

Merged Q1 Coarse Spaces for Schwarz Methods in 2D and 3D

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1 Introduction

Coarse spaces are a key ingredient in the development of scalable domain decomposition methods. Our goal here is to develop new 2D and 3D low-dimensional but robust coarse spaces based on the observation of the eigenmodes of the considered operator, namely here the RAS [4] iteration operator. A first step of this development was presented in [8] with the introduction, for the 2-D case, of the so-called `Half_Q1` coarse space, which will here be re-named as `Merged_2` coarse space in view of its 3D generalization. Our merged coarse spaces are obtained by taking a subset of the coarse space functions of the Q1 coarse space [6, 7], which is based, in 2D, on the Q1 bilinear finite element functions on rectangular elements, these elements being here the subdomains. This leads, in 2D, to four coarse grid points around each cross point of the non-overlapping decomposition (Fig. 2(a)). In 3D, the equivalent Q1 coarse space has eight coarse grid points around each cross point (Fig. 2(b)). The merged coarse spaces we consider will be described in section 2 for the 2D case and 3 for the 3D case. Our coarse spaces are implemented using PETSc [1, 2, 3] which, besides benefiting from its parallel features, enables us to compare our optimized two-level domain decomposition results, i.e., using the optimized RAS (ORAS [9, 12]) method, with multigrid options (GAMG, HYPRE). Our model problem is the Laplace problem $-\Delta u = 0$ or, adding advection with an upwind scheme, the non-symmetric problem

$$-\Delta u + \mathbf{a} \cdot \nabla u = 0 \tag{1}$$

using the 5-point finite difference scheme and homogeneous boundary conditions.

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2 Merged coarse spaces in 2D

In Fig. 1 we show the four largest eigenmodes (i.e., with eigenvalues closest to one in modulus) of the RAS iteration operator (i.e., $I - BA$ with A the stiffness matrix and B the RAS preconditioner) for the 2D Laplace model problem using a 2×2 algebraically non-overlapping decomposition (RAS then reduces to Block Jacobi), for a global 256×256 fine mesh resolution. The first pair of eigenmodes exhibits eigenvalues equal in modulus and of opposite signs, and similarly for the second pair. Moreover, the first two modes appear to be made of piecewise Q1 functions (as was first shown in [5]): denoting by q_1, q_2, q_3, q_4 the Q1 basis functions at a cross point (with the numbering convention of Fig. 2(a)), i.e., bilinear with value 1 at the cross point corner and 0 at the other corners of the subdomain (see the squared points in Fig. 2(c)), the first mode (continuous, with positive eigenvalue) appears to be $q_1 + q_2 + q_3 + q_4$, while the second mode (discontinuous, with negative eigenvalue) appears to be $q_1 - q_2 - q_3 + q_4$. Therefore the idea of the `Merged_2` coarse space (already introduced in [8] under the name `Half_Q1`) is to merge the four basis

