Communications in Mathematical Physics manuscript No.

(will be inserted by the editor)

# An invariance principle to Ferrari-Spohn diffusions

Dmitry Ioffe<sup>1</sup> \*, Senya Shlosman<sup>2</sup>, Yvan Velenik<sup>3</sup> \*\*

- William Davidson Faculty of Industrial Engineering and Management, Technion Israel Institute of Technology, Technion City, Haifa 32000, Israel E-mail: ieioffe@technion.ac.il
- <sup>2</sup> Aix Marseille Université, Université de Toulon, CNRS, CPT UMR 7332, 13288, Marseille, France, and Inst. of the Information Transmission Problems, RAS, Moscow, Russia E-mail: senya.shlosman@univ-amu.fr
- <sup>3</sup> Section de Mathématiques, Université de Genève, 2-4, rue du Lièvre, Case postale 64, 1211 Genève

E-mail: Yvan.Velenik@unige.ch

**Abstract:** We prove an invariance principle for a class of tilted 1+1-dimensional SOS models or, equivalently, for a class of tilted random walk bridges in  $\mathbb{Z}_+$ . The limiting objects are stationary reversible ergodic diffusions with drifts given by the logarithmic derivatives of the ground states of associated singular Sturm-Liouville operators. In the case of a linear area tilt, we recover the Ferrari-Spohn diffusion with log-Airy drift, which was derived in [12] in the context of Brownian motions conditioned to stay above circular and parabolic barriers.

**Key words.** Invariance principle – critical prewetting – entropic repulsion – random walk – Ferrari-Spohn diffusions

## 1. Introduction and Results

1.1. Physical motivations. We start with an informal description, in a restricted setting, of the effective interface model at the core of our study; a detailed description in the more general framework considered in the present work will be given in Subsection 1.3.

We consider a Gibbs random field  $(X_i)_{1 \leq i \leq N}$ , with  $X_i \in \mathbb{N}$  for all i, and effective Hamiltonian

$$H_{\lambda} = \sum_{i=1}^{N-1} \Phi(X_{i+1} - X_i) + \sum_{i=1}^{N} V_{\lambda}(X_i),$$

depending on a parameter  $\lambda > 0$ . Later, we shall allow rather general forms for the interaction  $\Phi$  and for the external potential  $V_{\lambda}$ . For the purpose of this

 $<sup>^{\</sup>star}\,$  DI was supported by the Israeli Science Foundation grants 817/09 and 1723/14.

 $<sup>^{\</sup>star\star}\,$  YV was partially supported by the Swiss National Science Foundation.

introductory section, let us however restrict the discussion to the physically very relevant case of  $V_{\lambda}(x) = \lambda x$ , and assume that  $\Phi$  is symmetric and grows fast enough: for example,  $\Phi(x) = x^2$ ,  $\Phi(x) = |x|$  or  $\Phi(x) = \infty \cdot \mathbf{1}_{|x| > R}$ . Let us denote by  $\mu_{N;\lambda}$  the corresponding Gibbs measure with boundary condition  $X_1 = X_N = 0$ .

With this choice, this model can be interpreted as follows. The random variable  $X_i$  models the height of an interface above the site i. This interface separates an equilibrium phase (above the interface) and a layer of unstable phase (delimited by the interface and the wall located at height 0). The parameter  $\lambda$  corresponds to the excess free energy associated to the unstable phase.

Of course, when  $\lambda = 0$ , the distribution of X is just that of a random walk, conditioned to stay positive, with distribution of jumps given by

$$p_x = \frac{e^{-\Phi(x)}}{\sum_y e^{-\Phi(y)}}. (1.1)$$

In particular, the field delocalizes as  $N \to \infty$ . When  $\lambda > 0$ , however, the field remains localized uniformly in N. We shall be mostly interested in the behavior as  $\lambda \downarrow 0$  (say, either after letting  $N \to \infty$ , or by assuming that  $N = N(\lambda)$  grows fast enough). In that case, one can prove that the typical width of the layer is of order  $\lambda^{-1/3}$  and that the correlation along the interface is of order  $\lambda^{-2/3}$  [1,14].

In the present work, we are interested in the scaling limit of the random field X as  $\lambda \downarrow 0$ , that is, in the limiting behavior of  $x_{\lambda}(t) = \lambda^{1/3} X_{[\lambda^{-2/3}t]}$ . As stated in Theorem A below, in the particular case considered here, the scaling limit is given by the diffusion on  $(0, \infty)$  with generator

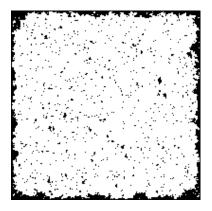
$$\frac{\sigma^2}{2}\frac{\mathrm{d}^2}{\mathrm{d}r^2} + \sigma^2 \frac{\varphi_0'}{\varphi_0} \frac{\mathrm{d}}{\mathrm{d}r} \,,$$

where (see (1.1))  $\sigma^2 = \sum_x x^2 p_x$  and  $\varphi_0 = \operatorname{Ai}(\sqrt[3]{\frac{2}{\sigma^2}} x - \omega_1)$  with  $-\omega_1$  the first zero of the Airy function Ai. This diffusion was first introduced by Ferrari and Spohn in the context of Brownian motions conditioned to stay above circular and parabolic barriers [12].

This scaling limit should be common to a wide class of systems, of which the following are but a few examples:

- Critical prewetting in the 2d Ising model: behavior of the film of unstable negatively magnetized layer induced by (-)-boundary conditions, in the presence of a positive bulk magnetic field [18]; see Figure 1.
- Interfacial adsorption at the interface between two equilibrium phases [13, 17].
- Geometry of the top-most layer of the 2 + 1-dimensional SOS model above a wall [7].
- Island of activity in kinetically constrained models [5].

 $\it 1.2.\ Limiting\ objects.$  Limiting objects are quantified in terms of Sturm-Liouville problems.



**Fig. 1:** A low-temperature two-dimensional Ising model in a box of sidelength N=200 with negative boundary condition and a positive magnetic field of the form h=c/N. When c is above a critical threshold, the bulk of the system is occupied by a positively magnetized phase, while the walls are wet by a film of (unstable) negatively magnetized phase [16]. For a slightly different geometry, it was shown in [18] that this film has a width (along the walls) of order  $h^{-1/3+o(1)}$  as  $h\downarrow 0$ .

A Sturm-Liouville problem. The basic space we shall work with is  $\mathbb{L}_2(\mathbb{R}_+)$ . The notations  $\|\cdot\|_2$  and  $\langle\cdot,\cdot\rangle_2$  are reserved for the corresponding norm and scalar product. Given  $\sigma>0$  and a non-negative function  $q\in\mathsf{C}^2(\mathbb{R}_+)$  which satisfies  $\lim_{r\to\infty}q(r)=\infty$ , consider the following family of singular Sturm-Liouville operators on  $\mathbb{R}_+$ :

$$\mathsf{L} = \mathsf{L}_{\sigma,q} = \frac{\sigma^2}{2} \frac{\mathrm{d}^2}{\mathrm{d}r^2} - q(r),\tag{1.2}$$

with zero boundary condition  $\varphi(0) = 0$ .

It is a classical result [10] that L possesses a complete orthonormal family  $\{\varphi_i\}$  of simple eigenfunctions in  $\mathbb{L}_2(\mathbb{R}_+)$  with eigenvalues

$$0 > -\zeta_0 > -\zeta_1 > -\zeta_2 > \dots; \lim \zeta_i = \infty.$$
 (1.3)

The eigenfunctions  $\varphi_i$  are smooth and  $\varphi_i$  has exactly i zeroes in  $(0, \infty)$ ,  $i = 0, 1, \ldots$ 

The domain of (the closure of) L in  $\mathbb{L}_2(\mathbb{R}_+)$  is

$$\mathcal{D}(\mathsf{L}) = \left\{ f = \sum_{i} a_{i} \varphi_{i} : \sum_{i} \zeta_{i}^{2} a_{i}^{2} < \infty \right\} \quad \text{and} \quad \mathsf{L}f = -\sum_{i} \zeta_{i} a_{i} \varphi_{i} \text{ for } f \in \mathcal{D}.$$

$$(1.4)$$

Clearly,  $\mathcal{D}(\mathsf{L})$  is dense in  $\mathbb{L}_2(\mathbb{R}_+)$ . Indeed, since any function  $f \in \mathbb{L}_2$  can be written as  $f = \sum_i a_i \varphi_i$ , the linear space of all finite linear combinations  $\mathcal{U} = \left\{ \sum_{i=0}^N a_i \varphi_i \right\} \subset \mathcal{D}(\mathsf{L})$  is dense in  $\mathbb{L}_2$ . For any function  $f = \sum_i a_i \varphi_i \in \mathcal{D}(\mathsf{L})$ ,  $\lim_{N \to \infty} \mathbb{L}\left(\sum_{i=0}^N a_i \varphi_i\right) = \mathbb{L}f$ . In particular,  $\mathcal{U}$  is a core for  $\mathsf{L}$ . If  $f = \sum_i a_i \varphi_i \in \mathcal{D}$  and  $\lambda > -\zeta_0$ , one has

$$\|(\lambda I - L) f\|_2 = \left\| \sum_i (\lambda + \zeta_i) a_i \varphi_i \right\|_2 \ge (\lambda + \zeta_0) \|f\|_2,$$
 (1.5)

which shows that  $L + \zeta_0 I$  is dissipative.

Furthermore, for any  $\lambda > -\zeta_0$ , Range  $(\lambda I - L) = \mathbb{L}_2(\mathbb{R}_+)$ , and L has a compact resolvent  $R_{\lambda} = (\lambda I - L)^{-1}$ . Indeed,

$$(\lambda \mathsf{I} - \mathsf{L}) \sum_{i} \frac{a_i}{\lambda + \zeta_i} \varphi_i = \sum_{i} a_i \varphi_i \quad \text{and} \quad \mathsf{R}_{\lambda} \left( \sum_{i} a_i \varphi_i \right) = \sum_{i} \frac{a_i}{\lambda + \zeta_i} \varphi_i.$$

By the Hille-Yosida theorem,  $L + \zeta_0 I$  generates a strongly continuous contraction semigroup  $\mathsf{T}^t$  on  $\mathbb{L}_2(\mathbb{R}_+)$ . Explicitly,

$$\mathsf{T}^t \left( \sum_i a_i \varphi_i \right) = \sum_i e^{-(\zeta_i - \zeta_0)t} a_i \varphi_i. \tag{1.6}$$

Ferrari-Spohn diffusions. Define

$$\mathsf{G}_{\sigma,q}\psi = \frac{1}{\varphi_0} \left( \mathsf{L} + \zeta_0 \right) \left( \psi \varphi_0 \right) = \frac{\sigma^2}{2} \frac{\mathrm{d}^2 \psi}{\mathrm{d}r^2} + \sigma^2 \frac{\varphi_0'}{\varphi_0} \frac{\mathrm{d}\psi}{\mathrm{d}r} = \frac{\sigma^2}{2\varphi_0^2} \frac{\mathrm{d}}{\mathrm{d}r} \left( \varphi_0^2 \frac{\mathrm{d}\psi}{\mathrm{d}r} \right). \tag{1.7}$$

The sub-indices  $\sigma$  and q will be dropped whenever there is no risk of confusion. We shall say that  $\mathsf{G}_{\sigma,q}$  is the generator of a Ferrari-Spohn diffusion on  $(0,\infty)$ . The diffusion itself is ergodic and reversible with respect to the measure  $\mathrm{d}\mu_0(r) = \varphi_0^2(r)\mathrm{d}r$ . In the sequel, we shall denote by  $\mathsf{S}_{\sigma,q}^t$  the corresponding semigroup,

$$S_{\sigma,q}^t \psi = \frac{1}{\varphi_0} \mathsf{T}^t(\psi \varphi_0), \tag{1.8}$$

and by  $\mathbb{P}_{\sigma,q}$  the corresponding path measure.

1.3. Random walks with tilted areas. Let  $p_y$  be an irreducible random walk kernel on  $\mathbb{Z}$ . The probability of a finite trajectory  $\mathbb{X}=(X_1,X_2,\ldots,X_k)$  is  $\mathsf{p}(\mathbb{X})=\prod_i p_{X_{i+1}-X_i}$ . Let  $\mathsf{u},\mathsf{v}\in\mathbb{N}$  and  $M,N\in\mathbb{Z}$  with  $M\leq N$ . Let  $\mathcal{P}_{M,N,+}^{\mathsf{u},\mathsf{v}}$  be the family of trajectories starting at  $\mathsf{u}$  at time M, ending at  $\mathsf{v}$  at time N and staying positive during the time interval  $\{M,\ldots,N\}$ . For N>0, we shall use a shorthand notations  $\mathcal{P}_{N,+}^{\mathsf{u},\mathsf{v}}=\mathcal{P}_{-N,N,+}^{\mathsf{u},\mathsf{v}}$  and  $\hat{\mathcal{P}}_{N,+}^{\mathsf{u},\mathsf{v}}=\mathcal{P}_{1,N,+}^{\mathsf{u},\mathsf{v}}$ .

Assumptions on p. Assume that

$$\sum_{\mathbf{z}} \mathbf{z} p_{\mathbf{z}} = 0 = \sum_{\mathbf{z}} \mathbf{z} p_{-\mathbf{z}} \quad \text{and } \mathbf{p} \text{ has finite exponential moments.}$$
 (1.9)

In the sequel, we shall use the notation

$$\sigma^2 = \sum_{\mathbf{z}} \mathbf{z}^2 p_{\mathbf{z}} = \sum_{\mathbf{z}} \mathbf{z}^2 p_{-\mathbf{z}} = \mathbb{V}\mathrm{ar}_{\mathbf{p}}(X) < \infty. \tag{1.10}$$

The model. Let  $\{V_{\lambda}\}_{{\lambda}>0}$  be a family of self-potentials,  $V_{\lambda}:\mathbb{R}_+\to\mathbb{R}_+$ . Given  $\lambda > 0$ , define the partition function

$$Z_{N,+,\lambda}^{\mathbf{u},\mathbf{v}} = \sum_{\mathbb{X} \in \mathcal{P}_{N,+}^{\mathbf{u},\mathbf{v}}} e^{-\sum_{-N}^{N} V_{\lambda}(X_{i})} \mathsf{p}(\mathbb{X}), \tag{1.11}$$

and, accordingly, the probability distribution

$$\mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}(\mathbb{X}) = \frac{1}{Z_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}} e^{-\sum_{-N}^{N} V_{\lambda}(X_i)} \mathsf{p}(\mathbb{X}). \tag{1.12}$$

The term  $\sum_{-N}^{N} V_{\lambda}(X_i)$  represents a generalized (non-linear) area below the trajectory X. It reduces to (a multiple of) the usual area when  $V_{\lambda}(x) = \lambda x$ . We make the following set of assumptions on  $V_{\lambda}$ :

Assumptions on  $V_{\lambda}$ . For any  $\lambda > 0$ , the function  $V_{\lambda}$  is continuous monotone increasing and satisfies

$$V_{\lambda}(0) = 0$$
 and  $\lim_{x \to \infty} V_{\lambda}(x) = \infty$ . (1.13)

In particular, the relation

$$H_{\lambda}^2 V_{\lambda}(H_{\lambda}) = 1 \tag{1.14}$$

determines unambiguously the quantity  $H_{\lambda}$ . Furthermore, we make the assumptions that  $\lim_{\lambda\to 0} H_{\lambda} = \infty$  and that there exists a function  $q \in C^2(\mathbb{R}^+)$  such

$$\lim_{\lambda \to 0} H_{\lambda}^2 V_{\lambda}(rH_{\lambda}) = q(r), \tag{1.15}$$

uniformly on compact subsets of  $\mathbb{R}_+$ . Note that  $H_{\lambda}$ , respectively  $H_{\lambda}^2$ , is the spatial, respectively temporal, scale for the invariance principle which is formulated below in Theorem A.

