Two field-theoretic viewpoints on the Fukaya-Morse A_∞ category

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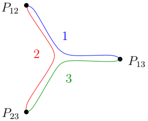
Plan

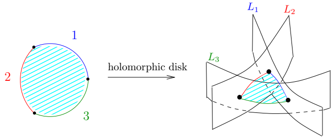
1. Fukaya-Morse A_∞ category (+“enhancement”)
2. Picture 1: effective field theory (second quantization)
3. Picture 2: HTQM (first quantization approach)

Fukaya A_∞ category and Morse degeneration

Fix X – a Riemannian manifold

	Fukaya-Morse A_∞ category on X	Fukaya A_∞category on T^*X
Ob:	$F_1, \ldots, F_N$	Lagrangians $L_a = \mathrm{graph}(\epsilon\, dF_a)$
Mor:	$\mathrm{Mor}(F_a, F_b) = \mathrm{Span}_{\mathbb{Z}}(\mathrm{Crit}(F_a - F_b))$	$\mathrm{Mor}(L_a, L_b) = \mathrm{Span}_{\mathbb{Z}}($intersection points of $L_a, L_b)$
$\circ$:	$P_{i,i+1} \in \mathrm{Crit}(F_i - F_{i+1})$ $m(P_{12}, P_{23}, \ldots, P_{N-1\,N}) = \sum_{P_{1N} \in \mathrm{Crit}(F_1 - F_N)} \#\mathcal{M}^{\mathrm{trees}}[P_{1N}]$	$p_{i,i+1} \in L_i \cap L_{i+1}$ $m(p_{12}, p_{23}, \ldots, p_{N-1\,N}) = \sum_{p_{1N} \in L_1 \cap L_N} \#\mathcal{M}^{\mathrm{hol.disks}}[p_{1N}]$



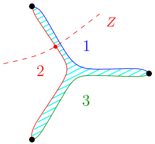


Enhancement

Enhancement: $\mathrm{Mor}(F_a, F_a) = C_\bullet(X)$
Composition maps:

$$m\colon \mathrm{Mor}(F_1, F_1)^{\otimes k_1} \otimes \mathrm{Mor}(F_1, F_2) \otimes \cdots \mathrm{Mor}(F_{N-1}, F_N) \otimes \mathrm{Mor}(F_N, F_N)^{\otimes k_N} \to \mathrm{Mor}(F_1, F_N)$$

$$m(\{Z_{1,\alpha}\}, P_{12}, \{Z_{2,\alpha}\}, \ldots, P_{N-1\,N}, \{Z_{N,\alpha}\}) = \sum_{P_{1N}} \#\mathcal{M} \cdot [P_{1N}]$$



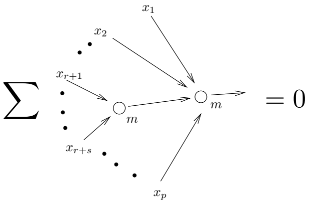
- differentials: $m(Z) = \partial Z$, $m(P_{12}) = d_{\mathrm{Morse}} P_{12}$
- $m(Z_1, Z_2) = Z_1 \cap Z_2$
- $m(Z_1, \ldots, Z_n) = 0$, $n \geq 3$

A_∞ relations

Let $x_i \in \mathrm{Mor}(F_{a_i}, F_{b_i})$, $i = 1 \dots p$, with $b_i = a_{i+1}$ - composable sequence of morphisms. Then:

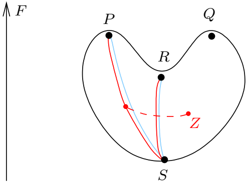
$$\sum_{r,s} m(x_1, \dots, x_r, m(x_{r+1}, \dots, x_{r+s}), x_{r+s+1}, \dots, x_p) = 0$$

Or:



Example of an A_∞ relation: heart-shaped sphere

Example: $X = S^2$, $N = 2$ functions, $F = F_1 - F_2$.