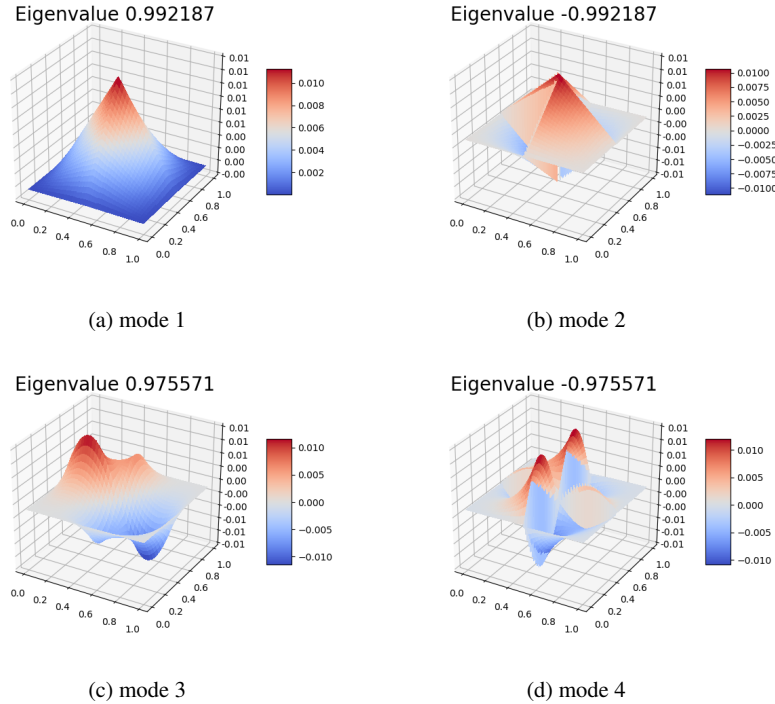


Fig. 1: Largest four modes of the RAS iteration operator for the Laplace problem using a 2×2 algebraically non-overlapping domain decomposition.

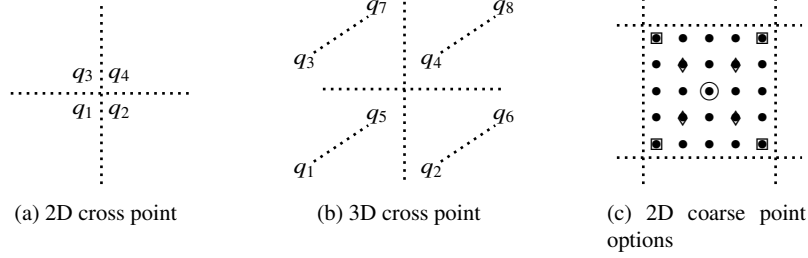


Fig. 2: (a) and (b): Numbering conventions of the Q1 basis functions at a 2D and 3D cross point. (c): Coarse grid point choice in 2-D for the Q1 (squares), Q1_fair (diamonds) and Middle (circles) options.

functions of the Q1 coarse space into these two combinations, yielding a coarse space which is half the size of the original Q1 one. For the two largest eigenmodes, this reduction of the coarse space size comes at no cost. Note that the **Merged_2** coarse space can equivalently be defined by taking the combinations $q_1 + q_4$ and $q_2 + q_3$ at each cross point.

We can further pursue the merging idea for the case where RAS is used with GMRES acceleration. Discontinuous eigenmodes of the RAS iteration operator $I - BA$ have negative eigenvalues close to -1 , see e.g. Fig. 1, which implies that in the preconditioned system BA they are close to 2 and thus can easily be handled by GMRES. This is in contrast to the continuous eigenmodes which have eigenvalues close to 1, and thus in the preconditioned system they are close to 0, which is hard for GMRES, and thus these modes must be handled by another mechanism, which here is the coarse space. Thus, we define the **Merged_1** coarse space as the one where only the $q_1 + q_2 + q_3 + q_4$ combination of the Q1 basis functions (i.e., the largest continuous mode) is kept, yielding a coarse space which is four times smaller than Q1 and half the size of **Merged_2**. Table 1(a) shows the number of iterations necessary to reach the zero solution of the model Laplace problem, when successively taking as initial guess the largest four modes of the RAS iteration operator displayed in Fig. 1, and using stationary iterations with or without coarse correction, or using GMRES acceleration. Using stationary iterations, the Q1 coarse space has the first two modes converge in one iteration. By construction, the **Merged_2** coarse space also contains the first two modes and has them converge in one iteration, while the **Merged_1** coarse space contains only the largest continuous mode. The last line of Table 1(a) shows that the GMRES acceleration, already without coarse correction, yields convergence in one iteration for the largest (and second largest) discontinuous mode. The **Merged_1** coarse space is thus meant to be used with GMRES acceleration, avoiding an unnecessary duplication of effects, at least for the first two modes.

