde{T} \mathbf{g}^{n-1}$, where $\tilde{T} \in \mathbb{R}^{(N-1) \times (N-1)}$ is given by

$$276 \quad \tilde{T} := -\frac{1}{4} \begin{bmatrix} \frac{N-2}{(N-1)^2} & 0 & \cdots & 0 & -\frac{N-2}{(N-1)^2} \\ \frac{N-4}{(N-1)^2} & 0 & \cdots & 0 & -\frac{N-4}{(N-1)^2} \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ -\frac{N-4}{(N-1)^2} & 0 & \cdots & 0 & \frac{N-4}{(N-1)^2} \\ -\frac{N-2}{(N-1)^2} & 0 & \cdots & 0 & \frac{N-2}{(N-1)^2} \end{bmatrix}.$$

277 Moreover, we have for all $N \geq 2$

$$278 \quad \rho(\tilde{T}) = \|\tilde{T}\|_\infty = \frac{N-2}{2(N-1)^2} < 1.$$

279 *Proof.* The subdomain solutions in the Laplace case are

$$280 \quad u_i^n(x) = g_{i-1}^n + \frac{g_i^n - g_{i-1}^n}{H}(x - x_{i-1}) \quad \text{for } i = 1, \dots, N,$$

281 where we set $g_0^n := g_N^n := 0$ for simplicity. Then, the a_i are given by

$$282 \quad a_i := \frac{2g_i^n - g_{i-1}^n - g_{i+1}^n}{2H}, \quad \text{for } i = 1, \dots, N-1.$$

283 It thus follows that $\hat{\mathbf{b}}$ is given by

$$284 \quad \hat{\mathbf{b}} := L^\dagger D \mathbf{g},$$

285 where

$$286 \quad D := \frac{1}{2H} \begin{bmatrix} 1 & 1 & -1 & 0 & \dots & \dots & 0 \\ -1 & 1 & 1 & -1 & 0 & \dots & 0 \\ 0 & -1 & 1 & 1 & -1 & \dots & 0 \\ \vdots & \ddots & \ddots & \ddots & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & -1 & 1 & 1 & -1 \\ 0 & \dots & \dots & 0 & -1 & 1 & 1 \end{bmatrix}.$$

287 Using [Lemma 3.2](#), we have

$$288 \quad L^\dagger D = \frac{1}{2H} \begin{bmatrix} \frac{1}{N-1} & 1 & 0 & 0 & \dots & -\frac{1}{N-1} \\ -\frac{N-2}{N-1} & 0 & 1 & 0 & \dots & \vdots \\ \frac{1}{N-1} & -1 & 0 & 1 & \dots & \vdots \\ \vdots & 0 & \ddots & \ddots & \ddots & -\frac{1}{N-1} \\ \frac{1}{N-1} & \vdots & \ddots & -1 & 0 & \frac{N-2}{N-1} \\ \frac{1}{N-1} & \dots & \dots & 0 & -1 & -\frac{1}{N-1} \end{bmatrix},$$

289 and hence

$$290 \quad \hat{b}_i := \frac{1}{2H} \left(\frac{1}{N-1} g_1^n + g_{i+1}^n - g_{i-1}^n - \frac{1}{N-1} g_{N-1}^n \right).$$

291 The formula for the local corrections ψ_i^n therefore becomes

$$292 \quad \psi_i^n(x) = \frac{1}{2H} \left(\frac{1}{N-1} g_1^n + 2g_i^n - 2g_{i-1}^n - \frac{1}{N-1} g_{N-1}^n \right) \left(x - \frac{x_i + x_{i-1}}{2} \right), \quad i = 2, \dots, N-1,$$

293 and for the first and last subdomains

$$294 \quad \begin{aligned} \psi_1^n(x) &= \frac{1}{2H} \left(\frac{2N-1}{N-1} g_1^n - \frac{1}{N-1} g_{N-1}^n \right) (x - x_0), \\ \psi_N^n(x) &= \frac{1}{2H} \left(-\frac{1}{N-1} g_1^n + \frac{2N-1}{N-1} g_{N-1}^n \right) (x_N - x). \end{aligned}$$

295 We then have

$$296 \quad \begin{aligned} u_i^{n+\frac{1}{2}}(x_i) &= \frac{1}{4} \left(-\frac{1}{N-1} g_1^n + 2g_i^n + 2g_{i-1}^n + \frac{1}{N-1} g_{N-1}^n \right), \\ u_i^{n+\frac{1}{2}}(x_{i-1}) &= \frac{1}{4} \left(\frac{1}{N-1} g_1^n + 2g_i^n + 2g_{i-1}^n - \frac{1}{N-1} g_{N-1}^n \right), \end{aligned}$$

297 and for the first and last subdomains

$$298 \quad \begin{aligned} u_1^{n+\frac{1}{2}}(x_1) &= \frac{1}{2} \left(-\frac{1}{N-1} g_1^n + \frac{1}{N-1} g_{N-1}^n \right), \\ u_N^{n+\frac{1}{2}}(x_{N-1}) &= \frac{1}{2} \left(\frac{1}{N-1} g_1^n - \frac{1}{N-1} g_{N-1}^n \right), \end{aligned}$$

299 which gives

$$300 \quad u_{i+1}^{n+\frac{1}{2}}(x_i) - u_i^{n+\frac{1}{2}}(x_i) = \frac{1}{2} \left(\frac{1}{N-1} g_1^n + g_{i+1}^n - g_{i-1}^n - \frac{1}{N-1} g_{N-1}^n \right), \quad i = 2, \dots, N-2,$$

301 and for the first and last subdomains, we have

$$302 \quad \begin{aligned} u_2^{n+\frac{1}{2}}(x_1) - u_1^{n+\frac{1}{2}}(x_1) &= \frac{1}{4} \left(\frac{2N+1}{N-1} g_1^n + 2g_2^n - \frac{3}{N-1} g_{N-1}^n \right), \\ u_N^{n+\frac{1}{2}}(x_{N-1}) - u_{N-1}^{n+\frac{1}{2}}(x_{N-1}) &= \frac{1}{4} \left(\frac{3}{N-1} g_1^n - 2g_{N-2}^n - \frac{2N+1}{N-1} g_{N-1}^n \right). \end{aligned}$$