Finally, we shall assume that there exist  $\lambda_0 > 0$  and a (continuous nondecreasing) function  $q_0 \geq 0$  with  $\lim_{r\to\infty} q_0(r) = \infty$  such that, for all  $\lambda \leq \lambda_0$ ,

$$H_{\lambda}^2 V_{\lambda}(rH_{\lambda}) \ge q_0(r) \text{ on } \mathbb{R}_+.$$
 (1.16)

1.4. Main result. It will be convenient to think about  $\mathbb X$  as being frozen outside  $\{-N,\ldots,N\}$ . In this way, we can consider  $\mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}$  as a distribution on the set of doubly infinite paths  $\mathbb{N}^{\mathbb{Z}}$ . We set  $h_{\lambda} = H_{\lambda}^{-1}$ . The paths are rescaled as follows: For  $t \in h_{\lambda}^2 \mathbb{Z}$ , define

$$x_{\lambda}(t) = h_{\lambda} X_{H_{\lambda}^{2}t} = \frac{1}{H_{\lambda}} X_{H_{\lambda}^{2}t}.$$
(1.17)

 $x_{\lambda}(t)$  is then extended to  $t\in\mathbb{R}$  by linear interpolation. In this way, we can talk about the induced distribution of  $\mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}$  on the space of continuous function C[-T, T], for any  $T \geq 0$ .

**Theorem A** Let  $\lambda_N$  be a sequence satisfying  $\lim_{N\to\infty}\lambda_N=0$ . Assume, furthermore, that  $\lim_{N\to\infty}Nh_{\lambda_N}^2=+\infty$ . Set  $x_N(\cdot)=x_{\lambda_N}(\cdot)$ . Fix any  $c\in(0,\infty)$ . Then, as  $N\to\infty$ , the distributions of  $x_N(\cdot)$  under  $\mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}$  are, uniformly in  $\mathsf{u},\mathsf{v}\leq cH_{\lambda_N}$ , weakly convergent to the stationary Ferrari-Spohn diffusion  $x_{\sigma,q}(\cdot)$  on  $\mathbb{R}_+$  with the generator  $\mathsf{G}_{\sigma,q}$  specified in (1.7).

Remark 1. The condition  $\lim_{N\to\infty} Nh_{\lambda_N}^2 = +\infty$  or, equivalently,  $H_{\lambda_N}^2 \ll N$  has a transparent meaning: N is the size of the system, whereas  $H_{\lambda_N}^2$  is the correlation length of the random walk with  $V_{\lambda_N}$ -area tilts. A precise statement of this sort is formulated in Proposition 5 in Subsection 3.4 below.

In the case  $V_{\lambda}(x) = \lambda x$ , the typical height  $H_{\lambda} = \lambda^{-\frac{1}{3}}$ , q(r) = r and the ground state  $\varphi_0$  is the rescaled Airy function:

$$\varphi_0 = \operatorname{Ai}(\chi r - \omega_1) \quad \text{and} \quad e_0 = \frac{\omega_1}{\chi},$$
(1.18)

where  $-\omega_1 = -2.33811...$  is the first zero of Ai and  $\chi = \sqrt[3]{\frac{2}{\sigma^2}}$ . Indeed, for  $\varphi_0$  defined as in (1.18),

$$\frac{\mathrm{d}^2}{\mathrm{d}r^2} \mathrm{Ai}(r) = r \mathrm{Ai}(r) \ \Rightarrow \ \frac{\mathrm{d}^2}{\mathrm{d}r^2} \varphi_0(r) = \chi^2 (\chi r - \omega_1) \varphi_0(r),$$

and one only needs to tune  $\chi$  in order to adjust to the expression (1.2) for L.

## 2. Proofs

The proof is a combination of probabilistic estimates based on an early paper [14] and rather straightforward functional analytic considerations. We shall first express the partition functions  $Z_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}$  and, subsequently, the probability distributions  $\mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}$  in terms of powers of a transfer operator  $\mathsf{T}_\lambda$ . For each  $\lambda$ , the operator  $\mathsf{T}_\lambda$  gives rise, via Doob's transform, to a stationary positive-recurrent Markov chain  $X^\lambda$  with path measure  $\mathbb{P}_\lambda$ . In the sequel, we shall refer to  $X^\lambda$  as to the ground-state chain. Following (1.17), the ground-state chains are rescaled as  $x_\lambda(t) = h_\lambda X_{H_\lambda^2 t}^\lambda$ .

The proof of Theorem A comprises three steps:

STEP 1. As  $\lambda \to 0$ , the finite-dimensional distributions of the rescaled ground-state chains  $x_{\lambda}$  converge to the finite-dimensional distributions of the Ferrari-Spohn diffusion  $x_{\sigma,q}$ . This is the content of Corollary 1 in Subsection 2.2.

STEP 2. Under our assumptions on p and on the family of potentials  $V_{\lambda}$ , the induced family of distributions  $\mathbb{P}_{\lambda}$  is tight on  $\mathsf{C}[-T,T]$  for any  $T<\infty$ . This is the content of Proposition 4 in Subsection 3.3.

STEP 3. Under the conditions of Theorem A, the following happens: For each  $T \geq 0$ , the variational distance between the induced distributions on  $\mathsf{C}[-T,T]$  of  $\mathbb{P}^{\mathsf{u},\mathsf{v}}_{N,+,\lambda_N}$  and of  $\mathbb{P}_{\lambda_N}$  tends to zero as  $N \to \infty$ . This is the content of Corollary 2 in Subsection 3.4.

Remark on constants.  $c_1, \nu_1, \kappa_1, c_2, \nu_2, \kappa_2, \ldots$  denote positive constants which may take different values in different Subsections, but are otherwise universal, in the sense that the corresponding bounds hold uniformly in the range of the relevant parameters.

2.1. The operator  $\mathsf{T}_{\lambda}$  and its Doob transform. In the sequel, we shall make a slight abuse of notation and identify the spaces  $\ell_p(\mathbb{N})$  with sub-spaces of  $\ell_p(\mathbb{Z})$ :

$$\ell_p(\mathbb{N}) = \{ \phi \in \ell_p(\mathbb{Z}) : \phi(\mathsf{x}) = 0 \text{ for } \mathsf{x} \le 0 \}.$$

For  $\lambda > 0$ , consider the operators  $\tilde{\mathsf{T}}_{\lambda}$  defined by

$$\tilde{\mathsf{T}}_{\lambda}[\phi](\mathsf{x}) = \sum_{\mathsf{y}} p_{\mathsf{y}-\mathsf{x}} e^{-\frac{1}{2}(V_{\lambda}(\mathsf{x}) + V_{\lambda}(\mathsf{y}))} \phi(\mathsf{y}). \tag{2.1}$$

In terms of  $\tilde{\mathsf{T}}_{\lambda}$ , the partition function can be expressed as

$$e^{\frac{1}{2}(V_{\lambda}(\mathsf{u})+V_{\lambda}(\mathsf{v}))}Z_{N,+,\lambda}^{\mathsf{u},\mathsf{v}} = \tilde{\mathsf{T}}_{\lambda}^{2N}[\mathbf{1}_{\mathsf{v}}](\mathsf{u}). \tag{2.2}$$

For each  $\lambda > 0$ , the operator  $\tilde{\mathsf{T}}_{\lambda}$  is positive on  $\ell_p(\mathbb{N})$  and compact from  $\ell_p(\mathbb{N})$  to  $\ell_q(\mathbb{N})$  for every  $p,q \in [1,\infty]$  (we use  $\ell_\infty(\mathbb{N})$  for the Banach space of functions  $\phi \in \mathbb{R}^{\mathbb{N}}$  which tend to zero as  $\mathsf{x} \to \infty$ ). Indeed,  $\{\mathbf{1}_{\mathsf{x}}\}_{\mathsf{x} \in \mathbb{N}}$  is a basis of  $\ell_p(\mathbb{N})$ . Since, for  $\mathsf{x}, \mathsf{y} > 0$ ,

$$\tilde{\mathsf{T}}_{\lambda}[\mathbf{1}_{\mathsf{x}}](\mathsf{y}) = \mathrm{e}^{-\frac{1}{2}(V_{\lambda}(\mathsf{x}) + V_{\lambda}(\mathsf{y}))} p_{\mathsf{x} - \mathsf{y}},$$

it follows that

$$\tilde{\mathsf{T}}_{\lambda}[\mathbf{1}_{\mathsf{x}}] = \mathrm{e}^{-\frac{1}{2}V_{\lambda}(\mathsf{x})} \sum_{\mathsf{y}} p_{\mathsf{x}-\mathsf{y}} \mathrm{e}^{-\frac{1}{2}V_{\lambda}(\mathsf{y})} \mathbf{1}_{\mathsf{y}} \ \Rightarrow \ \|\tilde{\mathsf{T}}_{\lambda}[\mathbf{1}_{\mathsf{x}}]\|_{p} \le \mathrm{e}^{-\frac{1}{2}V_{\lambda}(\mathsf{x})} \sum_{\mathsf{y}} p_{\mathsf{x}-\mathsf{y}} \mathrm{e}^{-\frac{1}{2}V_{\lambda}(\mathsf{y})}.$$

Hence, by the second condition in (1.13) and by the assumption on exponential tails of p in (1.9), the closure  $\{\tilde{T}_{\lambda}[1_x]\}$  is compact in any  $\ell_q(\mathbb{N})$  whenever  $\lambda > 0$ .

Since  $\tilde{\mathsf{T}}_\lambda$  is a positive operator (and since, e.g.,  $\sum_n 2^{-n} \tilde{\mathsf{T}}_\lambda^n$  is strictly positive and still compact), the Krein-Rutman Theorem [15, Theorem 6.3] applies, and  $\tilde{\mathsf{T}}_\lambda$  possesses a strictly positive leading eigenfunction  $\phi_\lambda$  (the same for all  $\ell_p(\mathbb{N})$  spaces, by compact embedding) of algebraic multiplicity one. Let  $E_\lambda$  be the corresponding leading eigenvalue. All the above reasoning applies to the adjoint operator  $\tilde{\mathsf{T}}_\lambda^*$  with matrix entries

$$\tilde{\mathsf{T}}_{\lambda}^{*}(\mathsf{x},\mathsf{y}) = \tilde{\mathsf{T}}_{\lambda}(\mathsf{y},\mathsf{x}) = \mathrm{e}^{-\frac{1}{2}(V_{\lambda}(\mathsf{x}) + V_{\lambda}(\mathsf{y}))} p_{\mathsf{x}-\mathsf{y}}. \tag{2.4}$$

Let  $\phi_{\lambda}^*$  be the Krein-Rutman eigenfunction (with the very same leading eigenvalue  $E_{\lambda}$ ) of  $\tilde{\mathsf{T}}_{\lambda}^*$ . As Theorem A indicates, the relevant spatial scale is given by  $h_{\lambda} = H_{\lambda}^{-1}$ . To fix notation, we normalize  $\phi_{\lambda}$  and  $\phi_{\lambda}^*$  as in (2.19) below, that is  $h_{\lambda} \sum_{x} \phi_{\lambda}(x)^2 = h_{\lambda} \sum_{x} (\phi_{\lambda}^*(x))^2 = 1$ .

It will be convenient to work with the following normalized version  $T_{\lambda}$  of  $\tilde{T}_{\lambda}$ : for  $x, y \in \mathbb{N}$ , let us introduce the kernel

$$\mathsf{T}_{\lambda}(\mathsf{x},\mathsf{y}) = \frac{1}{E_{\lambda}} \tilde{\mathsf{T}}_{\lambda}(\mathsf{x},\mathsf{y}) = \frac{1}{E_{\lambda}} \mathrm{e}^{-\frac{1}{2}(V_{\lambda}(\mathsf{x}) + V_{\lambda}(\mathsf{y}))} p_{\mathsf{y}-\mathsf{x}}. \tag{2.5}$$

In this way,  $\phi_{\lambda}$  and  $\phi_{\lambda}^{*}$  are the principal positive Krein-Rutman eigenfunctions of  $T_{\lambda}$  and  $T_{\lambda}^{*}$  with eigenvalue 1.

Ground-state chains. Define

$$\pi_{\lambda}(\mathsf{x},\mathsf{y}) = \frac{1}{\phi_{\lambda}(\mathsf{x})} \mathsf{T}_{\lambda}(\mathsf{x},\mathsf{y}) \phi_{\lambda}(\mathsf{y}) \quad \text{and} \quad \pi_{\lambda}^{*}(\mathsf{x},\mathsf{y}) = \frac{1}{\phi_{\lambda}^{*}(\mathsf{x})} \mathsf{T}_{\lambda}^{*}(\mathsf{x},\mathsf{y}) \phi_{\lambda}^{*}(\mathsf{y}). \tag{2.6}$$

 $\pi_{\lambda}$  and  $\pi_{\lambda}^{*}$  are irreducible Markov kernels on  $\mathbb{N}$ . The corresponding chains are positively recurrent and have the invariant probability measure  $\mu_{\lambda}(\mathsf{x}) = c_{\lambda}\phi_{\lambda}^{*}(\mathsf{x})\phi_{\lambda}(\mathsf{x})$ . As we shall prove below in Theorem 2,  $\lim_{\lambda \to 0} \frac{h_{\lambda}}{c_{\lambda}} = 1$ .

Notice that

$$\mu_{\lambda}(\mathsf{x})\pi_{\lambda}(\mathsf{x},\mathsf{y}) = \mu_{\lambda}(\mathsf{y})\pi_{\lambda}^{*}(\mathsf{y},\mathsf{x}) = c_{\lambda}\phi_{\lambda}^{*}(\mathsf{x})\mathsf{T}_{\lambda}(\mathsf{x},\mathsf{y})\phi_{\lambda}(\mathsf{y}). \tag{2.7}$$

In the sequel, we shall denote by  $\mathbb{P}_{\lambda}$  the stationary distribution on  $\mathbb{N}^{\mathbb{Z}}$  of the (direct) ground-state chain which corresponds to  $\pi_{\lambda}$ .

Variational description of  $E_{\lambda}$  and  $\mu_{\lambda}$ . Let us formulate a Donsker-Varadhan type formula for the principal eigenvalue  $E_{\lambda}$  of  $\tilde{\mathsf{T}}_{\lambda}$  or, equivalently, for the eigenvalue 1 of  $\mathsf{T}_{\lambda}$ .

#### Theorem 1.

$$1 = \sup_{\mu} \inf_{u \in \mathbb{U}_+} \sum_{\mathsf{x} \in \mathbb{N}} \mu(\mathsf{x}) \frac{\mathsf{T}_{\lambda}[u]}{u}(\mathsf{x}). \tag{2.8}$$

Above, the first supremum is over probability measures on N, and

$$\mathbb{U}_{+} = \left\{ u = v\phi_{\lambda} : 0 < \inf v \le \sup v < \infty \right\}. \tag{2.9}$$

Remark 2. Eventually, our proof of Theorem A will not rely on the variational formula (2.8). Theorem 1 and its consequence, Proposition 1, are formulated and proved in their own right, but also because they elucidate the type of variational convergence behind Theorem 2 below.