A_∞ relation: $x_1 = P \in \mathrm{Mor}(F_1, F_2)$, $x_2 = Z \in \mathrm{Mor}(F_2, F_2)$.

$$\boxed{m(d_{\mathrm{Morse}}(P), Z) + m(P, \partial Z) + d_{\mathrm{Morse}} m(P, Z) = 0}$$

Example: deformation of Morse differential by cycles

$N = 2$, $F = F_1 - F_2$. Fix $\{C_\alpha\}$ -cycles on X. $\{P_i\}$ – crit. points of F.
Generating function for compositions
$m\colon \mathrm{Mor}(F_1, F_2) \otimes \mathrm{Mor}(F_2, F_2)^{\otimes k} \to \mathrm{Mor}(F_1, F_2)$,

$$m_i^j(T) = \sum_{k\geq 1} \sum_{\alpha_1,\ldots,\alpha_k} \underbrace{\#\mathcal{M}(P_i, C_{\alpha_1}, \ldots, C_{\alpha_k}, P_j)}_{\#\text{grad traj } P_i \to P_j \text{ passing through cycles}} T_{\alpha_1} \cdots T_{\alpha_k}$$

T_α - generating parameters, $|T_\alpha| = 1 - \mathrm{codim} C_\alpha$.
A_∞ relations $\Rightarrow$ $\boxed{(d_{\mathrm{Morse}} + m(T))^2 = 0}$.

Explanation from HPT:

$$\begin{array}{c} K \circlearrowright \quad \Omega^\bullet(X), d + \sum_\alpha T_\alpha \delta_{C_\alpha} \\ p \downarrow \quad \uparrow i \\ MC(X, F) \;, d_{\mathrm{Morse}} + m(T) \end{array}$$

Morse contraction

- $i\colon \underset{\text{crit. point}}{P} \mapsto \delta_{\mathrm{Unstab}_P}$
- $p\colon \omega \mapsto \sum_P \left(\int_X \omega \wedge \delta_{\mathrm{Stab}_P} \right) \cdot [P]$
- $K = \int_0^\infty dt\, \iota_v\, e^{-t\mathcal{L}_v} : \quad \Omega^\bullet(X) \to \Omega^{\bullet-1}(X)$.
  v – gradient vector field.
  Integral kernel: $\delta_Y \in \Omega_{\mathrm{distr}}(X \times X)$;
  $Y = \{(x, y) \mid x = \mathrm{Flow}_t(v) \circ y \text{ for some } t > 0\}$

Picture 1a: homotopy transfer

$$\underline{K} \circlearrowright \quad V = \Omega^\bullet(X) \otimes \mathrm{Mat}_{N\times N} = \bigoplus_{a,b=1}^N \Omega^\bullet_{ab}(X) \quad \text{– dg algebra}$$

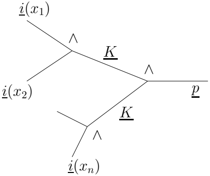
$$\underline{p}\downarrow \qquad \uparrow \underline{i}$$

$$\mathbb{M} = \bigoplus_{a,b} \mathbb{M}_{ab}, \qquad \mathbb{M}_{ab} = \begin{cases} MC(F_a - F_b), & a \neq b \\ \Omega^\bullet_{aa}, & a = b \end{cases}$$

$$\underline{i}, \underline{p}, \underline{K} = \begin{cases} \text{Morse contraction for } F_a - F_b, & a \neq b \\ \text{trivial } (\underline{i} = \underline{p} = \mathrm{id}, \underline{K} = 0), & a = b \end{cases}$$

Induced A_∞ algebra structure on $\mathbb{M}$:

$$m_n(x_1, \ldots, x_n) = \sum$$



inputs: $x_i \in \mathbb{M}_{a_i b_i}$ with $b_i = a_{i+1}$ $= \begin{cases} \text{Morse chain}, & a_i \neq b_i \\ \text{form/sing. chain } \delta_Z, & a_i = b_i \end{cases}$