303 Since \mathbf{c} is given by

$$304 \quad c_i = -\frac{N-i}{N-1} \sum_{k=1}^{i-1} \left(u_{k+1}^{n+\frac{1}{2}}(x_k) - u_k^{n+\frac{1}{2}}(x_k) \right) + \frac{i-1}{N-1} \sum_{k=i}^{N-1} \left(u_{k+1}^{n+\frac{1}{2}}(x_k) - u_k^{n+\frac{1}{2}}(x_k) \right),$$

305 and

$$\begin{aligned} \sum_{k=1}^{i-1} \left(u_{k+1}^{n+\frac{1}{2}}(x_k) - u_k^{n+\frac{1}{2}}(x_k) \right) &= u_2^{n+\frac{1}{2}}(x_1) - u_1^{n+\frac{1}{2}}(x_1) + \sum_{k=2}^{i-1} \left(u_{k+1}^{n+\frac{1}{2}}(x_k) - u_k^{n+\frac{1}{2}}(x_k) \right) \\ &= \frac{1}{4} \left(\frac{2N+1}{N-1} g_1^n + 2g_2^n - \frac{3}{N-1} g_{N-1}^n \right) \\ &\quad + \frac{1}{2} \sum_{k=2}^{i-1} \left(\frac{1}{N-1} g_1^n + g_{k+1}^n - g_{k-1}^n - \frac{1}{N-1} g_{N-1}^n \right) \\ 306 \quad &= \frac{1}{4} \left(\frac{2N+1}{N-1} g_1^n + 2g_2^n - \frac{3}{N-1} g_{N-1}^n \right) \\ &\quad + \frac{1}{2} \left(\frac{i-2}{N-1} g_1^n + g_i^n + g_{i-1}^n - g_2^n - g_1^n - \frac{i-2}{N-1} g_{N-1}^n \right) \\ &= \frac{1}{4} \left(\frac{2i-1}{N-1} g_1^n + 2g_i^n + 2g_{i-1}^n - \frac{2i-1}{N-1} g_{N-1}^n \right), \end{aligned}$$

307 and similarly,

$$\begin{aligned} \sum_{k=i}^{N-1} \left(u_{k+1}^{n+\frac{1}{2}}(x_k) - u_k^{n+\frac{1}{2}}(x_k) \right) &= \sum_{k=i}^{N-2} \left(u_{k+1}^{n+\frac{1}{2}}(x_k) - u_k^{n+\frac{1}{2}}(x_k) \right) + u_N^{n+\frac{1}{2}}(x_{N-1}) - u_{N-1}^{n+\frac{1}{2}}(x_{N-1}) \\ &= \frac{1}{2} \sum_{k=i}^{N-2} \left(\frac{1}{N-1} g_1^n + g_{k+1}^n - g_{k-1}^n - \frac{1}{N-1} g_{N-1}^n \right) \\ 308 \quad &+ \frac{1}{4} \left(\frac{3}{N-1} g_1^n - 2g_{N-2}^n - \frac{2N+1}{N-1} g_{N-1}^n \right) \\ &= \frac{1}{2} \left(\frac{N-i-1}{N-1} g_1^n + g_{N-1}^n + g_{N-2}^n - g_i^n - g_{i-1}^n - \frac{N-i-1}{N-1} g_{N-1}^n \right) \\ &\quad + \frac{1}{4} \left(\frac{3}{N-1} g_1^n - 2g_{N-2}^n - \frac{2N+1}{N-1} g_{N-1}^n \right) \\ &= \frac{1}{4} \left(\frac{2(N-i)+1}{N-1} g_1^n - 2g_i^n - 2g_{i-1}^n - \frac{2(N-i)+1}{N-1} g_{N-1}^n \right), \end{aligned}$$

309 we deduce for the components of \mathbf{c} the formula

$$310 \quad c_i = -\frac{1}{4} \left(\frac{N-2i+1}{(N-1)^2} g_1^n + 2g_{i-1}^n + 2g_i^n - \frac{N-2i+1}{(N-1)^2} g_{N-1}^n \right), \quad i = 2, \dots, N-2,$$

311 and hence

$$312 \quad g_i^{n+1} := -\frac{1}{4} \left(\frac{N-2i}{(N-1)^2} g_1^n - \frac{N-2i}{(N-1)^2} g_{N-1}^n \right), \quad i = 2, \dots, N-2,$$

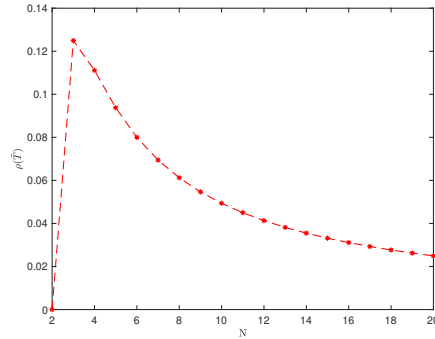


Fig. 5: Curve of the convergence factor in (3.12) with respect to N for $a = -1$, $b = 1$

313 which concludes the proof for the iteration matrix \tilde{T} . To prove that $\rho(\tilde{T}) = \|\tilde{T}\|_\infty$,
 314 define the vector $\mathbf{v} = [1 \quad \frac{N-2}{N-4} \quad \dots \quad -\frac{N-2}{N-4} \quad -1]^T$ (i.e. $v_k = \frac{N-2k}{N-2}$, $k = 1, \dots, N -$
 315 1). We clearly see that \mathbf{v} is an eigenvector of \tilde{T} associated to $\|\tilde{T}\|_\infty$. \square

316 **THEOREM 3.4.** *The two-level NNM for the Laplace problem in 1D given by Algo-*
 317 *rithm 3.2 is convergent, and satisfies the convergence estimate*

$$318 \quad (3.12) \quad \max_{1 \leq i \leq N-1} |g_i^n| \leq \left(\frac{H(b-a-2H)}{2(b-a-H)^2} \right)^n \max_{1 \leq i \leq N-1} |g_i^0|.$$

319 *Proof.* By Lemma 3.3, we have that $\rho(\tilde{T}) < 1$ for all $N \geq 2$, which shows that the
 320 algorithm is convergent. To prove the estimate, it suffice to take the infinity norm of
 321 \mathbf{g}^n and note that $N = (b-a)/H$. \square

322 We see from Theorem 3.4 that a coarse space with piecewise constant functions leads
 323 to a well posed and scalable algorithm. Indeed, Figure 5 shows that the convergence
 324 factor in (3.12) does not deteriorate for increasing N . This shows also that the
 325 constants are not sufficient for an optimal coarse space, i.e., to get a direct solver. In
 326 order to obtain a two-level algorithm that converges after the first coarse correction,
 327 one would need to choose the optimal coarse space X_d that consists of piecewise linear
 328 functions in the 1D Laplace case.