Mode number	1	2	3	4
Eigenvalue	0.99219	-0.99219	0.97557	-0.97557
<hr/>				
Stationary iterations				
no coarse correction	2349	2349	745	745
Q1	1	1	551	417
Merged_2	1	1	745	745
Merged_1	1	2349	745	745
<u>GMRES accelerated</u>				
no coarse correction	17	1	23	1
<hr/>				
(a) 2D case (256^2 fine mesh resolution)				
Eigenvalue	0.96875	-0.96875	0.93035	-0.93035
<hr/>				
Stationary iterations				
no coarse correction	581	581	256	256
Q1	1	2	175	141
Merged_2	1	2	256	256
Merged_1	1	581	256	256
<u>GMRES accelerated</u>				
no coarse correction	8	1	7	1
<hr/>				
(b) 3D case (64^3 fine mesh resolution)				

Table 1: Number of iterations to reach the zero solution of the Laplace model problem with the largest four RAS iteration operator eigenmodes taken as initial guess.

3 Merged coarse spaces in 3D

Our numbering convention of the Q1 basis functions at a 3D cross point is given in Fig. 2(b). We define our 3D merged coarse spaces using the same symmetry as the one used in 2D. That is, the **Merged_2** coarse space will use the following combinations as basis functions: $\sum_i q_i$ and $q_1 - q_2 - q_3 + q_4 - q_5 + q_6 + q_7 - q_8$ or, equivalently, $q_1 + q_4 + q_6 + q_7$ and $q_2 + q_3 + q_5 + q_8$. As for the **Merged_1** coarse space, it will use the sum $\sum_i q_i$ as unique basis function.

We consider again the four largest eigenmodes of the RAS iteration operator for the Laplace model problem, this time using a $2 \times 2 \times 2$ algebraically non-overlapping decomposition. These modes can not be plotted as in 2D but we have that, as in 2D, the first pair of eigenmodes exhibits eigenvalues equal in modulus and of opposite signs, and similarly for the second pair: the actual values are given on the second line of Table 1(b). Moreover, this table shows that the Q1 and **Merged_2** coarse spaces contain the first two modes, while the **Merged_1** coarse space contains only the first (positive) mode¹. This justifies the definition of our 3D merged coarse spaces. Moreover, here again the GMRES acceleration alone yields convergence in one iteration for

¹ The number of iterations for the mode number 2 with Q1 and **Merged_2** coarse spaces is 2 and not 1, but this is due only to its high sensitivity to the convergence criteria: taking a convergence criteria of $1.05\text{E-}8$ instead of $1\text{E-}8$ would have made the convergence achieved at iteration 1.

the largest (and second largest) negative mode, making the `Merged_1` coarse space a method of choice when this acceleration is used.

4 Numerical Results

We present numerical results of a weak scaling experiment for our model problem, using square domain decompositions (2×2 , 4×4 , ... and similarly in 3D). Each run starts from a random initial guess and converges to the zero solution with a relative convergence tolerance of $1.e-8$. Our results were obtained on the CPU partition of the *Adastra* machine hosted at CINES in Southern France (www.cines.fr), using one CPU per subdomain. We compare our merged coarse spaces to other coarse spaces, namely the *Nicolaides* coarse space [11] which uses a constant basis function on each subdomain, as well as, in 2D, to other coarse spaces using bilinear functions centered on coarse points depicted in Fig. 2(c): while the Q1 coarse space has coarse points in each corner of the subdomain, the `Q1_fair` coarse space has as many coarse points as the Q1 coarse space but equally distributed in space, and the `Middle` coarse space takes the fine mesh point in the middle of each subdomain as coarse grid point. Note that these coarse spaces were already considered in [8]. For an $N \times N$ domain decomposition, the asymptotic size of the coarse space is thus $4N^2$ for Q1 and `Q1_fair`, $2N^2$ for `Merged_2`, and N^2 for `Merged_1`, `Middle` and *Nicolaides*.

In Fig. 3(a) we show the number of iterations needed for convergence when RAS is used without GMRES acceleration on the algebraically non-overlapping Laplace problem. As expected, the `Merged_1` coarse space yields very poor results, while the `Merged_2` coarse space yields results which are not too far away from the Q1 ones, notably better than the `Q1_fair` methods which uses a twice bigger coarse space. With GMRES (Fig. 3(b)), the `Merged_1` results come very close to the `Merged_2` ones (they are in fact undistinguishable on the plot), as expected. Note that `Q1_fair` yields here the best results: it appears to take more advantage of the GMRES acceleration than the Q1 method and its merged variants.

