## Self-avoiding walk on the hexagonal lattice

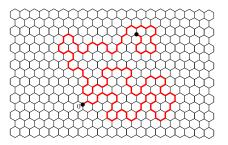
Hugo Duminil-Copin, Université de Genève

May 2010

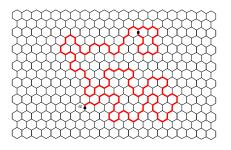
joint work with S. Smirnov



Consider the **hexagonal lattice**  $\mathbb{H}$ , and define  $c_n$  to be the number of self-avoiding walks of length n starting at the origin.



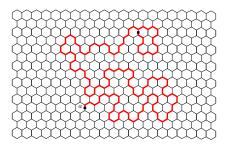
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• Combinatorial question: How many such trajectories are there? In other words, what is the value of  $c_n$ ?

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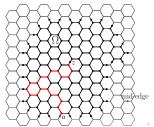
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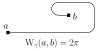
$$Z(x) := \sum_{\gamma : a \to \mathbb{H}} x^{-\ell(\gamma)}$$

We restrict ourself to *finite domains*  $\Omega$  and we weight walks by their winding.

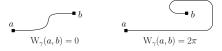


The winding  $W_{\Gamma}(a, b)$  of a curve  $\Gamma$  between a and b is the rotation (in radians) of the curve between a and b.





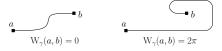
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With this definition, we can define the *parafermionic operator* for  $a \in \partial \Omega$  and  $z \in \Omega$ :

$$F(z) = F(a, z, x, \sigma) := \sum_{\gamma \subset \Omega: \ a \to z} e^{-i\sigma W_{\gamma}(a, z)} x^{-\ell(\gamma)}.$$

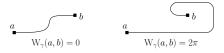
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### Lemma (Discrete integrals on elementary contours vanish)

If  $x = x_c$  and  $\sigma = \frac{5}{8}$ , then F satisfies the following relation for every vertex  $v \in V(\Omega)$ ,

$$(p-v)F(p) + (q-v)F(q) + (r-v)F(r) = 0$$

where p, q, r are the mid-edges of the three edges adjacent to v.

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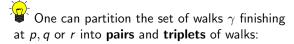
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In the first case:

$$\begin{split} c(\gamma_1) + c(\gamma_2) &= (q - v)e^{-i\sigma W_{\gamma_1}(a,q)}x_c^{-\ell(\gamma_1)} + (r - v)e^{-i\sigma W_{\gamma_2}(a,r)}x_c^{-\ell(\gamma_2)} \\ &= (p - v)e^{-i\sigma W_{\gamma_1}(a,p)}x_c^{-\ell(\gamma_1)}\left(e^{i\frac{2\pi}{3}}e^{-i\frac{5}{8}\cdot\frac{-4\pi}{3}} + e^{-i\frac{2\pi}{3}}e^{-i\frac{5}{8}\cdot\frac{4\pi}{3}}\right) = 0 \end{split}$$

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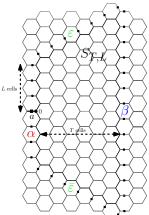
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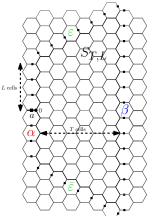
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In the second case:

$$c(\gamma_1) + c(\gamma_2) + c(\gamma_3)$$

$$= (p - \nu)e^{-i\sigma W_{\gamma_1}(a,p)} x_c^{-\ell(\gamma_1)} \left( 1 + x_c e^{i\frac{2\pi}{3}} e^{-i\frac{5}{8} \cdot \frac{-\pi}{3}} + x_c e^{-i\frac{2\pi}{3}} e^{-i\frac{5}{8} \cdot \frac{\pi}{3}} \right) = 0.$$





Sum the relation over all vertices in  $V(S_{T,L})$  (contour=boundary):

$$0 = -\sum_{z \in \alpha} F(z) + \sum_{z \in \beta} F(z) + e^{i\frac{2\pi}{3}} \sum_{z \in \varepsilon} F(z) + e^{-i\frac{2\pi}{3}} \sum_{z \in \bar{\varepsilon}} F(z)$$

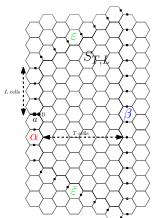


#### **Definition:**

$$A_{T,L}^{x} = \sum_{\gamma \subset S_{T,L}: a \to \alpha \setminus \{a\}} x^{-\ell(\gamma)},$$

$$B_{T,L}^{x} = \sum_{\gamma \subset S_{T,L}: a \to \beta} x^{-\ell(\gamma)},$$

$$E_{T,L}^{x} = \sum_{\gamma \subset S_{T,L}: a \to \epsilon \cup \bar{\epsilon}} x^{-\ell(\gamma)}.$$



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When  $x = x_c$ , we have

$$1 = \cos\left(\frac{3\pi}{8}\right) A_{T,L}^{x_c} + B_{T,L}^{x_c} + \cos\left(\frac{\pi}{4}\right) E_{T,L}^{x_c}$$

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$$0 = 1 - 1 = \cos\left(\frac{3\pi}{8}\right) \left(A_{T+1}^{x_c} - A_T^{x_c}\right) + B_{T+1}^{x_c} - B_T^{x_c}$$

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(3) From a walk entering in  $A_{T+1}^{x_c} - A_T^{x_c}$ , one can construct two excursions of width T+1, yielding  $A_{T+1}^{x_c} - A_T^{x_c} \le \left(B_{T+1}^{x_c}\right)^2$ . Thus,

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(4) It is an easy step to deduce that  $B_T^{x_c} \geq c/T$  and therefore

$$Z(x_c) \geq \sum_{r=1}^{\infty} B_T^{x_c} = +\infty.$$



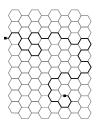
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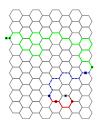
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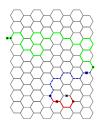
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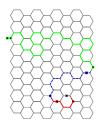
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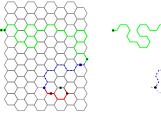
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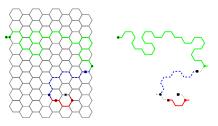
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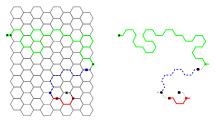
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(3) We obtain 
$$Z(x) \le 2 \sum_{\substack{T_{-i} < \dots < T_{-1} \\ T_{i} < \dots < T_{k}}} \left( \prod_{k=-i}^{j} B_{T_{k}}^{x} \right) = \prod_{T>0} (1 + B_{T}^{x})^{2} < +\infty$$

## Conjecture (Nienhuis)

There exists a constant  $A \in (0, +\infty)$  such that:

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## Conjecture (Lawler, Schramm, Werner)

The scaling limit of the self-avoiding walks is SLE(8/3).

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- similar observable for O(n) models: for  $n \in [-2,2]$ , infinite correlation length at  $\sqrt{2 + \sqrt{2 n}}$ .
- similar observable for random-cluster models: second (resp. first) order phase transition for  $q \in [1,4)$  (resp.  $q \in (4,+\infty)$ ), complete conformal invariance when q=2.

# Thank you