329 **3.2. Two-dimensional analysis.** We now analyze the convergence of the two-
 330 level NNM in 2D, both for the screened Laplace operator $\mathcal{L} = \eta^2 - \Delta$ and for the
 331 Laplacian $\mathcal{L} = -\Delta$.

332 **3.2.1. The screened Laplace operator $\mathcal{L} = \eta^2 - \Delta$.** We use for this case
 333 Dirichlet boundary conditions imposed on all boundaries of the original domain. In
 334 2D, the optimal coarse space X_d is now infinite dimensional, since the interface traces
 335 are now functions. The optimal coarse space is too big to be used in practice, and we
 336 need to choose finite dimensional approximation. In order to determine which func-
 337 tions to use in the approximation, we recall that in Subsection 2.2 we have expanded
 338 the subdomain solutions in a sine series based on the separation of variables approach.
 339 These sine functions are the eigenfunctions of the interface eigenvalue problem

$$340 \quad (3.13) \quad \begin{aligned} -\partial_{xx}\psi_i &= \lambda\psi_i \text{ in } \Gamma_i, \\ \psi_i &= 0 \text{ on } \partial\Gamma_i, \end{aligned}$$

341 the eigenpairs being $(\psi_i, \lambda) = \left(\sin\left(\frac{m\pi}{L}y\right), \frac{m^2\pi^2}{L^2} \right)$, $m \geq 1$. Note that interface eigen-
 342 value problems of the form (3.13) were also the essential ingredient in the construction
 343 of SHEM [19, 18, 21]. We now determine an optimized approximate coarse space using
 344 the interface eigenvalue problems in (3.13) for a decomposition as in Figure 3. Since
 345 the lowest eigenmode of (3.13) corresponds precisely to the most slowly converging
 346 mode of the NNM as we have seen in Theorem 2.4, we should use in an optimized
 347 coarse space the lowest modes of (3.13). One can then expect that the more modes we
 348 add, the more the convergence improves. This leads to our definition of the optimized
 349 spectral coarse space

$$350 \quad (3.14) \quad \tilde{X}_d := \left\{ v \in X_d : \forall i, \partial_x v_i(x_i, y) \in \text{span} \left\{ \sin\left(\frac{m\pi}{L}y\right), m = 1, \dots, J \right\} \right\},$$

351 where $J \geq 1$ can be suitably chosen to get a richer and richer coarse space, and the
 352 optimal one in the limit.

353 **THEOREM 3.5.** *The two-level NNM for the screened Laplace problem $(\eta^2 - \Delta)u =$
 354 0 in 2D given by Algorithm 3.1 with the optimized coarse space \tilde{X}_d defined in (3.14)
 355 satisfies the error bound*

$$356 \quad (3.15) \quad \max_{1 \leq i \leq N-1} \|g_i^n\|_2 \leq \frac{1}{\sinh^{2n}(k_{J+1}H)} \max_{1 \leq i \leq N-1} \|g_i^0\|_2.$$

357 *Proof.* At each iteration n , the intermediate solution $u_i^{n+1/2}$ can be written as

$$358 \quad u_i^{n+1/2}(x, y) = \sum_{m=1}^{\infty} \hat{u}_i^{n+1/2}(x, m) \sin(k_m y), \text{ for } i = 1, \dots, N.$$

359 Using Parseval's identity, we have for $\nu \in \tilde{X}_d$

$$360 \quad \int_0^L \left| (u_{i+1}^{n+\frac{1}{2}} + \nu)(x_i^+, y) - (u_i^{n+\frac{1}{2}} + \nu)(x_i^-, y) \right|^2 dy = \frac{L}{2} \sum_{m=1}^J \left| (\hat{u}_{i+1}^{n+\frac{1}{2}} + \hat{\nu})(x_i^+, m) - (\hat{u}_i^{n+\frac{1}{2}} + \hat{\nu})(x_i^-, m) \right|^2 \\ + \frac{L}{2} \sum_{m=J+1}^{\infty} \left| \hat{u}_{i+1}^{n+\frac{1}{2}}(x_i^+, m) - \hat{u}_i^{n+\frac{1}{2}}(x_i^-, m) \right|^2,$$

361 and hence

$$362 \quad U^n = \arg \min_{\nu \in \tilde{X}_d} \sum_{i=1}^{N-1} \sum_{m=1}^J \left| (\hat{u}_{i+1}^{n+\frac{1}{2}} + \hat{\nu})(x_i^+, m) - (\hat{u}_i^{n+\frac{1}{2}} + \hat{\nu})(x_i^-, m) \right|^2.$$

363 Since we are minimizing a finite dimensional quadratic, the quantity can be made
 364 zero. Hence, $g_i^{n+1}(m) = 0$ for $m \leq J$, and following the proof of Theorem 2.4, we get
 365 the stated bound. \square

366 Theorem 3.5 shows that no matter the value of H , we can obtain a convergent
 367 NNM by adding enough coarse space components, and thus turn a divergent NNM due
 368 to bad aspect ratio of the subdomains into a convergent one. Furthermore, the more
 369 we increase the number of coarse space components J , the faster the NNM becomes,
 370 and there are no other, more effective coarse space components to add, which explains
 371 the term optimized coarse space. Note also that the more modes we add, the better
 372 our approximation of the theoretically optimal coarse space X_d becomes, and in the
 373 discrete setting, we can actually reach the optimal coarse space, which is then finite
 374 dimensional as well.

375 **3.2.2. The Laplace operator $\mathcal{L} = -\Delta$ with Neumann boundary con-**
 376 **ditions.** If we have Neumann boundary conditions in 2D all around the original
 377 domain, then the local corrections ψ_i^n are not well defined as in the one-dimensional
 378 case, unless they satisfy the 2D consistency conditions

$$379 \quad (3.16) \quad \int_0^L \partial_x \psi_i^n(x_i, y) dy - \int_0^L \partial_x \psi_i^n(x_{i-1}, y) dy = 0, \text{ for } i = 2, \dots, N-1.$$

380 To obtain a well defined NNM, we generalize the ideas of [Subsection 3.1.2](#), and in-
 381 troduce new auxiliary functions $\hat{b}_i(y)$ along the interfaces $(0, L)$ such that [\(3.16\)](#) is
 382 satisfied, which is equivalent to

$$383 \quad (3.17) \quad \int_0^L \hat{b}_{i-1}(y) dy - \int_0^L \hat{b}_i(y) dy = \int_0^L a_{i-1}(y) dy + \int_0^L a_i(y) dy, \quad i = 1, \dots, N-1,$$

384 where, $a_i(y) := \frac{1}{2}(\partial_x u_i^n(x_i, y) - \partial_x u_{i+1}^n(x_i, y))$ on $(0, L)$, $i = 1, \dots, N-1$, and hence
 385 any function satisfying this condition is candidate for correcting [Algorithm 2.1](#).