*Proof.* As before, set  $h_{\lambda} = H_{\lambda}^{-1}$  and consider the following functional:

$$\mathcal{F}_{\lambda}(\mu) = \frac{1}{h_{\lambda}^2} \sup_{u \in \mathbb{U}_+} \sum_{\mathbf{x} \in \mathbb{N}} \mu(\mathbf{x}) \frac{(1 - \mathsf{T}_{\lambda})u}{u}(\mathbf{x}). \tag{2.10}$$

The coefficient  $h_{\lambda}^2$  plays no role in the proof, it just fixes the proper scaling. The claim of Theorem 1 will follow once we show that

$$\mathcal{F}_{\lambda}(\mu_{\lambda}) = 0$$
 and  $\mathcal{F}_{\lambda}(\mu) > 0$  whenever  $\mu \neq \mu_{\lambda}$ . (2.11)

Taking  $u \equiv \phi_{\lambda}$ , we readily infer that  $\mathcal{F}_{\lambda}$  is non-negative. In order to check the first statement in (2.11), we need to verify that

$$\sum_{\mathbf{x} \in \mathbb{N}} \mu_{\lambda}(\mathbf{x}) \frac{(1 - \mathsf{T}_{\lambda})u}{u}(\mathbf{x}) \le 0, \tag{2.12}$$

whenever  $u = v\phi_{\lambda}$  and  $v \not\equiv 1$ . Let us write v as  $v = e^{h}$ . Then, using the notation (2.6),

$$\sum_{\mathsf{x}\in\mathbb{N}}\mu_{\lambda}(\mathsf{x})\frac{(1-\mathsf{T}_{\lambda})u}{u}(\mathsf{x}) = \sum_{\mathsf{x}\in\mathbb{N}} \left(1-\mathrm{e}^{-h}\pi_{\lambda}\mathrm{e}^{h}\right)\mu_{\lambda}(\mathsf{x}) \leq \left(1-\mathrm{e}^{-\langle\mu_{\lambda},h\rangle+\langle\mu_{\lambda},\pi_{\lambda}h\rangle}\right) = 0,$$

where we used again Jensen's inequality and the invariance of  $\mu_{\lambda}$ :  $\mu_{\lambda}\pi_{\lambda}=\mu_{\lambda}$ . If  $\mu=g^2\mu_{\lambda}$  and g is bounded away from 0 and  $\infty$ , then, taking  $u=g\phi_{\lambda}$ ,

$$\sum_{\mathbf{x} \in \mathbb{N}} \mu(\mathbf{x}) \frac{(1 - \mathsf{T}_{\lambda})u}{u}(\mathbf{x}) = \frac{1}{2} \sum_{\mathbf{x}, \mathbf{y}} \hat{\pi}_{\lambda}(\mathbf{x}, \mathbf{y}) (g(\mathbf{x}) - g(\mathbf{y}))^{2} \mu_{\lambda}(\mathbf{x}), \tag{2.13}$$

where  $\hat{\pi}_{\lambda}$  is the symmetrized kernel,

$$\hat{\pi}_{\lambda}(x,y) = \frac{1}{2}(\pi_{\lambda}(x,y) + \pi_{\lambda}^{*}(x,y)). \tag{2.14}$$

By (2.7),

$$\mathcal{F}_{\lambda}(g^{2}\mu_{\lambda}) \geq \frac{c_{\lambda}}{4} \sum_{\mathbf{x},\mathbf{y}} \left(\frac{g(\mathbf{x}) - g(\mathbf{y})}{h_{\lambda}}\right)^{2} (\phi_{\lambda}^{*}(\mathbf{x})\phi_{\lambda}(\mathbf{y}) + \phi_{\lambda}^{*}(\mathbf{y})\phi_{\lambda}(\mathbf{x})). \tag{2.15}$$

We claim that (2.15) still holds when g is not bounded away from 0 and  $\infty$ . This will follow if we show that there exists a sequence  $u_{\ell} \in \mathbb{U}_+$  such that

$$\limsup_{\ell \to \infty} \sum_{\mathsf{x} \in \mathbb{N}} \mu_{\lambda}(\mathsf{x}) g^{2}(\mathsf{x}) \frac{\mathsf{T}_{\lambda}[u_{\ell}\phi_{\lambda}](\mathsf{x})}{u_{\ell}(\mathsf{x})\phi_{\lambda}(\mathsf{x})} \leq \sum_{\mathsf{x} \in \mathbb{N}} \mu_{\lambda}(\mathsf{x}) g^{2}(\mathsf{x}) \frac{\mathsf{T}_{\lambda}[g\phi_{\lambda}](\mathsf{x})}{g(\mathsf{x})\phi_{\lambda}(\mathsf{x})}$$
$$= c_{\lambda} \sum_{\mathsf{x} \in \mathbb{N}} \phi_{\lambda}^{*}(\mathsf{x}) g(\mathsf{x}) \mathsf{T}_{\lambda}[g\phi_{\lambda}](\mathsf{x}). \tag{2.16}$$

Assume that g is not bounded away from zero and consider  $g_n = g \vee \frac{1}{n}$ . Then,

$$\sum_{\mathbf{x} \in \mathbb{N}} \mu_{\lambda}(\mathbf{x}) g^2(\mathbf{x}) \frac{\mathsf{T}_{\lambda}[g_n \phi_{\lambda}](\mathbf{x})}{g_n(\mathbf{x}) \phi_{\lambda}(\mathbf{x})} \leq c_{\lambda} \sum_{\mathbf{x} \in \mathbb{N}} \phi^*(\mathbf{x}) g(\mathbf{x}) \mathsf{T}_{\lambda}[g_n \phi_{\lambda}](\mathbf{x}).$$

By a monotone convergence argument, the right-hand side above converges (as  $n \to \infty$ ) to  $c_{\lambda} \sum_{\mathbf{x} \in \mathbb{N}} \phi^*(\mathbf{x}) g(\mathbf{x}) \mathsf{T}_{\lambda}[g\phi_{\lambda}](\mathbf{x})$ . If, in addition, g is not bounded above, then consider  $g_{n,M} = g_n \land M \in \mathbb{U}_+$ . Define  $A_M = \{\mathbf{x} : g(\mathbf{x}) > M\}$ . Then,

$$\sum_{\mathbf{x} \in \mathbb{N}} \mu_{\lambda}(\mathbf{x}) g^2(\mathbf{x}) \frac{\mathsf{T}_{\lambda}[g_{n,M}\phi_{\lambda}](\mathbf{x})}{g_{n,M}(\mathbf{x})\phi_{\lambda}(\mathbf{x})} \leq \sum_{\mathbf{x} \not \in A_M} \mu_{\lambda}(\mathbf{x}) g^2(\mathbf{x}) \frac{\mathsf{T}_{\lambda}[g_n\phi_{\lambda}](\mathbf{x})}{g_n(\mathbf{x})\phi_{\lambda}(\mathbf{x})} + \sum_{\mathbf{x} \in A_M} \mu_{\lambda}(\mathbf{x}) g^2(\mathbf{x}).$$

Since  $\lim_{M\to\infty} \sum_{\mathsf{x}\in A_M} \mu_{\lambda}(\mathsf{x}) g^2(\mathsf{x}) = 0$ , the approximation procedure goes through as claimed in (2.16).  $\square$ 

As a byproduct, we obtain the following

**Proposition 1.** The functional  $\mathcal{F}_{\lambda}$  is convex and lower-semicontinuous. It has a unique minimum:

$$\mathcal{F}_{\lambda}(\mu) = 0 \iff \mu = \mu_{\lambda}. \tag{2.17}$$

Furthermore,  $\mu_{\lambda}$  is a quadratic minimum in the sense of (2.15).

2.2. Convergence of eigenfunctions, invariant measures and semigroups. It will be convenient to think about  $\mathsf{T}_{\lambda}$  and  $\pi_{\lambda}$  as acting on the rescaled spaces  $\ell_2(\mathbb{N}_{\lambda})$ , where

$$\mathbb{N}_{\lambda} = h_{\lambda} \mathbb{N}$$
 and the scalar product is  $\langle u, v \rangle_{2,\lambda} = h_{\lambda} \sum_{\mathbf{r} \in \mathbb{N}_{\lambda}} u(\mathbf{r}) v(\mathbf{r})$ . (2.18)

Accordingly, we rescale  $\phi_{\lambda}$  and  $\phi_{\lambda}^{*}$  in such a way that

$$\|\phi_{\lambda}\|_{2,\lambda} = \|\phi_{\lambda}^*\|_{2,\lambda} = 1. \tag{2.19}$$

We use the same notation  $\mu_{\lambda} = c_{\lambda} \phi_{\lambda} \phi_{\lambda}^{*}$  for the rescaled probability measure on  $\mathbb{N}_{\lambda}$ . In other words, the constants  $c_{\lambda}$  are defined via

$$\frac{1}{c_{\lambda}} = \sum_{\mathbf{r} \in \mathbb{N}_{\lambda}} \phi_{\lambda}(\mathbf{r}) \phi_{\lambda}^{*}(\mathbf{r}) \quad \text{or} \quad \frac{h_{\lambda}}{c_{\lambda}} = \langle \phi_{\lambda}, \phi_{\lambda}^{*} \rangle_{2,\lambda}, \tag{2.20}$$

where  $\phi_{\lambda}$  and  $\phi_{\lambda}^{*}$  are the principal eigenfunctions satisfying the normalization condition (2.19).

Remark 3. As in the case of  $\ell_p(\mathbb{N})$ , with a slight abuse of notation, we shall identify  $\ell_2(\mathbb{N}_{\lambda})$  with a closed linear sub-space of  $\ell_2(\mathbb{Z}_{\lambda})$ , where  $\mathbb{Z}_{\lambda} = h_{\lambda}\mathbb{Z}$ . Namely,

$$\ell_2(\mathbb{N}_{\lambda}) = \{ u \in \ell_2(\mathbb{Z}_{\lambda}) : u(r) = 0 \text{ for all } r \le 0 \}.$$
 (2.21)

In this way, if  $k_{\lambda}$  is a kernel on  $\mathbb{Z}_{\lambda}$ , then the operators

$$u(\cdot) \mapsto \mathbf{1}_{\cdot \in \mathbb{N}_{\lambda}} \sum_{\mathsf{s} \in \mathbb{Z}_{\lambda}} k_{\lambda}(\mathsf{s} - \cdot) u(\mathsf{s}) \quad \text{and} \quad u(\cdot) \mapsto \mathbf{1}_{\cdot \in \mathbb{N}_{\lambda}} \sum_{\mathsf{s} \in \mathbb{Z}_{\lambda}} k_{\lambda}(\cdot - \mathsf{s}) (u(\mathsf{s}) - u(\cdot)),$$

can be considered as operators on  $\ell_2(\mathbb{N}_{\lambda})$ . Accordingly,

$$\sum_{\mathsf{s},\mathsf{r}\in\mathbb{Z}_{\lambda}}k_{\lambda}(\mathsf{s}-\mathsf{r})(u(\mathsf{s})-u(\mathsf{r}))u(\mathsf{r})$$

is a quadratic form on  $\ell_2(\mathbb{N}_{\lambda})$ .

In the sequel, we shall write  $p_{\lambda}(\mathsf{r}) = p_{H_{\lambda}\mathsf{r}}$  for the rescaled random walk kernel on  $\mathbb{Z}_{\lambda}$ .

Convergence of Hilbert spaces. Let us fix a map  $\rho_{\lambda}: \mathbb{L}_2(\mathbb{R}_+) \to \ell_2(\mathbb{N}_{\lambda})$  with  $\|\rho_{\lambda}\| \leq 1$ . The specific choice is not really important; for instance, we may define

$$\rho_{\lambda} u(\mathbf{r}) = \frac{1}{h_{\lambda}} \int_{\mathbf{r} - h_{\lambda}}^{\mathbf{r}} u(s) ds.$$
 (2.22)

**Definition 1.** Let us say that a sequence  $u_{\lambda} \in \ell_2(\mathbb{N}_{\lambda})$  converges to  $u \in \mathbb{L}_2(\mathbb{R}_+)$ ,  $u = \text{s-}\lim u_{\lambda}$ , if

$$\lim_{\lambda \to 0} \|u_{\lambda} - \rho_{\lambda} u\|_{2,\lambda} = 0. \tag{2.23}$$

We shall write  $\lim_{\lambda\to 0}$  instead of s- $\lim_{\lambda\to 0}$  whenever no ambiguity arises.

Compactness of eigenfunctions. The following two probabilistic estimates will be proved in Section 3.

**Lemma 1.** Define  $e_{\lambda} = -H_{\lambda}^2 \log E_{\lambda}$ . Then,

$$0 < \liminf_{\lambda \to 0} \mathsf{e}_{\lambda} \le \limsup_{\lambda \to 0} \mathsf{e}_{\lambda} < \infty. \tag{2.24}$$

As we already noted, it follows from the compactness of  $\mathsf{T}_{\lambda}$  that the eigenfunctions  $\phi_{\lambda}$  and  $\phi_{\lambda}^{*}$  belong to  $\ell_{p}(\mathbb{N}_{\lambda})$  for any  $p \geq 1$  and  $\lambda > 0$ . The second probabilistic input is a tail estimate on  $\phi_{\lambda}$  and  $\phi_{\lambda}^{*}$ .

**Lemma 2.** There exist positive constants  $\nu_1$  and  $\nu_2$  such that

$$h_{\lambda} \sum_{\mathsf{r} \in \mathbb{N}_{\lambda}} \phi_{\lambda}(\mathsf{r}) \mathbf{1}_{\{\mathsf{r} > K\}} \le \nu_{1} e^{-\nu_{2} K H_{\lambda}(\sqrt{V_{\lambda}(H_{\lambda}K)} \wedge 1)} \le \nu_{1} e^{-\nu_{2} K (\sqrt{q_{0}(K)} \wedge H_{\lambda})}, \quad (2.25)$$

uniformly in K > 0 and  $\lambda \leq \lambda_0$ . The same holds for  $\phi_{\lambda}^*$ . In particular, both sequences  $\phi_{\lambda}$  and  $\phi_{\lambda}^*$  are bounded in  $\ell_1(\mathbb{N}_{\lambda})$ :

$$\limsup_{\lambda \to 0} h_{\lambda} \sum_{\mathbf{r} \in \mathbb{N}_{\lambda}} \phi_{\lambda}(\mathbf{r}) < \infty \quad and \quad \limsup_{\lambda \to 0} h_{\lambda} \sum_{\mathbf{r} \in \mathbb{N}_{\lambda}} \phi_{\lambda}^{*}(\mathbf{r}) < \infty. \tag{2.26}$$

Using the two lemmas above and Rellich's theorem (see, e.g., [2, Chapter 6]) on compact embeddings of the Sobolev spaces  $\mathbb{H}^1[a,b]$  into  $\mathbb{L}_2[a,b]$  for finite intervals [a,b], we shall prove the following

**Proposition 2.** Under our assumptions on  $V_{\lambda}$  and p, the sequence  $\phi_{\lambda}$  is sequentially compact (in the sense of s-convergence as described above in (2.23)).

*Proof.* The proof comprises two steps: We first show that we can restrict attention to the compactness properties of the functions  $\psi_{\lambda}$  defined in (2.27) below. We then check that the sequence  $\psi_{\lambda}$  satisfies the energy-type estimate (2.33), which enables a uniform control of both tails of  $\psi_{\lambda}$  and of their Sobolev norms over  $\mathbb{R}_+$ . In this way, sequential compactness follows by a standard diagonal argument.

STEP 1. In view of Lemma 2, rather than studying directly the functions  $\phi_{\lambda}$ , we can instead study the convergence properties of the functions

$$\psi_{\lambda}(\mathbf{r}) = e^{-\frac{1}{2}V_{\lambda}(H_{\lambda}\mathbf{r})}\phi_{\lambda}(\mathbf{r}). \tag{2.27}$$

Indeed, by (2.25), there exists a sequence  $\delta_{\lambda} \to 0$  such that

$$\lim_{\lambda \to 0} h_{\lambda} \sum_{\mathbf{r} \in \mathbb{N}_{\lambda}} \phi_{\lambda}^{2}(\mathbf{r}) \mathbf{1}_{\{V_{\lambda}(H_{\lambda}\mathbf{r}) > \delta_{\lambda}\}} = 0.$$
 (2.28)

So, (2.28) implies that the norm of the difference  $\|\phi_{\lambda} - \psi_{\lambda}\|_{2,\lambda}$  tends to zero, and hence  $\phi_{\lambda} - \psi_{\lambda}$  tends to zero in the sense of Definition 1.