Picture 1b: effective action

BF **theory:** $S = \int_X \langle B \overset{\wedge}{,} dA + \frac{1}{2}[A, A]\rangle$

Fields:

- $A \quad \in \quad \Omega^\bullet(X) \otimes \underbrace{(\mathrm{Mat}_{N\times N} \otimes \mathbb{A})}_{\mathfrak{g}}[1] \quad = \quad \bigoplus_{a,b} \Omega_{ab} \otimes \mathbb{A}[1]$

  where $\mathbb{A}$ = upper-triangular $\widetilde{N} \times \widetilde{N}$ matrices.
- $B \quad \in \quad \Omega^\bullet(X) \otimes \mathfrak{g}^*[d-2]$

Next: integrate out off-diagonal components A_{ab}, $a \neq b$ subject to gauge-fixing

$$\boxed{\iota_{v_{ab}} A_{ab} = 0}$$

– axial gauge but in different directions for different components! We induce the effective action on diagonal fields + remnants of off-diagonal fields.

Picture 1b: effective action cont'd

More explicitly:

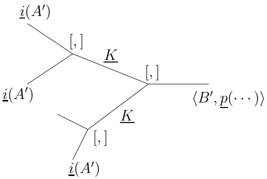
BV pushforward: $\underset{\text{fields}}{\mathcal{F}} \to \underset{\text{IR fields}}{\mathcal{F}'} = T^*[-1](\mathbb{M}\otimes\mathbb{A}[1])$

Splitting of fields: $\mathcal{F} = \mathcal{F}' \oplus \mathcal{F}''$ with $\mathcal{F}'' = \ker \underline{p} \oplus \ker \underline{i}^\vee$.

Gauge-fixing Lagrangian: $\mathcal{L} = \mathrm{im}(K) \oplus \mathrm{im}(K^\vee) \subset \mathcal{F}''$

$$e^{\frac{i}{\hbar}S_{\mathrm{eff}}(A',B')} = \int_{\mathcal{L}} \mathcal{D}A''\,\mathcal{D}B''\, e^{\frac{i}{\hbar}S(\underline{i}(A')+A'',\underline{p}^\vee(B')+B'')}$$

$$S_{\mathrm{eff}} = \sum_{\text{trees}}$$



$$= \sum_{n\geq 1} \frac{1}{n!}\langle B', l_n(A',\ldots,A')\rangle$$

$\{l_n\}$ – L_∞ algebra operations on $\mathbb{M}\otimes\mathbb{A}$.

L_∞ relations $\Leftrightarrow$ BV master equation $\{S_{\mathrm{eff}}, S_{\mathrm{eff}}\} = 0$

From L_∞ back to A_∞

Recovering A_∞ products on $\mathbb{M}$:

$$l_n(x_1 \otimes t_{12}, \ldots, x_n \otimes t_{n\,n+1}) = m_n(x_1, \ldots, x_n) \otimes t_{1\,n+1}$$

$t_{ij} \in \mathbb{A}$ matrix with (i,j)-entry 1 and all other entries 0.

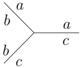
Picture 2: HTQM

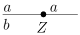
Topological quantum mechanics:

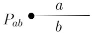
- Space of states: $\mathcal{H}_{ab} = \Omega^\bullet(X)$
  (for a particle of (a,b)-type, $a \neq b$).
- BRST operator $Q = d$.
- Hamiltonian $H = \mathcal{L}_{v_{ab}} = [Q, G]_+$.
- $G = \iota_{v_{ab}}$.
- Evolution operator (superpropagator):
  $U(t, dt) = e^{-t\,H - dt\,G} \quad \in \quad \Omega^\bullet(\mathbb{R}_+) \otimes \mathrm{End}(\mathcal{H}_{ab})$

HTQM on metric trees

HTQM on metric trees:

- 3-valent vertex 

$\sim \quad \mathcal{H}_{ab} \otimes \mathcal{H}_{bc} \xrightarrow{\wedge} \mathcal{H}_{ac}$
- 2-valent vertex 

$\sim$ operator $\mathcal{H}_{ab} \xrightarrow{\wedge \delta_Z} \mathcal{H}_{ab}$
- 1-valent vertex 

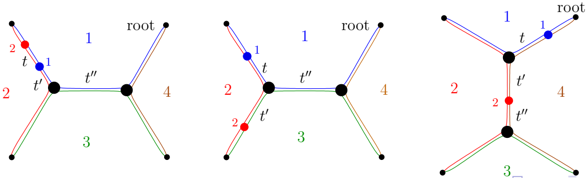
$\sim$ state $\delta_{\mathrm{Unstab}_{P_{ab}}}$
- (a, b)-edge of length $t \quad \sim \quad U_{ab}(t, dt)$

Out of these building blocks, we build a form on the space of metric trees:

$$I \in \Omega^\bullet(MT_{N;k_1,\dots,k_N}) \otimes \mathrm{Hom}(\mathrm{Mor}_{1,1}^{\otimes k_1} \otimes \mathrm{Mor}_{1,2} \otimes \cdots \otimes \mathrm{Mor}_{N-1,N} \otimes \mathrm{Mor}_{N,N}^{\otimes k_N}, \mathrm{Mor}_{1,N})$$

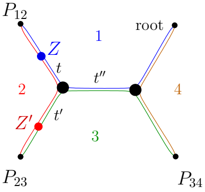
where $\mathrm{Mor}_{a,b} := \mathrm{Mor}(F_a, F_b)$.

Example: three top-cells in $MT_{4;1,1,0,0}$



Example

$$I\left(\text{[diagram]} \right) = \sum_{P_{14}\in \mathrm{Crit}(F_1-F_4)} \overline{I}\cdot [P_{14}]$$

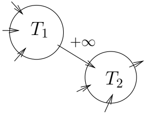


$$\overline{I} = \int_X \delta_{\mathrm{Stab}_{P_{14}}} \wedge U_{13}(t'',dt'')\Big(U_{12}(t,dt)(\delta_Z \wedge \delta_{\mathrm{Unstab}_{P_{12}}})\wedge$$

$$\wedge\, U_{23}(t',dt')(\delta_{Z'}\wedge \delta_{\mathrm{Unstab}_{P_{23}}})\Big) \wedge \delta_{\mathrm{Unstab}_{P_{34}}}$$

Properties of I

1. $(d_{MT} + Q)I = 0$
2. Factorization on IR boundary of MT:

$I\Big($ 

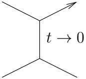
$\Big) = \left\langle I(T_2) \,\overset{\wedge}{,}\, I(T_1) \right\rangle$

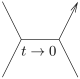
3. Period $\displaystyle\int_{MT} I = m$ – the composition map in Fukaya-Morse category.

A_∞ relations from IR factorization of HTQM

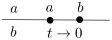
$$(d_{MT} + Q)I = 0 \qquad \Rightarrow \qquad \int_{\partial MT} I = -Q \underbrace{\int_{MT} I}_{m}$$

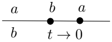
L.h.s. (contributions of boundary strata of MT):

a 

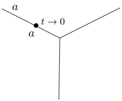
$+$ 

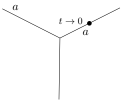
$= 0$

b 

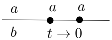
$+$ 

$= 0$

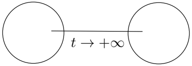
c 

$+$ 

$= 0$

d 

$\to$ terms $m(\cdots m(Z, Z') \cdots)$ in A_∞ relation.

e 

$\to$ terms $m(\cdots m(\underbrace{\cdots}_{\geq 2 \text{ colors}}) \cdots)$ in A_∞ relation.

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