386 Note that we can choose freely the $\hat{b}_i(y)$ as long as they satisfy [\(3.17\)](#). If chosen
 387 carefully, they can accelerate the convergence of the algorithm, but we will restrict
 388 ourselves first here to the case where all the auxiliary functions are kept constant
 389 along $(0, L)$. Then, using [\(3.17\)](#), we obtain the linear system

$$390 \quad (3.18) \quad L\hat{\mathbf{b}} = \hat{\mathbf{a}}$$

391 where $L \in \mathbb{R}^{N-2 \times N-1}$ is as in [Subsection 3.1.2](#), and $\hat{\mathbf{b}} \in \mathbb{R}^{N-1}$, $\hat{\mathbf{a}} \in \mathbb{R}^{N-2}$ are given
 392 by

$$393 \quad \hat{\mathbf{b}} := \begin{bmatrix} \hat{b}_1 \\ \vdots \\ \hat{b}_{N-1} \end{bmatrix}, \quad \hat{\mathbf{a}} := \begin{bmatrix} \hat{a}_1 \\ \vdots \\ \hat{a}_{N-1} \end{bmatrix},$$

394 and

$$395 \quad (3.19) \quad \hat{a}_i = \frac{1}{L} \left(\int_0^L a_i(y) dy + \int_0^L a_{i+1}(y) dy \right), \quad i = 1, \dots, N-2.$$

396 Again, as in the one dimensional case, the solutions of ψ_i are not unique, and we
 397 need to choose the most suitable constants. Following the ideas in [Subsection 3.1.2](#),
 398 we choose $\hat{\mathbf{c}} = [0, \hat{c}_2, \dots, \hat{c}_{N-1}, 0]^T$ such that

$$399 \quad (3.20) \quad \hat{\mathbf{c}} := \arg \min_{\mathbf{c}} \sum_{i=1}^{N-1} \int_0^L \left| (u_{i+1}^{n+\frac{1}{2}}(x_i, y) + c_{i+1}) - (u_i^{n+\frac{1}{2}}(x_i, y) + c_i) \right|^2 dy.$$

400 We thus obtain the two-level NNM for the Laplace case in 2D in [Algorithm 3.3](#).

401 **LEMMA 3.6.** *Let $\hat{\mathbf{g}}^n(m) = [\hat{g}_1^n(m), \hat{g}_2^n(m), \dots, \hat{g}_{N-1}^n(m)]^T \in \mathbb{R}^{N-1}$. Then for*
 402 *$N \geq 3$ we have*

$$403 \quad (3.22) \quad \begin{aligned} \hat{\mathbf{g}}^n(0) &= \tilde{T} \hat{\mathbf{g}}^{n-1}(0), & \text{for } m = 0, \\ \hat{\mathbf{g}}^n(m) &= T_m \hat{\mathbf{g}}^{n-1}(m), & \text{for } m \geq 1, \end{aligned}$$

404 where \tilde{T} and T_m are defined as in [Lemmas 2.1](#) and [2.3](#).

Algorithm 3.3 NNM for $\mathcal{L} = -\Delta$ in 2D

1. Set g_i^0 to zero or any inexpensive initial guess on the interfaces Γ_i
2. Repeat until convergence
 - (a) Solve the Dirichlet problems

$$(3.21) \quad \begin{aligned} -\Delta u_i^n &= f, \quad \text{in } \Omega_i, \\ u_i^n(x_{i-1}, y) &= g_{i-1}^n(y), u_i^n(x_i, y) = g_i^n(y) \quad \text{on } (0, L), \\ \partial_y u_i^n(x, 0) &= \partial_y u_i^n(x, L) = 0, \quad \text{on } (a, b). \end{aligned}$$

- (b) Solve the Neumann problems

$$\begin{aligned} -\Delta \psi_i^n &= 0, \quad \text{in } \Omega_i, \\ \partial_x \psi_i^n(x_{i-1}, y) &= -a_{i-1}(y) + \hat{b}_{i-1}(y), \quad \text{on } (0, L), \\ \partial_x \psi_i^n(x_i, y) &= a_i(y) + \hat{b}_i(y), \quad \text{on } (0, L), \\ \partial_y \psi_i^n(x, 0) &= \partial_y \psi_i^n(x, L) = 0, \quad \text{on } (a, b), \\ \int_{\Omega_i} \psi_i^n &= 0, \end{aligned}$$

where the $\hat{b}_i(y)$ are defined in (3.17).

- (c) Set for $i = 1, \dots, N$

$$\tilde{u}_i^{n+\frac{1}{2}} := u_i^{n+\frac{1}{2}} + \hat{c}_i,$$

where the \hat{c}_i are defined in (3.20).

- (d) Set for $i = 1, \dots, N$

$$g_i^{n+1}(y) := \left(\tilde{u}_{i+1}^{n+\frac{1}{2}}(x_i, y) + \tilde{u}_i^{n+\frac{1}{2}}(x_i, y) \right) / 2, \quad \text{on } (0, L).$$

405 *Proof.* In the case of Neumann conditions, u_i^n and ψ_i^n can be expanded in a cosine
406 series,

$$407 \quad u_i^n(x, y) = \sum_{m=0}^{\infty} \hat{u}_i^n(x, m) \cos(k_m y), \quad \psi_i^n(x, y) = \sum_{m=0}^{\infty} \hat{\psi}_i^n(x, m) \cos(k_m y),$$