In Figs. 3(c) and 3(d) we show the number of iterations and solution times for the model problem (1) with rotating advection $a_x = -10y$, $a_y = 10x$, and an algebraic overlap of 2, using the optimized RAS (ORAS) method based on optimized transmission conditions at subdomain interfaces [9, 12]. The ORAS method is implemented with a first-order accurate discretization of the normal derivative, which implies modifying only the diagonal components of the local system matrix, and the Robin parameters are the ones derived for the (symmetric) Laplace problem [6]. The number of iterations of the Q1, `Merged_2` and `Merged_1` options are close in Fig. 3(c), here with Q1 and `Merged_2` undistinguishable and `Merged_1` slightly above. Note that the number of iterations with 128^2 ($= 16,384$) subdomains (not on the plot) are the same as the ones shown here at 64^2 ($= 4,096$) subdomains.

In Fig. 3(d) we show solution times corresponding to Fig. 3(c), as well as coarse correction times per iteration. As expected, the smaller size of the merged coarse

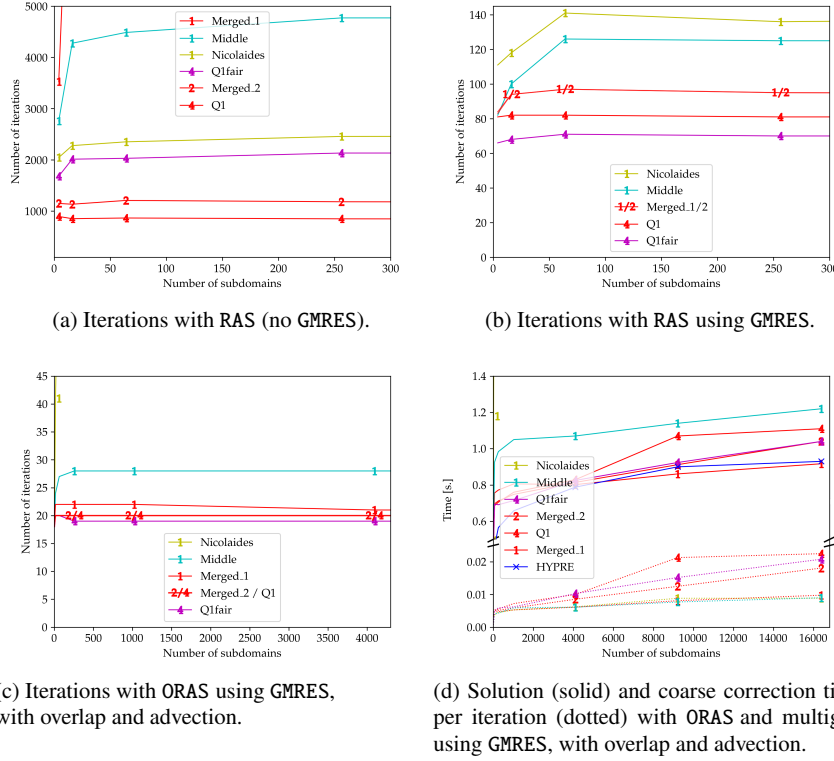
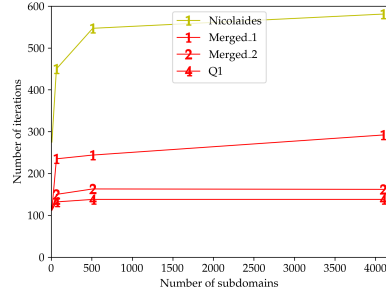


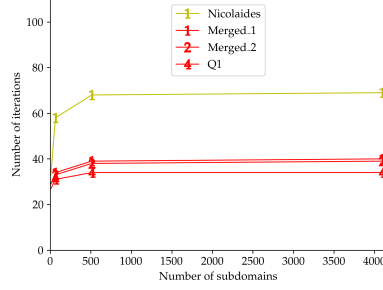
Fig. 3: Number of iterations (a)(b)(c) and solution times (d) for the model problem (1) in 2D using a 256×256 local mesh. In (a) and (b) $a_x = a_y = 0$ and the algebraic overlap is 0. In (c) and (d) $a_x = -10y$, $a_y = 10x$, and the algebraic overlap is 2.