STEP 2. In terms of  $\psi_{\lambda}$ , the eigenvalue equation  $\tilde{\mathsf{T}}_{\lambda}\phi_{\lambda}=E_{\lambda}\phi_{\lambda}$  reads as (recall Remark 3)

$$\sum_{\mathbf{s} \in \mathbb{Z}_{\lambda}} p_{\lambda}(\mathbf{s} - \mathbf{r})(\psi_{\lambda}(\mathbf{s}) - \psi_{\lambda}(\mathbf{r})) = (E_{\lambda} e^{V_{\lambda}(H_{\lambda}\mathbf{r})} - 1)\psi_{\lambda}(\mathbf{r}). \tag{2.29}$$

Multiplying both sides by  $-\psi_{\lambda}(\mathbf{r})$  and summing over  $\mathbf{r}$ , we get

$$h_{\lambda} \sum_{\mathbf{r}, \mathbf{s} \in \mathbb{Z}_{\lambda}} p_{\lambda}(\mathbf{s} - \mathbf{r}) \left( \frac{-\psi_{\lambda}(\mathbf{r})\psi_{\lambda}(\mathbf{s}) + \psi_{\lambda}^{2}(\mathbf{r})}{h_{\lambda}^{2}} \right) + h_{\lambda} \sum_{\mathbf{r} \in \mathbb{N}_{\lambda}} \frac{E_{\lambda} e^{V_{\lambda}(H_{\lambda}\mathbf{r})} - 1}{h_{\lambda}^{2}} \psi_{\lambda}^{2}(\mathbf{r}) = 0.$$

So, for the symmetrized kernel  $\hat{p}_{\lambda}(\mathsf{z}) = (p_{\lambda}(\mathsf{z}) + p_{\lambda}(-\mathsf{z}))/2$ , we obtain

$$h_{\lambda} \sum_{\mathbf{r}, \mathbf{s} \in \mathbb{Z}_{\lambda}} \hat{p}_{\lambda}(\mathbf{s} - \mathbf{r}) \left( \frac{\psi_{\lambda}(\mathbf{s}) - \psi_{\lambda}(\mathbf{r})}{h_{\lambda}} \right)^{2} + h_{\lambda} \sum_{\mathbf{r} \in \mathbb{N}_{\lambda}} \frac{E_{\lambda} e^{V_{\lambda}(H_{\lambda}\mathbf{r})} - 1}{h_{\lambda}^{2}} \psi_{\lambda}^{2}(\mathbf{r}) = 0. \quad (2.30)$$

In view of Lemma 1, we may assume that there exists  $\bar{\mathsf{e}} < \infty$  such that, possibly going to a subsequence, the limit  $\mathsf{e} = \lim_{\lambda \to 0} \mathsf{e}_{\lambda}$  exists and satisfies  $\mathsf{e} < \bar{\mathsf{e}}$ . So, we may assume that  $E_{\lambda} \geq \mathrm{e}^{-\bar{\mathsf{e}}h_{\lambda}^2}$ . Recall also our assumption (1.16) on the growth of  $V_{\lambda}$ . Let  $\bar{\mathsf{r}} = \sup \{\mathsf{r} : q_0(\mathsf{r}) < \bar{\mathsf{e}}\}$ . Then, (2.30) implies that

$$h_{\lambda} \sum_{\mathsf{r},\mathsf{s} \in \mathbb{Z}_{\lambda}} \hat{p}_{\lambda}(\mathsf{s} - \mathsf{r}) \Big( \frac{\psi_{\lambda}(\mathsf{s}) - \psi_{\lambda}(\mathsf{r})}{h_{\lambda}} \Big)^{2} + h_{\lambda} \sum_{\mathsf{r} \geq \bar{\mathsf{r}}} q_{0}(\mathsf{r}) \psi_{\lambda}^{2}(\mathsf{r}) \leq \bar{\mathsf{e}} \|\psi_{\lambda}\|_{2,\lambda}^{2}. \tag{2.31}$$

By construction,  $\|\psi_{\lambda}\|_{2,\lambda}^2 \leq 1$  and, as we have already mentioned, (2.28) implies that actually  $\lim_{\lambda \to 0} \|\psi_{\lambda}\|_{2,\lambda}^2 = 1$ .

Furthermore, since p is an irreducible kernel, there exists  $\delta > 0$  and a finite sequence of integer states  $x_0, x_1, \ldots, x_n$  with  $\hat{p}_{x_i - x_{i-1}} \geq \delta$ , which connects  $x_0 = 0$  to  $x_n = 1$ . Therefore,

$$\sum_{\mathbf{r},\mathbf{s}\in\mathbb{Z}_{\lambda}}\hat{p}_{\lambda}(s-r)\Big(\frac{\psi_{\lambda}(\mathbf{s})-\psi_{\lambda}(\mathbf{r})}{h_{\lambda}}\Big)^{2}\geq\frac{\delta}{n^{2}}\sum_{\mathbf{r}\in\mathbb{N}_{\lambda}}\Big(\frac{\psi_{\lambda}(\mathbf{r})-\psi_{\lambda}(\mathbf{r}-h_{\lambda})}{h_{\lambda}}\Big)^{2},\qquad(2.32)$$

where we use the elementary inequality

$$(z_0-z_1)^2+(z_1-z_2)^2+\ldots+(z_{n-1}-z_n)^2\geq \frac{1}{n}(z_0-z_n)^2,$$

valid for all real  $z_i$ . The additional 1/n in the prefactor  $1/n^2$  in (2.32) is due to the fact that each term  $(\psi_{\lambda}(s) - \psi_{\lambda}(r))^2$  is used in this way at most n times. Together with (2.31), this implies the existence of two finite positive constants  $c_1$  and  $c_2$  such that

$$c_1 h_{\lambda} \sum_{\mathbf{r} \in \mathbb{N}_{\lambda}} \left( \frac{\psi_{\lambda}(\mathbf{r}) - \psi_{\lambda}(\mathbf{r} - h_{\lambda})}{h_{\lambda}} \right)^2 + h_{\lambda} \sum_{\mathbf{r} > \bar{\mathbf{r}}} q_0(\mathbf{r}) \psi_{\lambda}^2(\mathbf{r}) \le c_2. \tag{2.33}$$

This is the desired energy estimate, which holds for all  $\lambda > 0$  small.

The rest of the proof is straightforward. Let  $\Psi_{\lambda}$  be the linear interpolation of  $\psi_{\lambda}$ : for  $r \in \mathbb{N}_{\lambda} \cup \{0\}$  and  $t \in [0, 1]$ ,

$$\Psi_{\lambda}(\mathbf{r} + th_{\lambda}) = (1 - t)\psi_{\lambda}(\mathbf{r}) + t\psi_{\lambda}(\mathbf{r} + h_{\lambda}).$$

The relation (2.33) and  $\lim_{r\to\infty} q_0(r) = \infty$  imply that  $\lim_{n\to\infty} \|\Psi_{\lambda} \mathbf{1}_{\{r>n\}}\|_2 = 0$ , uniformly in  $\lambda$  small. On the other hand, the very same (2.33) and Rellich's compact embedding theorem imply that, for any  $n < \infty$ , the family  $\Psi_{\lambda} \mathbf{1}_{\{r \le n\}}$  is subsequentially compact in  $\mathbb{L}_2[0, n]$ . Alternatively, (2.33) implies that the linear

interpolations  $\Psi_{\lambda}$  are uniformly continuous on [0, n] for each n fixed. We conclude that the family  $\Psi_{\lambda}$  is subsequentially compact in  $\mathbb{L}_2(\mathbb{R}_+)$ .

Remember how the map  $\rho_{\lambda}$  was defined in (2.22). Since  $\Psi_{\lambda}$  is the linear interpolation of  $\psi_{\lambda}$ , and since  $\lim_{\lambda \to 0} h_{\lambda} = 0$ , the energy estimate (2.33) evidently implies that

$$\lim_{\lambda \to 0} \|\psi_{\lambda} - \rho_{\lambda} \Psi_{\lambda}\|_{2,\lambda} = 0.$$

Hence,  $\psi_{\lambda}$  is subsequentially compact as well.  $\Box$ 

Convergence of semigroups. Possibly going to a subsequence, we can assume that  $\lim e_{\lambda} = e$ . We shall rely on Kurtz's semigroup convergence theorem [11, Theorem I.6.5]: Define

$$\mathsf{L}_{\lambda}f(\mathsf{r}) = \frac{\mathsf{T}_{\lambda} - \mathsf{I}}{h_{\lambda}^{2}}f(\mathsf{r}). \tag{2.34}$$

The following two statements are equivalent:

- (a) For any  $u \in \mathcal{U}$ , one can find a sequence  $u_{\lambda} \in \ell_2(\mathbb{N}_{\lambda})$  such that both  $\lim_{\lambda \to 0} u_{\lambda} = u$  and  $\lim_{\lambda \to 0} \mathsf{L}_{\lambda} u_{\lambda} = (\mathsf{L} + \mathsf{e})u$ .
- (b) If  $\lim_{\lambda \to 0} f_{\lambda} = f$ , then  $\lim_{\lambda \to 0} \mathsf{T}_{\lambda}^{\lfloor H_{\lambda}^{2} t \rfloor} f_{\lambda} = e^{(\mathsf{L} + \mathsf{el})t} f$ .

The above equivalence holds provided that the operators  $\mathsf{T}_\lambda$  are linear contractions (which is straightforward), and that  $\mathrm{e}^{(\mathsf{L}+\mathsf{el})t}$  is a strongly continuous semigroup with generator  $\mathsf{L}+\mathsf{el}$ , but that's exactly how it was constructed, see (1.6). Recall that the core  $\mathcal U$  consists of finite linear combinations of eigenfunctions  $\varphi_j$ . Equivalently, we might have considered  $\mathcal U'=\mathsf{C}_0^2[0,\infty)$ . Indeed, if  $\chi_0$  is a smooth function which is 1 on  $(-\infty,0]$  and 0 on  $[1,\infty)$  and if  $\chi_R(r)=\chi_0(r-R)$ , then, for any j,

$$\lim_{R \to \infty} \chi_R \varphi_j = \varphi_j \quad \text{and} \quad \lim_{R \to \infty} \mathsf{L}_{\sigma, q}(\chi_R \varphi_j) = \mathsf{L}_{\sigma, q} \varphi_j = -\zeta_j \varphi_j. \tag{2.35}$$

Above, both convergences are pointwise and in  $\mathbb{L}_2(\mathbb{R}_+)$ . In order to check the second claim in (2.35), just note that

$$\mathsf{L}_{\sigma,q}(\chi_R \varphi_j) = -\zeta_j \varphi_j + \frac{\sigma^2}{2} (\varphi_j \chi_R'' + 2\varphi_j' \chi_R'),$$

and the conclusion follows, since both  $\varphi_j$  and  $\varphi_i'$  belong to  $\mathbb{L}^2$ .

Consider, therefore,  $u \in \mathsf{C}_0^2[0,\infty)$ . Define  $u_\lambda(\mathsf{r}) = u(\mathsf{r})$ . Clearly,  $\lim_{\lambda \to 0} u_\lambda = u$ . On the other hand (see Remark 3),

$$E_{\lambda}\mathrm{e}^{V_{\lambda}(H_{\lambda}\mathsf{r})}\mathsf{L}_{\lambda}u_{\lambda}(\mathsf{r}) = \frac{1}{h_{\lambda}^{2}}\Bigl(\sum_{\mathsf{s}\in\mathbb{Z}_{\lambda}}p_{\lambda}(\mathsf{s}-\mathsf{r})\mathrm{e}^{\frac{V_{\lambda}(H_{\lambda}\mathsf{r})-V_{\lambda}(H_{\lambda}\mathsf{s})}{2}}u(\mathsf{s}) - E_{\lambda}\mathrm{e}^{V_{\lambda}(H_{\lambda}\mathsf{r})}u_{\lambda}(\mathsf{r})\Bigr)$$

$$=\frac{1}{h_{\lambda}^2}\sum_{\mathbf{s}}p_{\lambda}(\mathbf{s}-\mathbf{r})\left(\mathrm{e}^{\frac{V_{\lambda}(H_{\lambda}\mathbf{r})-V_{\lambda}(H_{\lambda}\mathbf{s})}{2}}u(\mathbf{s})-u(\mathbf{r})\right)+\frac{1-E_{\lambda}\mathrm{e}^{V_{\lambda}(H_{\lambda}\mathbf{r})}}{h_{\lambda}^2}u(\mathbf{r}). \quad (2.36)$$

Choose R such that  $\operatorname{supp}(u) \in [0, R]$ . Possibly going to a sub-sequence assume that  $e = \lim e_{\lambda}$  exists. Then, by our assumptions on  $V_{\lambda}$ , the second term converges to (e - q(r))u(r), uniformly in  $r \in [0, R]$ .

As for the first term in (2.36), note that, since  $u_{\lambda}(\mathsf{s}) \equiv 0$  for  $\mathsf{s} > R$  and  $p_{\lambda}(\mathsf{s} - \mathsf{r}) \leq \mathrm{e}^{-cH_{\lambda}|\mathsf{s} - \mathsf{r}|}$ , we may restrict attention to  $\mathsf{r}, \mathsf{s} \leq R + 1$ . But then, again by our assumptions on  $V_{\lambda}$ , the quantity

$$|(V_{\lambda}(H_{\lambda}\mathsf{s}) - V_{\lambda}(H_{\lambda}\mathsf{r}))p_{\lambda}(\mathsf{s} - \mathsf{r})| = h_{\lambda}^{2} |(q(\mathsf{s}) - q(\mathsf{r}))p_{\lambda}(\mathsf{s} - \mathsf{r})| + o(h_{\lambda}^{2}) = o(h_{\lambda}^{2}).$$

Finally, by our assumptions (1.9) and (1.10) on the underlying random walk,

$$\lim_{\lambda \to 0} \frac{1}{h_{\lambda}^2} \sum_{\mathbf{s} \in \mathbb{Z}_{\lambda}} p_{\lambda}(\mathbf{s} - \cdot)(u(\mathbf{s}) - u(\cdot)) = \frac{\sigma^2}{2} u''(\cdot),$$

in the sense of Definition 1.

We have proved:

**Proposition 3.** Under our assumptions on  $V_{\lambda}$  and p, the following convergence (in the sense of Definition 1), holds uniformly in t on compact subsets of  $\mathbb{R}_+$ : If  $\lim_{k\to\infty} e_{\lambda_k} = e$  and  $\lim_{k\to\infty} f_{\lambda_k} = f$ , then

$$\lim_{k \to \infty} \mathsf{T}_{\lambda_k}^{\lfloor H_{\lambda_k}^2 t \rfloor} f_{\lambda_k} = \mathrm{e}^{(\mathsf{L} + \mathsf{el})t} f. \tag{2.37}$$

Convergence of eigenvalues and eigenfunctions.

**Theorem 2.** Under our assumptions on  $V_{\lambda}$  and p,

$$\zeta_0 = \lim_{\lambda \to 0} \mathsf{e}_{\lambda}, \quad \varphi_0 = \lim_{\lambda \to 0} \phi_{\lambda} = \lim_{\lambda \to 0} \phi_{\lambda}^* \quad and \quad \lim_{\lambda \to 0} \frac{c_{\lambda}}{h_{\lambda}} = 1.$$
 (2.38)

*Proof.* By Lemma 1, the set  $\{e_{\lambda}\}$  is bounded and, by Proposition 2, the set  $\{\phi_{\lambda}\}$  is sequentially compact. Let  $\lambda_k \searrow 0$  be a sequence such that both  $e = \lim_{k \to \infty} e_{\lambda_k}$  and  $\varphi = \lim_{k \to \infty} \phi_{\lambda_k}$  exist. Then Proposition 3 implies that

$$\varphi = e^{(\mathsf{L} + \mathsf{el})t} \varphi.$$

By compactness,  $\|\varphi\|_2 = 1$ . In other words,  $\varphi$  is a non-negative normalized  $\mathbb{L}_2(\mathbb{R}_+)$ -eigenfunction of L with eigenvalue -e. Which means that  $\varphi = \varphi_0$  and  $e = \zeta_0$ . Exactly the same argument applies to  $\phi_*^*$ .

e =  $\zeta_0$ . Exactly the same argument applies to  $\phi_{\lambda}^*$ . By construction (see (2.20)),  $1 \equiv c_{\lambda} \sum_{\mathbf{r} \in \mathbb{N}_{\lambda}} \phi_{\lambda}(\mathbf{r}) \phi_{\lambda}^*(\mathbf{r}) = \frac{c_{\lambda}}{h_{\lambda}} \langle \phi_{\lambda}, \phi_{\lambda}^* \rangle_{2,\lambda}$ . Since, by the second assertion of (2.38),  $\lim_{\lambda \to 0} \langle \phi_{\lambda}, \phi_{\lambda}^* \rangle_{2,\lambda} = \|\varphi_0\|_2^2 = 1$ , the last claim of Theorem 2 follows as well.  $\square$ 

Convergence of finite-dimensional distributions. Recall our notations  $\mathbb{P}_{\lambda}$  and  $\mathbb{P}_{\sigma,q}$  for the path measures of the ground-state chain  $X_n$  and the Ferrari-Spohn diffusion x(t). Recall also our rescaling of the ground-state chain:  $x_{\lambda}(t) = h_{\lambda} X_{|H_{\lambda}^2 t|}$ .