408 and we thus have as before a sequence of one dimensional problems for each mode m .
409 For the case $m = 0$, the Fourier coefficients $\hat{u}_i^n(x, 0)$, $\hat{\psi}_i^n(x, 0)$ satisfy

$$410 \quad \begin{aligned} -\partial_{xx} \hat{u}_i^n(x, 0) &= 0, & -\partial_{xx} \hat{\psi}_i^n(x, 0) &= 0, \\ \hat{u}_i^n(x_{i-1}, 0) &= \hat{g}_{i-1}^n(0), & \hat{\psi}_i^n(x_{i-1}, 0) &= a_{i-1}(0), \\ \hat{u}_i^n(x_i, 0) &= \hat{g}_i^n(0), & \hat{\psi}_i^n(x_i, 0) &= a_i(0), \end{aligned}$$

411 where $a_i(m) = (\partial_x \hat{u}_i^n(x_i, m) - \partial_x \hat{u}_{i+1}^n(x_i, m)) / 2$, $m \geq 0$, $i = 1, \dots, N-1$. From (3.19),
412 we find that $\hat{a}_i = a_i(0) + a_{i+1}(0)$, $i = 1, \dots, N-2$, and since $\int_{\Omega_i} u_i(x, y) dx dy =$
413 $\int_{x_{i-1}}^{x_i} \int_0^L u_i(x, y) = 0$, we deduce that $\int_{x_{i-1}}^{x_i} \hat{u}_i^n(x, 0) dx = 0$, $i = 2, \dots, N-1$. Moreover,

414 we remark that \hat{c} satisfies

$$\begin{aligned}
415 \quad \hat{c} &= \operatorname{argmin}_c \sum_{i=1}^{N-1} \int_0^L \left| (u_{i+1}^{n+\frac{1}{2}}(x_i, y) + c_{i+1}) - (u_i^{n+\frac{1}{2}}(x_i, y) + c_i) \right|^2 dy, \\
&= \operatorname{argmin}_c \sum_{i=1}^{N-1} L \left| (\hat{u}_{i+1}^{n+\frac{1}{2}}(x_i, 0) + c_{i+1}) - (\hat{u}_i^{n+\frac{1}{2}}(x_i, 0) + c_i) \right|^2 + \frac{L}{2} \sum_{m=1}^{\infty} \left| \hat{u}_{i+1}^{n+\frac{1}{2}}(x_i, m) - \hat{u}_i^{n+\frac{1}{2}}(x_i, m) \right|^2, \\
&= \operatorname{argmin}_c \sum_{i=1}^{N-1} \left| (\hat{u}_{i+1}^{n+\frac{1}{2}}(x_i, 0) + c_{i+1}) - (\hat{u}_i^{n+\frac{1}{2}}(x_i, 0) + c_i) \right|^2.
\end{aligned}$$

416 Hence, the Fourier coefficients $\hat{u}_i^n(x, 0)$ and $\hat{\psi}_i^n(x, 0)$ have the same iterates as in the 1D
417 Algorithm 3.2 for Laplace equation. Using then Lemma 3.6, we get the first recurrence
418 relation of (3.22). For $m \geq 1$, the iterates $\hat{u}_i^n(x, m)$ and $\hat{\psi}_i^n(x, m)$ are the same as in
419 the 1D Algorithm 2.1 for the screened Laplace equation, and using Lemma 2.3, we
420 get the second recurrence relation of (3.22). \square

421 **THEOREM 3.7.** *If the ratio of subdomain width H and height L satisfies $\frac{H}{L} >$
422 $\frac{\ln(1+\sqrt{2})}{\pi}$, then the two-level NNM for the Laplace problem with Neumann boundary
423 conditions in 2D given by Algorithm 3.3 converges, and satisfies the error estimate*

$$424 \quad (3.23) \quad \max_{1 \leq i \leq N-1} \|g_i^n\|_2 \leq \max \left\{ \frac{H(b-a-2H)}{2(b-a-H)^2}, \frac{1}{\sinh^2(k_1 H)} \right\}^n \max_{1 \leq i \leq N-1} \|g_i^0\|_2.$$

425 *Proof.* From Lemma 3.6, we have that Algorithm 3.3 converges iff $\rho(\tilde{T}) < 1$, and
426 $\rho(\tilde{T}_m) < 1$ for $m \geq 1$. Since we already know that $\rho(\tilde{T}) < 1$ as shown in Lemma 3.3,
427 it suffices that $\rho(\tilde{T}_m) \leq \|T_m\|_\infty \leq \|T_1\|_\infty < 1$, which is satisfied if $\sinh(k_1 H) > 1$,
428 or equivalently $k_1 H > \ln(1 + \sqrt{2})$. To show the error bound, it suffices to use again
429 Parseval's identity $\|g_i^n\|_2^2 = L \hat{g}_i(0)^2 + \frac{L}{2} \sum_{m=1}^{\infty} \hat{g}_i(m)^2$, and follow the same steps as
430 in the proof of Theorem 2.4. \square

431 **Theorem 3.7** shows that piecewise constant functions are sufficient to have a well
432 defined and convergent iterative NNM provided an assumption on the aspect ratio of
433 the subdomain geometry is verified.

434 **3.3. Three-dimensional analysis.** In this part, we analyze the convergence
435 of the two-level NNM for the three-dimensional screened Laplace operator $\mathcal{L} = \eta^2 -$
436 Δ . As in the two-dimensional case, the optimal coarse space X_d defined in (3.2)
437 is of infinite dimension. We propose to approximate this coarse space using the
438 eigenvalue problem defined in Equation (3.13), except that now the interfaces Γ_i
439 represent surfaces in the yz -plane. The solution of this problem is given by the
440 eigenpairs $(\psi_i, \lambda) = \left(\sin\left(\frac{m\pi}{L}y\right) \sin\left(\frac{m'\pi}{L'}z\right), \frac{m^2\pi^2}{L^2} + \frac{m'^2\pi^2}{L'^2} \right)$. This yields the definition
441 of the optimized coarse space \tilde{X}_d in 3D,
(3.24)

$$442 \quad \tilde{X}_d := \left\{ v \in X_d : \forall i, \partial_x v_i(x_i, y, z) \in \operatorname{span} \left\{ \sin\left(\frac{m\pi}{L}y\right) \sin\left(\frac{m'\pi}{L'}z\right), (m, m') \in \mathcal{I}_{J,J'} \right\} \right\},$$

443 where $\mathcal{I}_{J,J'} := \{1, \dots, J\} \times \{1, \dots, J'\}$, and $J, J' \geq 1$ are positive integers that are
444 used to enrich the approximate coarse space.