spaces reduces the coarse correction time per iteration compared to the original Q1 coarse space (Merged_1, Middle and Nicolaides are nearly undistinguishable, in accordance to their same asymptotic size) and, since the merging process comes at a very moderate price in terms of iterations, it results in a total time to solution gain. Moreover, our Merged_1 results appear comparable to the ones obtained with the algebraic multigrid HYPRE/BoomerAMG option [10] available through PETSc (with optimized tuning form [13]). However, if one includes setup costs, the multigrid option keeps its advantage: for instance at 128^2 cores, the Merged_1 setup cost (which notably includes the local LU factorizations) is .93s, while it is only .03s for the HYPRE option.

In Figs. 4(a) and 4(b) we show the number of iterations for the Laplace problem in 3D, with and without GMRES acceleration, using an algebraic overlap of 2 and a 64^3 local mesh. As expected, the Merged_1 coarse space performs poorly (even if not as bad as in 2D) without GMRES, but yields results very similar to Merged_2 with GMRES.



(a) Iterations with RAS (no GMRES).



(b) Iterations with RAS using GMRES.

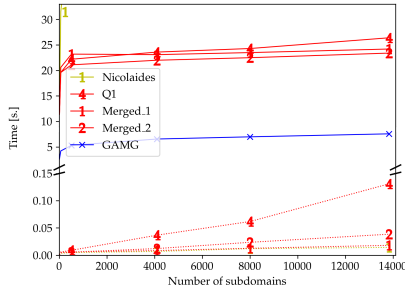
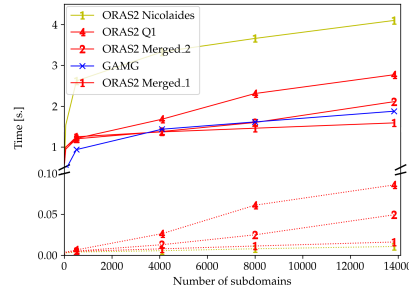
(c) Solution time with ORAS and multigrid, using GMRES (64^3 local mesh).(d) Solution time with ORAS and multigrid, using GMRES (32^3 local mesh).

Fig. 4: Number of iterations (a)(b) and solution times (c)(d) for the Laplace problem in 3D using a 64^3 (a)(b)(c) or 32^3 (d) local mesh and an algebraic overlap of 2.

Moreover, the GMRES Merged_1 and Merged_2 results are close to the Q1 result, showing again that the decrease in coarse space size resulting from the merging process comes at a low cost in terms of iterations. In Fig. 4(c) we show time to solution results for the same Laplace problem but with ORAS, as well as the corresponding coarse correction times per iteration. As in 2D and as expected, the Merged_1 coarse space has its coarse correction time very close to the Nicolaides one and smaller than Merged_2, itself smaller than Q1. In turn, the merged coarse spaces tend to perform better than the Q1 one even if, at 24^3 ($= 13,824$) cores, the Merged_2 coarse space remains a bit better than the Merged_1. The solution time with PETSc's native algebraic multigrid preconditioner GAMG (with smoothed aggregation and CG eigenvalue estimator [2]) is also given, with here a clearly better performance. However, the algebraic multigrid solution time performance (outside of setup costs²) is approached more closely by our merged coarse space if one considers a smaller local mesh, for instance 32^3 , as shown in Fig. 4(d) (with the expected hierarchy:

² At 24^3 cores: 1.26s. (ORAS with Merged_1) vs. 0.02s. (GAMG) with the 32^3 local mesh, 23.19s. vs. 0.03s. with the 64^3 local mesh.

Merged_1 faster than Merged_2, itself faster than Q1). The number of subdomains, thus the number of cross points and, in turn, our merged coarse space size, is then relatively larger with respect to the global number of fine mesh points. This suggests that the efficiency of the Merged_Q1 coarse space compared to multigrid options could be improved by using more than one subdomain per CPU core.

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