**Corollary 1.** For any k, any  $0 < s_1 < s_2 < \cdots < s_k$  and for any collection of bounded continuous functions  $u_0, \ldots, u_k \in \mathsf{C}_\mathsf{b}(\mathbb{R}_+)$ ,

$$\lim_{\lambda \to 0} \mathbb{E}_{\lambda} \{ u_0(x_{\lambda}(0)) u_1(x_{\lambda}(s_1)) \cdots u_k(x_{\lambda}(s_k)) \}$$

$$= \mathbb{E}_{\sigma, q} \{ u_0(x(0)) u_1(x(s_1)) \cdots u_k(x(s_k)) \}. \quad (2.39)$$

*Proof.* Set  $s_0 = 0$  and  $t_i = s_i - s_{i-1}$ . Since  $\mu_{\lambda} = c_{\lambda} \phi_{\lambda}^* \phi_{\lambda}$  and in view of the expressions (2.6) for transition probabilities  $\pi_{\lambda}$  of the ground-state chain, the rightmost asymptotic relation in (2.38), and (1.8) for Ferrari-Spohn semigroups, the target formula (2.39) can be written as

$$\lim_{\lambda \to 0} h_{\lambda} \sum_{\mathbf{r} \in \mathbb{N}_{\lambda}} \phi_{\lambda}^{*}(\mathbf{r}) u_{\lambda,0}(\mathbf{r}) \mathsf{T}_{\lambda}^{\lfloor H_{\lambda}^{2} t_{1} \rfloor} \left( u_{\lambda,1} \mathsf{T}_{\lambda}^{\lfloor H_{\lambda}^{2} t_{2} \rfloor} \left( u_{\lambda,2} \cdots \mathsf{T}_{\lambda}^{\lfloor H_{\lambda}^{2} t_{k} \rfloor} \left( u_{\lambda,k} \phi_{\lambda} \right) \cdots \right) \right) (\mathbf{r})$$

$$= \int_{0}^{\infty} \varphi_{0}(r) u_{0}(r) \mathsf{T}^{t_{1}} \left( u_{1} \mathsf{T}^{t_{2}} \left( u_{2} \cdots \mathsf{T}^{t_{k}} \left( u_{k} \varphi_{0} \right) \cdots \right) \right) (r) \, \mathrm{d}r, \quad (2.40)$$

where  $u_{\lambda,i}$  and  $u_i$  coincide on  $\mathbb{N}_{\lambda}$ . Theorem 2 implies that  $\lim_{\lambda \to 0} \phi_{\lambda} = \varphi_0$  and  $\lim_{\lambda \to 0} \phi_{\lambda}^* = \varphi_0$ . Hence, by induction, (2.40) is a consequence of Proposition 3 and the following two elementary facts:

- (a) If  $\lim v_{\lambda} = v$  and  $u_{\lambda}(\mathbf{r}) = u(\mathbf{r})$  with u being a bounded continuous function, then  $\lim_{\lambda \to 0} v_{\lambda} u_{\lambda} = v u$ .
- (b) If  $\lim u_{\lambda} = u$  and  $\lim v_{\lambda} = v$ , then  $\lim_{\lambda \to 0} \langle u_{\lambda}, v_{\lambda} \rangle_{2,\lambda} = \langle u, v \rangle_{2}$ .  $\square$

#### 3. Probabilistic tools

The derivations of the probabilistic estimates given below are based on the techniques and ideas developed in [14]. Nevertheless, because our setting is slightly different and for completeness, we provide detailed proofs. In addition, one of the needed claims from [14] (Theorem 1.2 therein) contains a mistake, which we correct here.

Recall our notation  $\hat{\mathcal{P}}_{N,+}^{\mathsf{u},\mathsf{v}} = \mathcal{P}_{1,N,+}^{\mathsf{u},\mathsf{v}}$ . As before, given a path  $\mathbb{X} = (X_1,\ldots,X_N)$ , set  $\mathsf{p}(\mathbb{X}) = \prod_{i=1}^{N-1} p_{X_{i+1}-X_i}$ . Define  $\hat{\mathcal{P}}_{N,+}^{\mathsf{u},\mathscr{S}} = \cup_{\mathsf{v}\in\mathbb{Z}_+} \hat{\mathcal{P}}_{N,+}^{\mathsf{u},\mathsf{v}}$  and consider the partition functions

$$\hat{Z}_{N,+,\lambda}^{\mathsf{u},\varnothing} = \sum_{\mathbb{X} \in \hat{\mathcal{P}}_{N,+}^{\mathsf{u},\varnothing}} e^{-\sum_{i=1}^{N} V_{\lambda}(X_{i})} \, \mathsf{p}(\mathbb{X}).$$

More generally, given any subset  $\mathcal{C} \subset \hat{\mathcal{P}}_{N,+}^{\mathsf{u},\varnothing}$ , we denote by

$$\hat{Z}_{N,+,\lambda}^{\mathsf{u},\varnothing}[\mathcal{C}] = \sum_{\mathbb{X} \in \mathcal{C}} e^{-\sum_{i=1}^N V_\lambda(X_i)} \, \mathsf{p}(\mathbb{X}),$$

the partition function restricted to paths satisfying the constraint C.

3.1. Proof of Lemma 1. The proof will rely on the following identity which, exactly as (2.2), is straightforward from the very definition of  $\tilde{\mathsf{T}}_{\lambda}$  in (2.1):

$$\hat{Z}_{N,+,\lambda}^{\mathsf{u},\varnothing} = \mathrm{e}^{-\frac{1}{2}V_{\lambda}(\mathsf{u})} \tilde{\mathsf{T}}_{\lambda}^{N}[f_{\lambda}](\mathsf{u}),$$

where  $f_{\lambda}(x) = e^{-\frac{1}{2}V_{\lambda}(x)}$ . Since  $f_{\lambda}$  is positive,

$$\log E_{\lambda} = \lim_{N \to \infty} \frac{1}{N} \log \hat{Z}_{N,+,\lambda}^{\mathsf{u},\varnothing}, \tag{3.1}$$

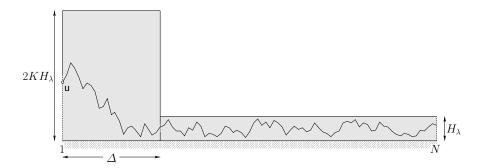


Fig. 2: The construction for the lower bound in the proof of Lemma 1.

for all  $u \in \mathbb{Z}^+$ . In particular, the claim of Lemma 1 will follow from lower and upper bounds on  $\hat{Z}_{N,+,\lambda}^{u,\varnothing}$  for finite values of N and  $\lambda$ . In the sequel, we shall allow rather general values of the boundary condition u. Of course, to derive the claim of Lemma 1, we could as well take u=0.

We shall compare the tilted partition functions  $\hat{Z}_{N,+,\lambda}^{\mathsf{u},\varnothing}$  and  $\hat{Z}_{N,+,0}^{\mathsf{u},\varnothing}$ . The latter equals to the probability that the random walk starting at  $\mathsf{u}$  stays positive for first N steps of its life. This probability is evidently non-decreasing with  $\mathsf{u}$  and, as is well known (see for instance [3]), it is of order  $N^{-1}$  for  $\mathsf{u} = 1$ . In particular,

$$\lim_{N \to \infty} \frac{1}{N} \log \hat{Z}_{N,+,0}^{\mathsf{u},\varnothing} = 0, \tag{3.2}$$

uniformly in  $u \in \mathbb{N}$ .

Lower bound on  $\hat{Z}_{N,+,\lambda}^{\mathsf{u},\varnothing}$  and upper bound on  $\mathsf{e}_\lambda$ . We claim that there exist finite constants  $\overline{\mathsf{e}}$  and  $c_1$  such that, for any  $K \geq 1$  fixed,

$$\hat{Z}_{N,+,\lambda}^{\mathsf{u},\varnothing} \ge e^{-\overline{\mathsf{e}}NH_{\lambda}^{-2} - c_1K\sqrt{q(2K)}} \, \hat{Z}_{N,+,0}^{\mathsf{u},\varnothing},\tag{3.3}$$

uniformly in  $\lambda$  small,  $0 \le \mathsf{u} \le KH_{\lambda}$  and

$$N \gg \Delta = \Delta(K, \lambda) = \frac{KH_{\lambda}^2}{\sqrt{q(2K)}}.$$
 (3.4)

In view of (3.1) and (3.2), this implies that  $e_{\lambda} = -H_{\lambda}^2 \log E_{\lambda} \leq \overline{e}$  for all  $\lambda$  sufficiently small.

In order to check (3.3), we restrict the partition function to trajectories made of two pieces (see Figure 2). The left part is used to bring the interface below  $H_{\lambda}$ ; in the remaining piece, the interface remains inside a tube of height  $H_{\lambda}$ .

We consider the events  $^1$ 

$$\begin{split} \mathcal{D}_L &= \big\{ \mathbb{X} \in \hat{\mathcal{P}}_{N,+}^{\mathtt{u},\varnothing} \, : \, \max_{i \in \{1,\dots,\Delta\}} X_i \leq 2KH_{\lambda}, X_{\Delta} \in \big[\frac{1}{3}H_{\lambda}, \frac{2}{3}H_{\lambda}\big] \big\}, \\ \mathcal{D}_M &= \big\{ \mathbb{X} \in \hat{\mathcal{P}}_{N,+}^{\mathtt{u},\varnothing} \, : \, \max_{i \in \{\Delta,\dots,N\}} X_i \leq H_{\lambda} \big\}. \end{split}$$

 $<sup>^1</sup>$  Here and several times in the sequel, we assume numbers like  $\varDelta$  to be integers whenever it is desirable.

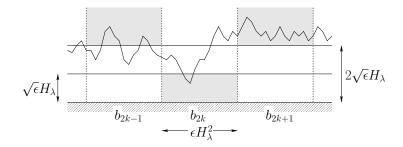


Fig. 3: The event  $\mathcal{B}_k$  occurs if the path visits both leftmost and rightmost shaded areas. The event  $C_k$  occurs if, in addition, it also visits the third one.

Then

$$\hat{Z}_{N,+,\lambda}^{\mathrm{u},\varnothing} \geq e^{-\Delta V_{\lambda}(2KH_{\lambda})-(N-\Delta)V_{\lambda}(H_{\lambda})} \operatorname{p}(\mathcal{D}_{L}\cap\mathcal{D}_{M}\,|\,\hat{\mathcal{P}}_{N,+}^{\mathrm{u},\varnothing})\,\hat{Z}_{N,+,0}^{\mathrm{u},\varnothing}.$$

By the assumptions (1.14) and (1.15),  $H_{\lambda}^2 V_{\lambda}(rH_{\lambda}) < 2q(r)$  uniformly in  $r \in$ [0,2K], for all  $\lambda$  sufficiently small. Hence, for such  $\lambda$ , the exponent in the righthand side is bounded below by  $e^{-2K\sqrt{q(2K)}-2NH_{\lambda}^{-2}}$ 

It remains to estimate  $p(\mathcal{D}_L \cap \mathcal{D}_M | \mathcal{P}_{N,+}^{u,\varnothing})$ . By the invariance principle (for a random walk conditioned to stay positive; see first [4, Theorem 1] and then [8, Theorem 1.1]),

$$\liminf_{N \to \infty} p(\mathcal{D}_L \,|\, \mathcal{P}_{N,+}^{\mathsf{u},\varnothing}) \ge e^{-c_2 K \sqrt{q(2K)}}.$$

for some absolute constant  $c_2 > 0$ , provided that  $\lambda$  be small enough. On the other hand, letting  $\mathcal{D}' = \{\sup_{0 \leq i \leq H_{\lambda}^2} X_i \leq H_{\lambda}\} \cap \{X_{\lceil H_{\lambda}^2 \rceil} \in [\frac{1}{3}H_{\lambda}, \frac{2}{3}H_{\lambda}]\},$ it follows from the Markov property that

$$\inf_{\ell \in \left[\frac{1}{3}H_{\lambda}, \frac{2}{3}H_{\lambda}\right]} \mathsf{p}(\mathcal{D}_{M} \mid X_{L} = \ell, X_{i} \geq 0 \; \forall \Delta \leq i \leq N)$$

$$\geq \left\{ \inf_{\ell \in \left[\frac{1}{3}H_{\lambda}, \frac{2}{3}H_{\lambda}\right]} \mathsf{p}(\mathcal{D}' \mid X_{0} = \ell, X_{i} \geq 0 \; \forall \Delta \leq i \leq H_{\lambda}^{2}) \right\}^{\lceil N/H_{\lambda}^{2} \rceil}$$

$$\geq e^{-c_{3}NH_{\lambda}^{-2}}.$$

Upper bound on  $\hat{Z}_{N,+,\lambda}^{u,\varnothing}$  and lower bound on  $e_{\lambda}$ . We claim that there exist  $\bar{\lambda}>0$ and a positive constant e such that

$$\hat{Z}_{N,+,\lambda}^{\mathrm{u},\varnothing} \le e^{-\underline{\mathbf{e}}NH_{\lambda}^{-2}}\hat{Z}_{N,+,0}^{\mathrm{u},\varnothing}, \tag{3.5}$$

uniformly in  $\mathbf{u} \geq 0$ ,  $N \geq H_{\lambda}^2$  and  $\lambda < \bar{\lambda}$ .

Let us fix some small  $\epsilon > 0$  (which does not have to be very small; one can optimize over it at the end of the proof). The idea behind the proof is that a typical trajectory has many disjoint segments of the length at least  $\epsilon H_{\lambda}^2$ , which are at a distance at least  $\sqrt{\epsilon}H_{\lambda}$  from the wall.

We partition the interval  $\{1,\ldots,N\}$  into  $N_{\lambda}$  disjoint intervals  $b_1,\ldots,b_{N_{\lambda}}$  of length  $\epsilon H_{\lambda}^2$  and, possibly, one additional shorter rightmost interval.

We say that the event  $\mathcal{B}_k$  occurs if (see Figure 3)

$$\max_{i \in b_{2k-1}} X_i > 2\sqrt{\epsilon} H_{\lambda} \quad \text{and} \quad \max_{i \in b_{2k+1}} X_i > 2\sqrt{\epsilon} H_{\lambda}. \tag{3.6}$$

Let us denote by G the number of indices k for which the event  $\mathcal{B}_k$  occurs. It follows from the CLT that there exists  $\kappa_1 > 0$  such that

$$\inf_{\mathsf{v},\mathsf{w}\geq 0} \mathsf{p}\left(\max_{1\leq i\leq \epsilon H_{\lambda}^{2}} X_{i} > 2\sqrt{\epsilon}H_{\lambda} \mid \mathcal{P}_{1,\epsilon H_{\lambda}^{2},+}^{\mathsf{v},\mathsf{w}}\right) > \kappa_{1}. \tag{3.7}$$

Observe that the events  $\{\max_{i \in b_{2j-1}} X_i > 2\sqrt{\epsilon}H_{\lambda}\}, j = 1, \ldots, N_{\lambda}/2$ , are conditionally independent given the trajectories in the intervals  $b_{2k}$ . As a result, (3.7) implies that there exists  $\kappa_2 > 0$  such that

$$\mathsf{p}(G \leq \frac{1}{8}\kappa_1^2 N_\lambda \,|\, \hat{\mathcal{P}}_{N,+}^{\mathsf{u},\varnothing}) \leq e^{-\kappa_2 N_\lambda}, \tag{3.8}$$

uniformly in  $\boldsymbol{u}$ .

Similarly, let us say that the event  $C_k$  occurs if  $B_k$  occurs and (see Figure 3)

$$\min_{i \in b_{2k}} X_i < \sqrt{\epsilon} H_{\lambda},$$

and let us denote by G' the number of indices such that  $C_k$  occurs.