445 **THEOREM 3.8.** *The two-level NNM for the screened Laplace problem $(\eta^2 - \Delta)u =$
446 0 in 3D given by Algorithm 3.1 with the optimized coarse space \tilde{X}_d defined in (3.24)*

447 *satisfies the error bound*

$$448 \quad (3.25) \quad \max_{1 \leq i \leq N-1} \|g_i^n\|_2 \leq \frac{1}{\sinh^{2n}(k^* H)} \max_{1 \leq i \leq N-1} \|g_i^0\|_2,$$

$$449 \quad \text{where } k^* := \min_{(m,m') \in \mathbb{Z}_{>0}^2 \setminus \mathcal{I}_{J,J'}} k_{m,m'}.$$

450 *Proof.* We proceed as in the proof of [Theorem 3.5](#). The intermediate solution
451 $u_i^{n+1/2}$ can be written as

$$452 \quad u_i^{n+1/2}(x, y, z) = \sum_{m,m' \geq 1} \hat{u}_i^{n+1/2}(x, m) \sin\left(\frac{m\pi}{L} y\right) \sin\left(\frac{m'\pi}{L'} z\right), \quad \text{for } i = 1, \dots, N.$$

453 Using Parseval's identity, we have for $\nu \in \tilde{X}_d$

$$\begin{aligned} & \iint_{[0,L] \times [0,L']} \left| (u_{i+1}^{n+\frac{1}{2}} + \nu)(x_i^+, y, z) - (u_{i+1}^{n+\frac{1}{2}} + \nu)(x_i^-, y, z) \right|^2 dy dz = \\ & = \frac{LL'}{4} \sum_{(m,m') \in \mathcal{I}_{J,J'}} \left| (\hat{u}_{i+1}^{n+\frac{1}{2}} + \hat{\nu})(x_i^+, m, m') - (\hat{u}_{i+1}^{n+\frac{1}{2}} + \hat{\nu})(x_i^-, m, m') \right|^2 \\ & + \frac{LL'}{4} \sum_{(m,m') \in \mathbb{Z}_{>0}^2 \setminus \mathcal{I}_{J,J'}} \left| \hat{u}_{i+1}^{n+\frac{1}{2}}(x_i^+, m, m') - \hat{u}_{i+1}^{n+\frac{1}{2}}(x_i^-, m, m') \right|^2, \end{aligned}$$

455 and hence

$$456 \quad U^n = \arg \min_{\nu \in \tilde{X}_d} \sum_{i=1}^{N-1} \sum_{(m,m') \in \mathcal{I}_{J,J'}} \left| (\hat{u}_{i+1}^{n+\frac{1}{2}} + \hat{\nu})(x_i^+, m, m') - (\hat{u}_{i+1}^{n+\frac{1}{2}} + \hat{\nu})(x_i^-, m, m') \right|^2.$$

457 Since we are minimizing a finite dimensional quadratic, the quantity can be made
458 zero. Hence, $g_i^{n+1}(m, m') = 0$ for $m \leq J$, $m' \leq J'$, and following the proof of [Theorem 2.6](#),
459 we get the stated bound. \square

460 **4. Numerical experiments without cross points.** We illustrate now our
461 convergence results with numerical experiments. We start with the one dimensional
462 case of the screened Poisson problem $(\eta^2 - \Delta)u = f$ on the domain $\Omega := (-1, 1)$ with
463 $\eta = \sqrt{2}$ and $f(x) = 1$. We use centered finite differences with mesh size $\Delta x = 10^{-4}$,
464 and run the one-level algorithm with $N = 20$ subdomains. [Figure 6](#) (left) shows that
465 without coarse correction the algorithm fails to converge, and with the optimal coarse
466 space we get convergence after one coarse correction on the interfaces, which implies
467 convergence after the second subdomain solve in volume in our implementation. In
468 [Figure 6](#) (right) we show the convergence behavior for the Poisson case, where the
469 NNM needs already a constant coarse space to be well posed. This method is also
470 convergent, but we see that the constants are not sufficient to have convergence after
471 the one coarse correction: the optimal coarse space consists of linear functions, and
472 leads to convergence after one coarse correction (dashed curve).

473 In 2D, we decompose the domain $\Omega := (-1, 1) \times (0, 1)$ into $N = 10$ subdomains,
474 and run Algorithm [3.1](#) on the screened Poisson problem $\eta u - \Delta u = 1$ with $\eta = 2$, dis-
475 cretized by centered finite differences using the mesh size $\Delta x = \Delta y = 5 \cdot 10^{-3}$. We can
476 see in [Figure 7](#) (left) that NNM without coarse correction is divergent. Our proposed
477 approximations of the optimal coarse space make the iterative NNM convergent, and

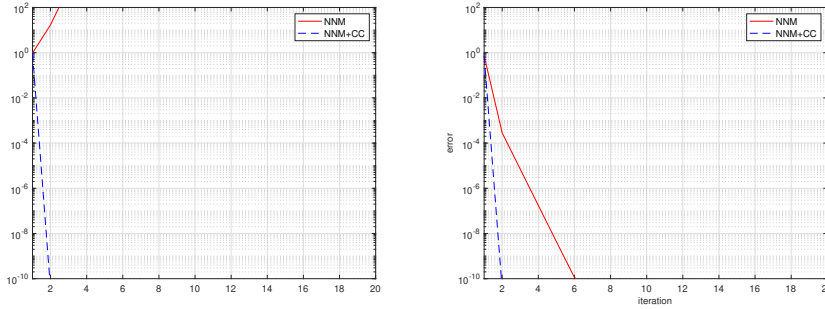


Fig. 6: 1D case: error as a function of iteration for the one and new optimal two-level NNM for the screened Poisson problem with $\eta = 2$ (left), and for the classical two-level NNM with constant coarse space and the new optimal two-level NNM for the Poisson problem (right)

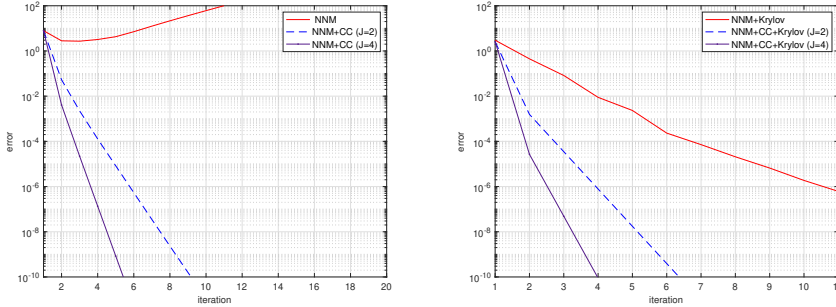


Fig. 7: Left: Convergence of NNM and coarse corrected NNM in 2D. Right: Same but using Krylov (GMRES)