The occurrence of  $C_k$  enforces a downward fluctuation at least as large as  $\sqrt{\epsilon}H_{\lambda}$  on a time interval of length at most  $3\epsilon H_{\lambda}^2$ . The functional CLT [9, Theorem 2.4] implies that such an event has probability at most  $\kappa_3$ , for some  $\kappa_3 < 1$ , uniformly in  $\lambda$  small. This implies that there exists  $\kappa_4 > 0$ , such that

$$p(G' \ge \frac{1+\kappa_3}{2}g \mid \hat{\mathcal{P}}_{N,+}^{\mathsf{u},\varnothing}; G = g) \le e^{-\kappa_4 g}, \tag{3.9}$$

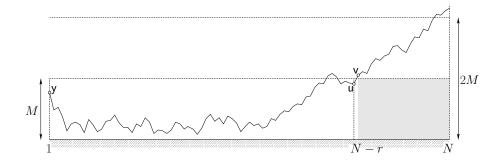
uniformly in u and g. Altogether, (3.8) and (3.9) yield

$$\mathsf{p}\left(G - G' \le \frac{1 - \kappa_3}{8} \kappa_1^2 N_\lambda \mid \hat{\mathcal{P}}_{N,+}^{\mathsf{u},\varnothing}\right) \le e^{-\kappa_5 N_\lambda} \le e^{-\kappa_6 \epsilon^{-1} H_\lambda^{-2} N}. \tag{3.10}$$

The quantity G - G' provides a lower bound on the number of disjoint intervals  $b_{2k}$  of length  $\epsilon H_{\lambda}^2$  such that  $\min_{i \in b_{2k}} X_i \geq \sqrt{\epsilon} H_{\lambda}$ . Therefore,

$$\sum_{i=1}^{N} V_{\lambda}(X_i) \ge (G - G')\epsilon H_{\lambda}^2 V(\sqrt{\epsilon}H_{\lambda}) \ge (G - G')\epsilon q_0(\sqrt{\epsilon}) \ge \kappa_7 q_0(\sqrt{\epsilon}) H_{\lambda}^{-2} N,$$

whenever  $G - G' \ge \kappa_7 N_\lambda$ . Take  $\kappa_7 = \frac{1-\kappa_3}{8} \kappa_1^2$ . The conclusion (3.5) follows from (3.10) and (3.11).  $\square$ 



**Fig. 4:** The last exit decomposition in (3.14). After time N-r, the path cannot visit the shaded area and has to end up above level 2M.

3.2. Proof of Lemma 2. We shall prove Lemma 2 only for  $\phi_{\lambda}^*$ . The proof for  $\phi_{\lambda}$  is a literal repetition for reversed walks. For the sake of notations, we shall think of  $T_{\lambda}$  in (2.5) as acting on non-rescaled spaces  $\ell_2(\mathbb{N})$ , with the norm

$$\langle u, v \rangle_{2,\lambda} = h_{\lambda} \sum_{\mathbf{r} \in \mathbb{N}} u(\mathbf{r}) v(\mathbf{r}).$$

Compare with (2.18).

Similarly, we shall think of  $\phi_{\lambda}$  and  $\phi_{\lambda}^{*}$  as of functions on  $\mathbb{N}$ . Recall the normalizing constant  $c_{\lambda}$  which was introduced in (2.20), and recall that  $\mu_{\lambda}(\mathsf{x}) = c_{\lambda}\phi_{\lambda}(\mathsf{x})\phi_{\lambda}^{*}(\mathsf{x})$  is the invariant measure of the positively recurrent chain on  $\mathbb{N}$  with transition probabilities  $\pi_{\lambda}$  specified in (2.6). Define  $g_{M}(\mathsf{x}) = \mathbf{1}_{\{x>M\}} = \sum_{\mathsf{x}>M} \mathbf{1}_{\mathsf{x}}$ . Then,

$$\lim_{N \to \infty} \mathsf{T}_{\lambda}^{N} g_{M}(\cdot) = \phi_{\lambda}(\cdot) \lim_{N \to \infty} \pi_{\lambda}^{N} \left[ \sum_{\mathsf{x} > M} \frac{\mathbf{1}_{\mathsf{x}}}{\phi_{\lambda}(\mathsf{x})} \right] (\cdot) = \phi_{\lambda}(\cdot) \sum_{\mathsf{x} > M} \frac{\mu_{\lambda}(\mathsf{x})}{\phi_{\lambda}(\mathsf{x})}$$
$$= c_{\lambda} \phi_{\lambda}(\cdot) \sum_{\mathsf{x} > M} \phi_{\lambda}^{*}(\mathsf{x}), \tag{3.12}$$

for any  $\lambda > 0$ .

For  $\mathsf{v} > M$  and  $k \ge 0$  let  $\mathcal{Q}_{k,+}^{\mathsf{v},M}$  be the family of k-step paths  $\mathbb{X} = (\mathsf{x}_0,\ldots,\mathsf{x}_k)$  which start at  $\mathsf{v}$ ,  $\mathsf{x}_0 = \mathsf{v}$ , stay above level M, and end up above level 2M,  $\mathsf{x}_k > 2M$ . We employ the notation (see (2.5))

$$\mathsf{T}_{\lambda}^{k} \{ \mathcal{Q}_{k,+}^{\mathsf{v},M} \} = \sum_{\mathbb{X} \in \mathcal{Q}_{k,+}^{\mathsf{v},M}} \prod_{1}^{k} \mathsf{T}_{\lambda}(\mathsf{x}_{i-1},\mathsf{x}_{i}). \tag{3.13}$$

By convention,  $\mathsf{T}^0_{\lambda}\{\mathcal{Q}^{\mathsf{v},M}_{0,+}\} = \mathbf{1}_{\{\mathsf{v}>2M\}}.$ 

Let us fix  $0 < \mathsf{y} < M$  and consider paths  $\mathbb{X} \in \hat{\mathcal{P}}_{N,+}^{\mathsf{y},\emptyset}$  ending up above level  $2M, X_N > 2M$ . By the last exit decomposition from  $\{1, \dots, M\}$  (see Figure 4),

$$\mathsf{T}_{\lambda}^{N} g_{2M}(\mathsf{y}) = \sum_{\mathsf{u} \leq M} \sum_{\mathsf{v} > M} \sum_{r=1}^{N-1} \mathsf{T}_{\lambda}^{N-r-1} \left[ \mathbf{1}_{\mathsf{u}} \right] (\mathsf{y}) \mathsf{T}_{\lambda} (\mathsf{u}, \mathsf{v}) \mathsf{T}_{\lambda}^{r-1} \left\{ \mathcal{Q}_{r-1,+}^{\mathsf{v}, M} \right\} \\
= \sum_{\mathsf{u} \leq M} \sum_{\mathsf{v} > M} \sum_{r=1}^{N-1} \mathsf{T}_{\lambda}^{N-r-1} \left[ \mathbf{1}_{\mathsf{u}} \right] (\mathsf{y}) E_{\lambda}^{-1} p_{\mathsf{v}-\mathsf{u}} e^{-\frac{V_{\lambda} (\mathsf{u}) + V_{\lambda} (\mathsf{v})}{2}} \mathsf{T}_{\lambda}^{r-1} \left\{ \mathcal{Q}_{r-1,+}^{\mathsf{v}, M} \right\}.$$
(3.14)

Taking the limit  $N \to \infty$ , we infer from (3.12), (3.14) and positivity of  $\phi_{\lambda}$  and  $c_{\lambda}$  that

$$\sum_{\mathsf{x}>2M} \phi_{\lambda}^*(\mathsf{x}) = \sum_{\mathsf{u}\leq M} \phi_{\lambda}^*(\mathsf{u}) \sum_{\mathsf{v}>M} E_{\lambda}^{-1} p_{\mathsf{v}-\mathsf{u}} \mathrm{e}^{-\frac{V_{\lambda}(\mathsf{u}) + V_{\lambda}(\mathsf{v})}{2}} \sum_{r>0} \mathsf{T}_{\lambda}^r \{\mathcal{Q}_{r,+}^{\mathsf{v},M}\}. \tag{3.15}$$

Let us try to derive an upper bound on

$$\max_{\mathbf{u} \le M} \sum_{\mathbf{v} > M} E_{\lambda}^{-1} p_{\mathbf{v} - \mathbf{u}} e^{-\frac{V_{\lambda}(\mathbf{u}) + V_{\lambda}(\mathbf{v})}{2}} \sum_{r > 0} \mathsf{T}_{\lambda}^{r} \{ \mathcal{Q}_{r,+}^{\mathbf{v}, M} \}. \tag{3.16}$$

For simplicity, we shall prove directly the second inequality in (2.25). The arguments rely on the lower bound  $H_{\lambda}^2 V_{\lambda}(M) \geq q_0(h_{\lambda}M)$ . If instead we keep track of the original quantity  $H_{\lambda}^2 V_{\lambda}(M)$ , then the first inequality in (2.25) will follow.

By Lemma 1 and by our assumptions on p and  $V_{\lambda}$ ,

$$E_{\lambda}^{-1} p_{\mathsf{v}-\mathsf{u}} \mathrm{e}^{-\frac{V_{\lambda}(\mathsf{u}) + V_{\lambda}(\mathsf{v})}{2}} \le \exp\left\{-c_1 h_{\lambda}^2 \left(q_0(h_{\lambda} M) - 1\right) - c_2(\mathsf{v} - M)\right\}. \tag{3.17}$$

On the other hand, Lemma 1 and crude estimates on the values of the potential V above level M and of the hitting probability of the half-line  $\{2M, 2M+1, \ldots\}$  by an r-step random walk which starts at  $\mathsf{v}$  imply that, for r>1,

$$\mathsf{T}_{\lambda}^{r}\{\mathcal{Q}_{r,+}^{\mathsf{v},M}\} \leq E_{\lambda}^{-r} \mathrm{e}^{-rh_{\lambda}^{2}q_{0}(h_{\lambda}M)} \mathsf{p}(X_{r} > 2M - \mathsf{v}) 
\leq \exp\left\{-c_{3}h_{\lambda}^{2}r\left(q_{0}(h_{\lambda}M) - 1\right) - c_{4}\frac{(2M - \mathsf{v})_{+}^{2}}{r} \wedge (2M - \mathsf{v})_{+}\right\}.$$
(3.18)

Indeed, the second term in the exponent on the right-hand side above follows from the exponential Markov inequality  $p(X_r > a) \le e^{-c_4 \frac{a_+^2}{r} \wedge a_+}$ .

The right-hand sides of both (3.17) and (3.18) are already independent of u. Let us sum over v > M. If r = 0, then v has to satisfy v > 2M. Using (3.17),

$$\sum_{\mathbf{v} > 2M} E_{\lambda}^{-1} p_{\mathbf{v} - \mathbf{u}} e^{-\frac{V_{\lambda}(\mathbf{u}) + V_{\lambda}(\mathbf{v})}{2}} \le \exp\left\{-c_5 \left(h_{\lambda}^2 \left(q_0(h_{\lambda} M) - 1\right) + M\right)\right\}. \tag{3.19}$$

For r > 0, we take advantage of both upper bounds (3.17) and (3.18) above:

$$\sum_{\mathsf{v}>M} e^{-c_2(\mathsf{v}-M) - c_4 \frac{(2M-\mathsf{v})_+^2}{r} \wedge (2M-\mathsf{v})_+} \le e^{-c_6 \frac{M^2}{r} \wedge M}. \tag{3.20}$$

Putting things together for  $M = H_{\lambda}K$ , we conclude that the expression in (3.16) is bounded above by

$$e^{-c_7 H_{\lambda} K} + \sum_{r \ge 1} \exp \left\{ -c_8 \left( h_{\lambda}^2 r q_0(K) + \frac{(H_{\lambda} K)^2}{r} \wedge H_{\lambda} K \right) \right\} \le c_{10} e^{-c_9 K (\sqrt{q_0(K)} \wedge H_{\lambda})}, \tag{3.21}$$

uniformly in K > 0 and  $\lambda$  sufficiently small.

Coming back to (3.15), we infer: For any K > 0 fixed,

$$\sum_{\mathsf{x}>H_{\lambda}K} \phi_{\lambda}^{*}(\mathsf{x}) \le c_{11} e^{-c_{12}K} \sqrt{q_{0}(K)} \sum_{\mathsf{x}} \phi_{\lambda}^{*}(\mathsf{x}), \tag{3.22}$$

for all  $\lambda < \lambda_0(K)$ . Notice that  $\sum_{\mathbf{x}} \phi_{\lambda}^*(\mathbf{x}) < \infty$  by compactness of  $\mathsf{T}_{\lambda}$ .

Let us return to our basic rescaling (2.19) of  $\phi_{\lambda}$  and  $\phi_{\lambda}^{*}$  as unit norm elements of  $\ell_{2}(\mathbb{N}_{\lambda})$ . The bound (3.22) can be rewritten as

$$h_{\lambda} \sum_{\mathsf{r}>K} \phi_{\lambda}^{*}(\mathsf{r}) \le c_{11} e^{-c_{12}K\sqrt{q_{0}(K)}} \|\phi_{\lambda}^{*}\|_{1,\lambda}.$$
 (3.23)

Since  $h_{\lambda} \sum_{\mathbf{r} \leq K} \phi_{\lambda}^{*}(\mathbf{r}) \leq \sqrt{K} \|\phi_{\lambda}^{*}\|_{2,\lambda} = \sqrt{K}$ , we conclude that  $\{\|\phi_{\lambda}^{*}\|_{1,\lambda}\}$  is a bounded sequence. The bound (2.26) and, in view of (3.23), also (2.25) follow.  $\square$ 

3.3. Tightness of  $(x_{\lambda}, \mathbb{P}_{\lambda})$ . Fix any  $T < \infty$  and consider the family of rescaled processes  $x_{\lambda}$  defined in (1.17). Precisely,  $x_{\lambda}$  is the linear interpolation of the rescaled stationary ergodic ground-state chain  $X = X^{\lambda}$  with transition probabilities  $\pi_{\lambda}$  defined in (2.6) and invariant distribution  $\mu_{\lambda}$ . With a slight abuse of notation, we shall continue to use  $\mathbb{P}_{\lambda}$  for the induced distribution of  $x_{\lambda}(\cdot)$  on  $\mathbb{C}[-T,T]$ .

**Proposition 4.** The family  $(x_{\lambda}, \mathbb{P}_{\lambda})$  is tight on C[-T, T].

Proof of Proposition 4. Recall that the invariant measure  $\mu_{\lambda}$  at  $\lambda > 0$  is given by  $\mu_{\lambda} = c_{\lambda}\phi_{\lambda}\phi_{\lambda}^{*}$ . Thus, by Theorem 2, the sequence  $\{x_{\lambda}(0)\}$  is tight. It remains to show that, for each  $\epsilon, \nu > 0$ , there exists  $\delta > 0$  such that

$$\mathbb{P}_{\lambda}\left(\max_{0 < t < \delta} |x_{\lambda}(t) - x_{\lambda}(0)| > \epsilon\right) \le \nu \delta. \tag{3.24}$$

For any event  $\mathcal{A} \in \sigma(X_n: 0 \leq n \leq \delta H_{\lambda}^2)$ , and for any  $\mathsf{x}, \mathsf{y} \in \mathbb{N}$ , let us define  $\mathcal{A}_{\mathsf{x},\mathsf{y}} = \left\{ \mathbb{X} \subset \mathcal{A}: X_0 = \mathsf{x}, X_{\delta H_{\lambda}^2} = \mathsf{y} \right\}$ . As in (3.13), we employ the notation

$$\mathsf{T}_{\lambda}^{\delta H_{\lambda}^{2}}\{\mathcal{A}_{\mathsf{x},\mathsf{y}}\} = \sum_{\mathbb{X}\in\mathcal{A}_{\mathsf{x},\mathsf{y}}} \prod_{1}^{\delta H_{\lambda}^{2}} \mathsf{T}_{\lambda}(\mathsf{x}_{i-1},\mathsf{x}_{i}). \tag{3.25}$$

for the corresponding restricted partition function. In this way,

$$\mathbb{P}_{\lambda}(\mathcal{A}) = c_{\lambda} \sum_{\mathsf{x}, \mathsf{y} \in \mathbb{N}} \phi_{\lambda}^{*}(\mathsf{x}) \mathsf{T}_{\lambda}^{\delta H_{\lambda}^{2}} \{ \mathcal{A}_{\mathsf{x}, \mathsf{y}} \} \phi_{\lambda}(\mathsf{y}). \tag{3.26}$$

Since  $V_{\lambda} \geq 0$ ,

$$\mathsf{T}_{\lambda}^{\delta H_{\lambda}^{2}}\{\mathcal{A}_{\mathsf{x},\mathsf{y}}\} \leq \mathrm{e}^{-\delta H_{\lambda}^{2}\log E_{\lambda}}\mathsf{p}(\mathcal{A}_{\mathsf{x},\mathsf{y}}).$$

By Lemma 1,  $\{-\mathbf{e}_{\lambda} = H_{\lambda}^2 \log E_{\lambda}\}$  is a bounded sequence. In the case of

$$\mathcal{A} = \left\{ \max_{0 \le n \le \delta H_{\lambda}^2} |X_n - X_0| > \epsilon H_{\lambda} \right\},\,$$

the upper bound on the probabilities

$$p(\mathcal{A}_{x,y}) \le \kappa_1 \frac{h_{\lambda}}{\sqrt{\delta}} e^{-\kappa_2 \frac{\epsilon^2}{\delta} \wedge (\epsilon H_{\lambda})}$$
(3.27)

holds uniformly in  $\delta, \epsilon > 0$ ,  $x, y \in \mathbb{N}$  and  $\lambda$  small. By Theorem 2,  $c_{\lambda}/h_{\lambda}$  is bounded. Putting things together, we infer that

$$\begin{split} \frac{1}{\delta} \mathbb{P}_{\lambda} \Big( \max_{0 \leq t \leq \delta} |x_{\lambda}(t) - x_{\lambda}(0)| > \epsilon \Big) &\leq \kappa_{3} \mathrm{e}^{-\kappa_{2} \frac{\epsilon^{2}}{\delta} \wedge (\epsilon H_{\lambda}) - \frac{3}{2} \log \delta} h_{\lambda}^{2} \sum_{\mathsf{x}, \mathsf{y} \in \mathbb{N}} \phi_{\lambda}^{*}(\mathsf{x}) \phi_{\lambda}(\mathsf{y}) \\ &= \kappa_{3} \mathrm{e}^{-\kappa_{2} \frac{\epsilon^{2}}{\delta} \wedge (\epsilon H_{\lambda}) - \frac{3}{2} \log \delta} \|\phi_{\lambda}^{*}\|_{1,\lambda} \|\phi_{\lambda}\|_{1,\lambda}, \quad (3.28) \end{split}$$

uniformly in  $\delta, \epsilon > 0$  and  $\lambda$  small. Since  $\|\phi_{\lambda}\|_{1,\lambda}$  and  $\|\phi_{\lambda}^*\|_{1,\lambda}$  are bounded and since  $H_{\lambda}$  is bounded away from zero, we are home.  $\square$ 

3.4. Asymptotic ground-state structure of  $\mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}$ . Let us fix C>0 and T>1. For  $\lambda>0$ ,  $\mathsf{u},\mathsf{v}\leq CH_\lambda$  and  $N>2TH_\lambda^2$ , we are going to compare the restriction  $\mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}$  of  $\mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}$  to the  $\sigma$ -algebra

$$\mathcal{F}_{\lambda,T} = \sigma(X_i : -TH_{\lambda}^2 \le i \le TH_{\lambda}^2)$$

with the restriction  $\mathbb{P}_{\lambda}^{T}$  of  $\mathbb{P}_{\lambda}$  to  $\mathcal{F}_{\lambda,T}$ .

**Proposition 5.** There exists  $c_1 > 0$  and  $K = K(C,T) < \infty$  such that

$$\|\mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v},T} - \mathbb{P}_{\lambda}^{T}\|_{\mathrm{Var}} \le 2e^{-c_{1}NH_{\lambda}^{-2}},\tag{3.29}$$

uniformly in  $\lambda$  small,  $N > (T+K)H_{\lambda}^2$  and  $u, v \leq CH_{\lambda}$ . Above,  $\|\cdot\|_{Var}$  is the total variational norm.

As an immediate consequence, we deduce the following

Corollary 2. Let  $\lambda_N$  be a sequence satisfying the assumptions of Theorem A. Let  $C, T < \infty$  be fixed and assume that the sequences  $u_N, v_N \in \mathbb{N}$  satisfy  $u_N, v_N \leq CH_{\lambda_N}$ . Consider the sequence of processes  $x_N(\cdot) = x_{\lambda_N}(\cdot)$  defined via linear interpolation from (1.17). With a slight abuse of notation, let  $\mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v},\mathsf{v}}$  and  $\mathbb{P}_{\lambda_N}^T$  denote the induced distributions on  $\mathsf{C}[-T,T]$  of  $\mathbb{P}_{N,+,\lambda_N}^{\mathsf{u},\mathsf{v}}$  and, respectively, of the direct ground-state chain measure  $\mathbb{P}_{\lambda_N}$ . Then,

$$\lim_{N \to \infty} \|\mathbb{P}_{N,+,\lambda_N}^{\mathsf{u,v,T}} - \mathbb{P}_{\lambda_N}^T\|_{\mathsf{Var}} = 0. \tag{3.30}$$

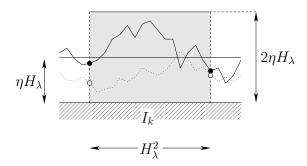


Fig. 5: The interval  $I_k$  is  $\eta$ -good if both paths  $X^1$  and  $X^2$  stay inside the shaded area and take values smaller than  $\eta H_{\lambda}$  at the boundaries of the interval.

Proof (Proof of Proposition 5). We shall use a coupling argument, considering two independent copies  $X^1$  and  $X^2$  of the process, with possibly different boundary conditions.

The first step is to show that we can typically find many pieces of tubes of length  $H^2_{\lambda}$  and height of order  $H_{\lambda}$  inside which both paths are confined.

Let us first split the interval of length 2N + 1 into

$$m = |(2N+1)/H_{\lambda}^{2}|$$

consecutive disjoint intervals  $I_1, I_2, \dots, I_m$ , of length  $H^2_{\lambda}$  (plus, possibly, a final interval of shorter length). We say that the interval  $I_k$  is  $\eta$ -good if (see Figure 5)

$$\max_{i \in I_k} X_i^1 < 2\eta H_\lambda, \quad \max_{i \in I_k} X_i^2 < 2\eta H_\lambda$$

and the values of  $X_i$ ; i = 1, 2, at the end-points of  $I_k$  are less than  $\eta H_{\lambda}$ .

**Lemma 3.** Given the realizations of the two paths  $X^1$  and  $X^2$ , let us denote by M the number of  $\eta$ -good intervals of the form  $I_{3k+2}$ ,  $0 \le k < m/3$ . Then, there exist  $\eta$ ,  $c_2 > 0$ ,  $\rho > 0$  and  $K_0 < \infty$  such that

$$\mathbb{P}_{N+\lambda}^{0,0} \otimes \mathbb{P}_{N+\lambda}^{\mathsf{u},\mathsf{v}}(M < \rho \frac{m}{3}) \leq e^{-c_2 N H_{\lambda}^{-2}},$$

uniformly in  $\lambda$  small,  $0 \le u, v \le CH_{\lambda}$  and  $N \ge K_0H_{\lambda}^2$ .

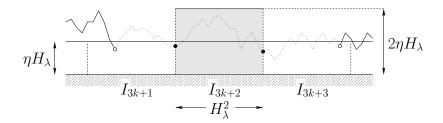
*Proof.* We first show that it is very unlikely for  $X^1$  or  $X^2$  to stay far away from the wall for a long time. Indeed, let us write B for the number of intervals  $I_k$  such that  $\min_{i \in I_k} X_i > \eta H_{\lambda}$ . Then, for any  $\epsilon > 0$ , there exists  $\eta(C, \epsilon)$  such that for all  $\eta > \eta(C, \epsilon)$ ,

$$\mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}(B > \epsilon m) \le e^{-c_3 N H_{\lambda}^{-2}},\tag{3.31}$$

for some constant  $c_3 > 0$ , uniformly in  $0 \le \mathsf{u}, \mathsf{v} \le CH_\lambda$ . Indeed, on the event  $B > \epsilon m$ ,

$$\sum_{i=-N}^{N} V_{\lambda}(X_i) \ge \epsilon m H_{\lambda}^2 V_{\lambda}(\eta H_{\lambda}) \ge \epsilon q_0(\eta) (2N H_{\lambda}^{-2} - 1),$$

which provides an upper bound on  $Z_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}[B>\epsilon m].$ 



**Fig. 6:** When  $\min_{i \in I_{3k+1}} X_i^1 < \eta H_{\lambda}$  and  $\min_{i \in I_{3k+3}} X_i^1 < \eta H_{\lambda}$ , there is a uniformly (in  $\lambda$ ) positive probability that  $\max_{i \in I_{3k+2}} X_i^1 < 2\eta H_{\lambda}$  while taking values smaller than  $\eta H_{\lambda}$  at the boundary of  $I_{3k+2}$  (black dots). The white vertices correspond to the position of the path at times  $\ell_k^1$  and  $r_k^1$ .

Remark 4. A similar argument applies for the stationary measure  $\mathbb{P}_{\lambda}$ . This means that we may derive our target (3.29) for  $\mathbb{P}_{\lambda}(\cdot | X_{-N}, X_N \leq \eta H_{\lambda})$  instead of deriving it for  $\mathbb{P}_{\lambda}$  itself.

The claim (3.31) then follows by using the lower bound (3.3) on the partition function (and taking  $\eta$  large enough).

Let us say that the triple  $(I_{3k+1}, I_{3k+2}, I_{3k+3})$  is potentially  $\eta$ -good if (see Figure 6)

$$\min_{i \in I_{3k+1}} X_i^1 < \eta H_{\lambda}, \quad \min_{i \in I_{3k+1}} X_i^2 < \eta H_{\lambda}, \quad \min_{i \in I_{3k+3}} X_i^1 < \eta H_{\lambda}, \quad \min_{i \in I_{3k+3}} X_i^2 < \eta H_{\lambda}.$$

Let us denote by  $\tilde{M}$  the number of potentially  $\eta$ -good triples. We deduce from (3.31) that, for any  $\epsilon > 0$ , we can find  $\eta$  such that

$$\mathbb{P}_{N,+,\lambda}^{0,0} \otimes \mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}(\tilde{M} \le (1-\epsilon)\frac{m}{3}) \le e^{-c_4 N H_{\lambda}^{-2}},$$

for some  $c_4 > 0$ .

Now, given a potentially  $\eta$ -good triple  $(I_{3k+1}, I_{3k+2}, I_{3k+3})$ , let

$$\ell_k^1 = \min \left\{ i \in I_{3k+1} \, : \, X_i^1 < \eta H_{\lambda} \right\}, \, r_k^1 = \max \left\{ i \in I_{3k+3} \, : \, X_i^1 < \eta H_{\lambda} \right\}.$$

Conditionally on  $X^1_{\ell^1_k}$  and  $X^1_{r^1_k}$ , the probability that  $X^1_i \leq 2\eta H_\lambda$  for all  $\ell^1_k < i < r^1_k$  and that both walks sit below  $\eta H_\lambda$  at the end-points of  $I_{3k+2}$  is bounded away from zero, uniformly in  $\lambda$ . Indeed, uniformly in  $\mathbf{x}, \mathbf{y} < \eta H_\lambda, \ n \leq 3H_\lambda^2$  and  $1 \leq k < m \leq n$ ,

$$\begin{split} \mathbb{P}_{n,+,\lambda}^{\mathbf{x},\mathbf{y}} \left( \max_{i} X_{i}^{1} \leq 2\eta H_{\lambda}; X_{k}^{1} < \eta H_{\lambda}; X_{m}^{1} < \eta H_{\lambda} \right) \\ & \geq e^{-6q(2\eta)} \, \mathbb{P}_{n,+,0}^{\mathbf{x},\mathbf{y}} \left( \max_{i} X_{i}^{1} \leq 2\eta H_{\lambda}; X_{k}^{1} < \eta H_{\lambda}; X_{m}^{1} < \eta H_{\lambda} \right), \end{split}$$

since  $nV_{\lambda}(2\eta H_{\lambda}) \leq 6q(2\eta)$  (and  $\hat{Z}_{n,+,\lambda}^{\mathsf{x},\mathsf{y}} \leq 1$ ). That the latter probability is bounded below is a consequence of the invariance principle.

The claim of the lemma now follows easily, since, conditionally on the pieces of paths between  $r_{k-1}^1$  and  $\ell_k^1$ , these events are independent (and since the same argument can be made independently for  $X^2$ ).  $\square$ 

Now that we know that we can find  $O(NH_{\lambda}^{-2})$   $\eta$ -good intervals, the main observation is that, inside each such interval, there is a uniformly positive probability that the two paths meet. Let us make this more precise:

Definition. For  $\frac{n}{3} \leq m \leq n$ , let  $\mathcal{R}_{n,m}^{\mathsf{x},\mathsf{y}}[\eta]$  be the set of paths  $\mathbb{X} = (X_1,\ldots,X_m)$ satisfying  $X_1 = x$ ,  $X_m = y$  and  $0 < X_i < 2\eta\sqrt{n}, i = 1, ..., n$ . We shall employ the short-hand notation  $\mathcal{R}_n^{x,y}[\eta] = \mathcal{R}_{n,n}^{x,y}[\eta]$ ; the argument  $\eta$  will often be dropped when no ambiguity arises. We also set  $\mathcal{R}_{n,m}^{\mathsf{x},\mathsf{y},+} = \mathcal{R}_{n,m}^{\mathsf{x},\mathsf{y}}[\infty]$ .

**Proposition 6.** There exists  $\eta_0 > 0$  such that the following happens: For every  $\eta \geq \eta_0$ , one can find  $n_0 = n_0(\eta) \in \mathbb{N}$  and  $p = p(\eta) > 0$  such that

$$\mathbb{P}_{n,+,0}^{\mathsf{x},\mathsf{y}} \otimes \mathbb{P}_{n,+,0}^{\mathsf{z},\mathsf{w}} (\exists i \, : \, X_i^1 = X_i^2 \, | \, \mathcal{R}_n^{\mathsf{x},\mathsf{y}} \times \mathcal{R}_n^{\mathsf{z},\mathsf{w}}) \geq p, \tag{3.32}$$

uniformly in  $n \geq n_0$  and  $x, y, z, w \in (0, \eta \sqrt{n}] \cap \mathbb{N}$ .

Proposition 6 is a statement about random walks with transition probabilities p. We relegate the proof to the Appendix and proceed with the proof of Proposition 5.

First of all pick  $n=H_{\lambda}^2$  and note that, in view of Assumption (1.15), the following happens: For any  $x,y\leq \eta H_{\lambda}$  and any path  $\mathbb{X}\in\mathcal{R}_n^{\mathsf{x},\mathsf{y}}[\eta]$ , the value of the potential satisfies

$$0 \le \sum_{1}^{n} V_{\lambda}(X_i) \le nV_{\lambda}(2\eta\sqrt{n}) = H_{\lambda}^{2}V_{\lambda}(2\eta H_{\lambda}) \le 2q(2\eta), \tag{3.33}$$

for all  $\lambda$  sufficiently small. In fact, (3.33) was precisely the reason to introduce the notion of  $\eta$ -good intervals. An immediate consequence of (3.32) and (3.33) is that

$$\mathbb{P}^{\mathsf{x},\mathsf{y}}_{H^{2}_{\lambda},+,\lambda}\otimes\mathbb{P}^{\mathsf{z},\mathsf{w}}_{H^{2}_{\lambda},+,\lambda}(\exists i\,:\,X^{1}_{i}=X^{2}_{i}\,|\,\mathcal{R}^{\mathsf{x},\mathsf{y}}_{H^{2}_{\lambda}}\times\mathcal{R}^{\mathsf{z},\mathsf{w}}_{H^{2}_{\lambda}})\geq p\mathrm{e}^{-4q(2\eta)},\tag{3.34}$$

for all  $\lambda$  sufficiently small. The formula (3.34) provides a uniform lower bound

on probability of coupling inside an  $\eta$ -good interval. Consider the product measures  $\mathbb{P}^{0,0}_{N,+,\lambda_N}\otimes\mathbb{P}^{\mathsf{u},\mathsf{v}}_{N,+,\lambda}$ . Let  $\mathcal{M}$  be the event that the paths  $X^1$  and  $X^2$  meet both on the left and on the right of the segment [-T,T]. It follows from Lemma 3 and (3.34) that there exist K=K(C,T) and  $c_5 > 0$  such that

$$\mathbb{P}_{N,+,\lambda}^{0,0} \otimes \mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}(\mathcal{M}) \ge 1 - e^{-c_5 N H_{\lambda}^{-2}}, \tag{3.35}$$

uniformly in  $\lambda$  small,  $\mathsf{u},\mathsf{v} \leq CH_\lambda$  and  $N > (K+T)H_\lambda$ . For  $\ell < -TH_\lambda^2, r > TH_\lambda^2$  and  $\mathsf{x},\mathsf{y} \in \mathbb{N}$ , let  $\mathcal{M}_{\ell,r}^{\mathsf{x},\mathsf{y}} \subset \mathcal{M}$  be the event that  $\ell$ is the leftmost meeting point of  $X^1$ ,  $X^2$ , and  $X^1_\ell = X^2_\ell = \mathsf{x}$ , whereas r is the rightmost meeting point of  $X^1$ ,  $X^2$ , and  $X^1_r = X^2_r = \mathsf{y}$ . In this notation,  $\mathcal{M}$  is the disjoint union,  $\mathcal{M} = \cup \mathcal{M}^{\mathsf{x},\mathsf{y}}_{\ell,r}$ .