478 the more we enrich the optimized coarse space \tilde{X}_d , the faster the convergence becomes.
 479 On the right in [Figure 7](#) we show the corresponding results with Krylov acceleration.
 480 We see that GMRES also makes the one-level NNM convergent, but our new coarse
 481 corrected NNM still performs better, both as iterative solver and as preconditioner.
 482 We further observe that with enough enrichment of the coarse space, our new method
 483 as an iterative solver is almost as fast as when used as a preconditioner. In 3D, we
 484 decompose the domain $\Omega := (0, 1) \times (0, 1) \times (0, 1)$ into $N = 10$ bricks, and run [Algo-](#)
 485 [rithm 3.1](#) on the screened Poisson problem $\eta u - \Delta u = 1$ with $\eta = 2$, discretized by
 486 centered finite differences using the mesh size $\Delta x = \Delta y = \Delta z = 5 \cdot 10^{-2}$. We can see
 487 from [Figure 8](#) (left) that NNM without coarse correction is divergent. The addition of
 488 the coarse correction makes the method convergent. Moreover, enriching the coarse
 489 space makes the method faster. As in the two-dimensional case, we next use the one-
 490 and two-level methods as preconditioners for a Krylov subspace method (GMRES).
 491 We observe from [Figure 8](#) (right) that GMRES makes the one-level method converge,
 492 and improves the convergence of the two-level method. Moreover, we see that in both
 493 experiments, the two-level method is faster than the one-level method, and the more
 494 we increase the size of the coarse space, the faster the two-level method becomes.

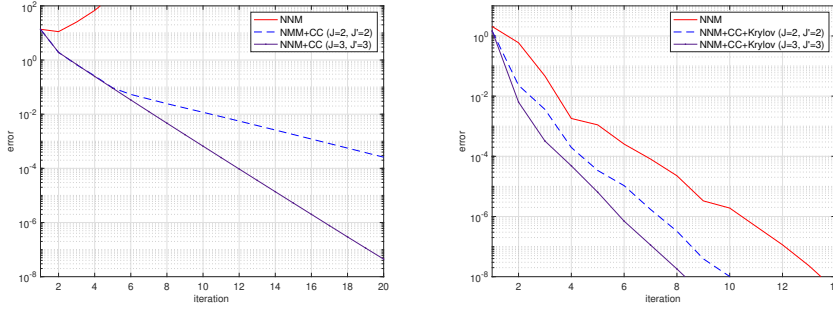


Fig. 8: Left: Convergence of NNM and coarse corrected NNM in 3D. Right: Same but using Krylov (GMRES)

495 We note also that this approach of constructing the coarse space is not restricted
 496 only to problems with constant coefficients. Indeed, the same coarse correction can
 497 be used for problems with non-constant coefficients provided that we choose the appropriate
 498 coarse functions. In general, the latter are obtained by solving an interface
 499 eigenvalue problem on the interfaces; cf. [19, 18]. For the Laplacian example, it can
 500 be shown that these eigenfunctions are exactly the sine functions; cf. [21].

501 **5. Discrete analysis of the two-level NNM including cross points.** We
 502 now present an optimal coarse space correction at the discrete level. Let \mathcal{T}_h be a
 503 triangulation of Ω such that $\mathcal{T}_{h,i} := \mathcal{T}_h \cap \bar{\Omega}_i$ is also a triangulation of Ω_i . We denote
 504 by $x_{i,j,k}$ the nodes shared by $\mathcal{T}_{h,i} \cap \mathcal{T}_{h,j}$. We define the discrete harmonic basis $\varphi_{i,j,k}$
 505 for i, j such that $\Omega_i \cap \Omega_j \neq \emptyset$ as the unique functions such that

- 506 1. their support is in the adjacent subdomains, $\text{supp } \varphi_{i,j,k} \subset \Omega_i \cup \Omega_j$,
- 507 2. their normal derivatives are continuous, $\varphi_{i,k} := \varphi_{i,j,k}|_{\Omega_i}$ and $\varphi_{j,k} := \varphi_{i,j,k}|_{\Omega_j}$
 508 verify $\frac{\partial \varphi_{i,k}}{\partial n_i}(x_{i,j,k'}) = \delta_{k,k'}$, $\frac{\partial \varphi_{j,k}}{\partial n_j}(x_{i,j,k'}) = -\delta_{k,k'}$,
- 509 3. $\varphi_{i,k}$ and $\varphi_{j,k}$ verify the discrete variational formulation of (2.1) in Ω_i and Ω_j
 510 respectively.

511 The discrete optimal coarse space X_h is then defined as

$$512 \quad (5.1) \quad X_h := \text{span}\{\varphi_{i,j,k} : \text{for all } i, j \text{ s.t. } \Omega_i \cap \Omega_j \neq \emptyset\}.$$

513 We now prove the optimality of this discrete coarse space, i.e convergence of the
 514 two-level NNM at the discrete level after one coarse correction.

515 **THEOREM 5.1.** *Algorithm 3.1 with the discrete coarse space X_h defined in (5.1)*
 516 *converges after one coarse correction.*

517 *Proof.* We proceed using an argument of dimension: consider the map

$$518 \quad T: X_h \mapsto \mathbb{R}^{d_1 + \dots + d_N}$$

$$u \mapsto (u_i(x_{i,j,k}) - u_j(x_{i,j,k})),$$

519 where $d_i := \#\{k : x_{i,j,k} \in \Omega_i \cap \Omega_j \neq \emptyset\}$, $u_i := u|_{\Omega_i}$ and $u_j := u|_{\Omega_j}$. We claim that T
 520 is a one-to-one correspondence. In fact, since the $\varphi_{i,j,k}$ are linearly independent, we
 521 have that $\dim(X_h) = d_1 + \dots + d_N$. Moreover, if $v \in \text{Ker}(T)$, then $v_i := v|_{\Omega_i}$ are
 522 continuous for both the Neumann and Dirichlet trace at the discrete level and they

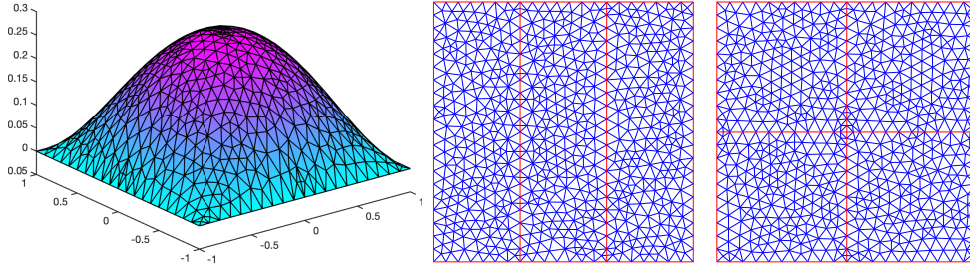


Fig. 9: Exact monodomain solution of (5.2) (left), and a 1×3 partition without a cross point of the unit square (middle), and a 2×2 partition with a cross point (right).