Let  $A \in \mathcal{F}_{\lambda,T}$ . Then,

$$\begin{split} \mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}(\mathcal{A}) &= \mathbb{P}_{N,+,\lambda}^{0,0} \otimes \mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}(\Omega \times \mathcal{A}) \\ &= \mathbb{P}_{N,+,\lambda}^{0,0} \otimes \mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}(\Omega \times \mathcal{A};\mathcal{M}^c) + \sum_{\substack{\ell,r\\\mathsf{x},\mathsf{y}}} \mathbb{P}_{N,+,\lambda}^{0,0} \otimes \mathbb{P}_{N,+,\lambda}^{\mathsf{u},\mathsf{v}}(\Omega \times \mathcal{A};\mathcal{M}^{\mathsf{x},\mathsf{y}}_{\ell,r}). \end{split}$$

However,

$$\mathbb{P}^{0,0}_{N,+,\lambda}\otimes\mathbb{P}^{\mathsf{u},\mathsf{v}}_{N,+,\lambda}(\varOmega\times\mathcal{A};\mathcal{M}^{\mathsf{x},\mathsf{y}}_{\ell,r})=\mathbb{P}^{0,0}_{N,+,\lambda}\otimes\mathbb{P}^{\mathsf{u},\mathsf{v}}_{N,+,\lambda}(\mathcal{A}\times\varOmega;\mathcal{M}^{\mathsf{x},\mathsf{y}}_{\ell,r}).$$

Therefore,

$$\left|\mathbb{P}^{\mathsf{u},\mathsf{v}}_{N,+,\lambda}(\mathcal{A}) - \mathbb{P}^{0,0}_{N,+,\lambda}(\mathcal{A})\right| \leq \mathbb{P}^{0,0}_{N,+,\lambda} \otimes \mathbb{P}^{\mathsf{u},\mathsf{v}}_{N,+,\lambda}(\mathcal{M}^c),$$

which, in view of Remark 4 and (3.35), implies (3.29).  $\square$ 

## A. Proof of Proposition 6

We shall employ here our original notation p for the path measure of the underlying random walk. Our argument is based on the second moment method, which is put to work using the following input from [6,9]:

Bounds on probabilities of random walks to stay positive. There exists  $\eta_0$ , such that for any  $\eta \geq \eta_0$  the following happens: One can find  $n_0 = n_0(\eta)$ ,  $c_1 = c_1(\eta)$  and  $c_2 = c_2(\eta)$  such that

$$c_1 \frac{\mathsf{x}\mathsf{y}}{n^{3/2}} \le \mathsf{p}(\mathcal{R}_{n,m}^{\mathsf{x},\mathsf{y}}\left[\eta\right]) \le c_2 \frac{\mathsf{x}\mathsf{y}}{n^{3/2}},\tag{A.1}$$

uniformly  $n \geq n_0$ ,  $\frac{n}{3} \leq m \leq n$  and  $1 \leq \mathsf{x}, \mathsf{y} \leq \eta \sqrt{n}$ . Indeed, in view of the invariance principle for random walk bridges [9, Theorem 2.4], the restriction  $X_i \leq 2\eta \sqrt{n}$  may be removed from  $\mathcal{R}_{n,m}^{\mathsf{x},\mathsf{y}}[\eta]$  in the following sense: There exists  $c = c(\eta) \in [1, \infty)$ , such that

$$1 \le \frac{\mathsf{p}(\mathcal{R}_{n,m}^{\mathsf{x},\mathsf{y},+})}{\mathsf{p}(\mathcal{R}_{n,m}^{\mathsf{x},\mathsf{y}}[\eta])} \le c(\eta),\tag{A.2}$$

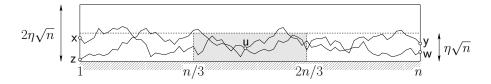
uniformly in  $n \geq n_0$ ,  $\frac{n}{3} \leq m \leq n$  and  $x, y \leq \eta \sqrt{n}$ . Hence, the two-sided inequality (A.1) can be verified along the lines of the proof of Theorem 4.3 in [6], where a stronger asymptotic statement was derived for a more restricted range of parameters.

Let  $1 \le x, y, z, w \le \eta \sqrt{n}$  and consider now the product measure,

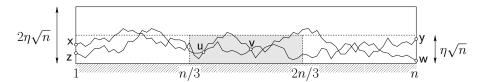
$$\mathbf{p} \otimes \mathbf{p} \left( \cdot \, | \, \mathcal{R}_{n}^{\mathsf{x},\mathsf{y}} \left[ \boldsymbol{\eta} \right] \times \mathcal{R}_{n}^{\mathsf{z},\mathsf{w}} \left[ \boldsymbol{\eta} \right] \right).$$

Let  $\mathcal{N}$  be the number of intersections of the two replicas  $X^1$  and  $X^2$  during the time interval  $\left[\frac{n}{3}, \frac{2n}{3}\right]$  below level  $\eta\sqrt{n}$ . Precisely,

$$\mathcal{N} = \# \left\{ \ell \in \left[ \frac{n}{3}, \frac{2n}{3} \right] : X_{\ell}^{1} = X_{\ell}^{2} \le \eta \sqrt{n} \right\} = \sum_{\ell = \frac{n}{3}}^{\frac{2n}{3}} \sum_{\mathbf{u} < \eta \sqrt{n}} \mathbf{1}_{\left\{ X_{\ell}^{1} = X_{\ell}^{2} = \mathbf{u} \right\}}.$$
 (A.3)



**Fig. 7:** The decomposition of  $p \otimes p(\mathcal{N} \mid \mathcal{R}_n^{x,y} \times \mathcal{R}_n^{z,w})$ .



**Fig. 8:** The decomposition of  $p \otimes p(\mathcal{N}^2 \mid \mathcal{R}_n^{x,y} \times \mathcal{R}_n^{z,w})$ .

Lower bound on the expectation  $p \otimes p(\mathcal{N} \mid \mathcal{R}_n^{x,y} \times \mathcal{R}_n^{z,w})$ . The expectation

$$\mathsf{p} \otimes \mathsf{p}(\mathcal{N} \,|\, \mathcal{R}_n^{\mathsf{x},\mathsf{y}} \times \mathcal{R}_n^{\mathsf{z},\mathsf{w}}) = \sum_{\ell = \frac{n}{2}}^{\frac{2n}{3}} \sum_{\mathsf{u} < \eta \sqrt{n}} \varPhi_n(\ell,\mathsf{u};\mathsf{x},\mathsf{y},\mathsf{z},\mathsf{w}),$$

where (see Figure 7)

$$\varPhi_n(\ell,\mathbf{u};\mathbf{x},\mathbf{y},\mathbf{z},\mathbf{w}) = \frac{\mathbf{p}(\mathcal{R}_{n,\ell}^{\mathbf{x},\mathbf{u}})\mathbf{p}(\mathcal{R}_{n,\ell}^{\mathbf{z},\mathbf{u}})\mathbf{p}(\mathcal{R}_{n,n-\ell}^{\mathbf{u},\mathbf{y}})\mathbf{p}(\mathcal{R}_{n,n-\ell}^{\mathbf{u},\mathbf{w}})}{\mathbf{p}(\mathcal{R}_{n,\ell}^{\mathbf{x},\mathbf{y}})\mathbf{p}(\mathcal{R}_{n,\ell}^{\mathbf{z},\mathbf{w}})}.$$

By (A.1),

$$\Phi_n(\ell,\mathsf{u};\mathsf{x},\mathsf{y},\mathsf{z},\mathsf{w}) \geq c_3 \frac{\mathsf{u}^4}{n^3},$$

uniformly in all the arguments in question. Consequently,

$$p \otimes p(\mathcal{N} \mid \mathcal{R}_n^{x,y} \times \mathcal{R}_n^{z,w}) \ge c_4(\eta)\sqrt{n},$$
 (A.4)

also uniformly in  $x, \ldots, w \le \eta \sqrt{n}$ .

Upper bound on the expectation  $p \otimes p(\mathcal{N}^2 \mid \mathcal{R}_n^{x,y} \times \mathcal{R}_n^{z,w})$ . The expectation

$$\mathsf{p} \otimes \mathsf{p}(\mathcal{N}^2 \,|\, \mathcal{R}_n^{\mathsf{x},\mathsf{y}} \times \mathcal{R}_n^{\mathsf{z},\mathsf{w}}) \leq \sum_{\substack{\ell, m = \frac{n}{3} \\ \ell < m}}^{\frac{2n}{3}} \sum_{\mathsf{u},\mathsf{v} \leq \eta \sqrt{n}} \varPsi_n(\ell,\mathsf{u};m,\mathsf{v};\mathsf{x},\mathsf{y},\mathsf{z},\mathsf{w}),$$

where (see Figure 8)

$$\varPsi_n(\ell,\mathbf{u};m,\mathbf{v};\mathbf{x},\mathbf{y},\mathbf{z},\mathbf{w}) = \frac{\mathsf{p}(\mathcal{R}_{n,\ell}^{\mathbf{x},\mathbf{u}})\mathsf{p}(\mathcal{R}_{n,\ell}^{\mathbf{z},\mathbf{u}})\mathsf{p}_{m-\ell}(\mathbf{u},\mathbf{v})^2\mathsf{p}(\mathcal{R}_{n,n-m}^{\mathbf{v},\mathbf{y}})\mathsf{p}(\mathcal{R}_{n,n-m}^{\mathbf{v},\mathbf{w}})}{\mathsf{p}(\mathcal{R}_{n,\ell}^{\mathbf{x},\mathbf{y}})\mathsf{p}(\mathcal{R}_{n,\ell}^{\mathbf{z},\mathbf{w}})}. \tag{A.5}$$

Above,  $p_r(u, v)$  is a short-hand notation for  $p(X_r = v | X_0 = u)$ . The inequality in (A.5) is due to the fact that we ignore the positivity condition on the interval  $\ell, \ldots, m$ .

By (A.1),

$$\varPsi_n(\ell, \mathbf{u}; m, \mathbf{v}; \mathbf{x}, \mathbf{y}, \mathbf{z}, \mathbf{w}) \leq c_6(\eta) \frac{\mathsf{p}_{m-\ell}(\mathbf{u}, \mathbf{v})^2}{n},$$

uniformly in all the arguments in question. Consequently,

$$\sum_{\substack{\ell,m=\frac{n}{3}\\\ell\leq m}}^{\frac{2n}{3}}\sum_{\mathbf{u},\mathbf{v}\leq \eta\sqrt{n}}\varPsi_n(\ell,\mathbf{u};m,\mathbf{v};\mathbf{x},\mathbf{y},\mathbf{z},\mathbf{w})\leq c_7\sqrt{n}\sum_{r=0}^{\frac{n}{3}}\sum_{\mathbf{v}}\mathsf{p}_r(0,\mathbf{v})^2.$$

The double sum on the right-hand side above is equal to the expectation of the number of intersections of two independent p-walks during the first  $\frac{n}{3}$  steps of their life. It is bounded above by  $c_8\sqrt{n}$ . We conclude that

$$p \otimes p(\mathcal{N}^2 \mid \mathcal{R}_n^{\mathsf{x},\mathsf{y}} \times \mathcal{R}_n^{\mathsf{z},\mathsf{w}}) \le c_9 n. \tag{A.6}$$

The lower and upper bounds (A.4) and (A.6) imply the existence of  $\nu = \nu(\eta) > 0$ , such that

$$p \otimes p(\mathcal{N}^2 \mid \mathcal{R}_n^{x,y} \times \mathcal{R}_n^{z,w}) \le \nu[p \otimes p(\mathcal{N} \mid \mathcal{R}_n^{x,y} \times \mathcal{R}_n^{z,w})]^2, \tag{A.7}$$

uniformly in  $x, y, z, w \leq \eta \sqrt{n}$ . Set  $\mathcal{E} = p \otimes p(\mathcal{N} \mid \mathcal{R}_n^{x,y} \times \mathcal{R}_n^{z,w})$ . By the Paley-Zygmund inequality,

$$p \otimes p(\mathcal{N} > \alpha \mathcal{E} \mid \mathcal{R}_n^{x,y} \times \mathcal{R}_n^{z,w}) \ge \frac{(1-\alpha)^2}{\nu},$$
 (A.8)

for every  $\alpha \leq 1$ . (3.32) follows.

## References

- D. B. Abraham and E. R. Smith. An exactly solved model with a wetting transition. J. Statist. Phys., 43(3-4):621–643, 1986.
- Robert A. Adams and John J. F. Fournier. Sobolev spaces, volume 140 of Pure and Applied Mathematics (Amsterdam). Elsevier/Academic Press, Amsterdam, second edition, 2003.
- 3. L. Alili and R. A. Doney. Wiener-Hopf factorization revisited and some applications. Stochastics Stochastics Rep., 66(1-2):87–102, 1999.
- 4. J. Bertoin and R. A. Doney. On conditioning a random walk to stay nonnegative. *Ann. Probab.*, 22(4):2152–2167, 1994.
- T. Bodineau, V. Lecomte, and C. Toninelli. Finite size scaling of the dynamical free-energy in a kinetically constrained model. J. Stat. Phys., 147(1):1–17, 2012.
- M. Campanino, D. Ioffe, and O. Louidor. Finite connections for supercritical Bernoulli bond percolation in 2D. Markov Process. Related Fields, 16(2):225–266, 2010.
   P. Caputo, E. Lubetzky, F. Martinelli, A. Sly, and F. L. Toninelli. The shape of the
- P. Caputo, E. Lubetzky, F. Martinelli, A. Sly, and F. L. Toninelli. The shape of the (2+1)D SOS surface above a wall. C. R. Math. Acad. Sci. Paris, 350(13-14):703-706, 2012.
- 8. F. Caravenna and L. Chaumont. Invariance principles for random walks conditioned to stay positive. *Ann. Inst. Henri Poincaré Probab. Stat.*, 44(1):170–190, 2008.
- 9. F. Caravenna and L. Chaumont. An invariance principle for random walk bridges conditioned to stay positive. *Electron. J. Probab.*, 18:no. 60, 32, 2013.
- E. A. Coddington and N. Levinson. Theory of ordinary differential equations. McGraw-Hill Book Company, Inc., New York-Toronto-London, 1955.

- 11. S. N. Ethier and T. G. Kurtz. *Markov processes*. Wiley Series in Probability and Mathematical Statistics: Probability and Mathematical Statistics. John Wiley & Sons Inc., New York, 1986. Characterization and convergence.
- 12. P. L. Ferrari and H. Spohn. Constrained Brownian motion: fluctuations away from circular and parabolic barriers. *Ann. Probab.*, 33(4):1302–1325, 2005.
- 13. M. E. Fisher. Walks, walls, wetting, and melting. J. Statist. Phys., 34(5-6):667–729, 1984.
- 14. O. Hryniv and Y. Velenik. Universality of critical behaviour in a class of recurrent random walks. *Probab. Theory Related Fields*, 130(2):222–258, 2004.
- 15. M. G. Kreĭn and M. A. Rutman. Linear operators leaving invariant a cone in a Banach space. *Amer. Math. Soc. Translation*, 1950(26):128, 1950.
- 16. R. H. Schonmann and S. B. Shlosman. Constrained variational problem with applications to the Ising model. *J. Statist. Phys.*, 83(5-6):867–905, 1996.
- 17. W. Selke, D. A. Huse, and D. M. Kroll. Interfacial adsorption in the two-dimensional Blume-Capel model. *Journal of Physics A: Mathematical and General*, 17(15):3019, 1984.
- 18. Y. Velenik. Entropic repulsion of an interface in an external field. *Probab. Theory Related Fields*, 129(1):83–112, 2004.