523 satisfy the homogeneous counterpart of (2.1) inside each subdomain, hence $v = 0$.
 524 Thus, by the rank-nullity which affirms that $\text{rank}(T) + \dim(\ker(T)) = \dim(X_h)$,
 525 we conclude that T is a one-to-one correspondence, and in particular is onto. Now,
 526 let $(U_i^0)_{1 \leq i \leq N}$ be the coarse correction, and set $\tilde{u}_i^1 := u_i^{\frac{1}{2}} + U_i^0$ the iterate after the
 527 coarse correction. The latter is chosen such that it satisfies the discrete form of the
 528 minimum jump condition in (3.3). Since T is onto, it is possible to choose U^0 in
 529 the preimage of T such that \tilde{u}_i^1 for $i = 1, \dots, N$ completely cancels the jump, i.e.,
 530 $q(\tilde{u}^1) = 0$. Therefore, the \tilde{u}_i^1 are continuous across the subdomains for both Dirichlet
 531 and Neumann traces. Moreover, the \tilde{u}_i^1 satisfy the discrete form of (2.1) inside each
 532 subdomain. It follows that \tilde{u}^1 is the discrete monodomain solution of (2.1). \square

533 We test this new, optimal two-level NNM on a variable coefficient diffusion prob-
 534 lem,

$$535 \quad (5.2) \quad \begin{aligned} -\nabla \cdot (a(x, y) \nabla u(x, y)) &= 1, \text{ if } (x, y) \in (-1, 1) \times (-1, 1), \\ u(x, y) &= 0, \text{ if } (1 - x^2)(1 - y^2) = 0, \end{aligned}$$

536 where $a(x, y) := 1 + x^2y^2$ is a continuously varying function along the interfaces.
 537 The exact monodomain solution is shown in Figure 9 (left). We choose two different
 538 decompositions as shown in Figure 9, one without cross points (middle), and one
 539 with a cross point (right). We mention that in the finite element setting, the discrete
 540 normal derivatives are computed as in [20]. The convergence results in Table 1 show
 541 that Algorithm 3.1 is truly optimal since it leads to convergence after one coarse
 542 correction both with and without cross point. For the 2×2 cross point case, Table 1
 543 shows that the one level NNM iteration diverges. The reason for this divergence was
 544 identified in [5] to be the fact that NNM is not well posed in a functional setting
 545 at the continuous level in the presence of cross points. Despite these difficulties,
 546 Algorithm 3.1 still converges in one iteration; the coarse correction is still able to
 547 extract the right traces and combine them to find the exact solution.

548 **6. Conclusion.** We designed, analyzed, and tested a new coarse space correction
 549 for the Neumann-Neumann iterative method. We described the method for general
 550 second order elliptic PDEs, and analyzed it for the Poisson and the screened Poisson
 551 equation. We proved that the new coarse space is truly optimal, i.e the method con-
 552 verges after one coarse correction. In two and three dimensions, this optimal coarse
 553 space is however too high dimensional, and we introduced an optimized spectral ap-
 554 proximation which can make an otherwise divergent iterative Neumann-Neumann

Its.	1×3 example		2×2 example with cross point	
	Algorithm 2.1	Algorithm 3.1	Algorithm 2.1	Algorithm 3.1
1	5.817e-01	2.741e-01	4.754e-01	9.582e-01
2	2.358e-01	8.881e-16	9.145e-01	1.382e-14
3	9.682e-02	4.440e-16	2.316e+00	4.996e-16
4	3.971e-02	4.440e-16	5.947e+00	4.440e-16
5	1.629e-02	4.718e-16	1.527e+01	7.216e-16
6	6.683e-03	4.996e-16	3.924e+01	5.551e-16
7	2.741e-03	4.163e-16	1.007e+02	3.885e-16

Table 1: Convergence of NNM and coarse corrected NNM for (5.2) using the optimal discrete coarse space defined in (5.1).

555 method convergent. The more we enrich our optimized coarse space, the faster the
556 convergence becomes, and it improves both as an iterative solver and as a preconditioner.
557 Our results in the discrete setting open also up the field for optimized coarse
558 spaces at cross points in Neumann-Neumann methods, and other domain decomposition
559 methods. Cross points in domain decomposition methods are currently an active
560 field of research, especially for non-overlapping domain decomposition methods: for
561 non-overlapping Schwarz methods, see e.g. [16, 17, 20, 28, 7, 6], and for the Neumann-
562 Neumann method, see [4, 5]. For classical overlapping Schwarz methods, cross points
563 are also interesting, since the partition of unity can slightly influence the convergence
564 there [13], and for the additive Schwarz preconditioner, cross points lead to the coloring
565 influencing the condition number estimate [32], and divergence when the additive
566 Schwarz method is used as a stationary iteration [12, Section 3.2], an issue that can
567 also be addressed with the coarse space, see [21].

568 Compared to the classical coarse space in the balancing Neumann-Neumann
569 method which is using a constant per subdomain, our coarse space can be made
570 richer and we showed precisely how this enhances the convergence. One then has to
571 solve however a larger coarse problem and there is thus a trade-off, like in the recent
572 algebraic coarse spaces GenEO, ACMS, or SHEM for Schwarz methods. An advantage
573 of our approximate coarse space is that it is defined a priori for constant coefficient
574 problems, and we know it at the continuous level at which the Neumann-Neumann
575 method is defined. How to deal with approximations of the optimal coarse space in
576 the variable coefficient case, and do enrichment in the presence of cross points will be
577 described in future work.

578

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