

# High-order integrators for sampling the invariant distribution of stochastic ordinary and partial differential equations.

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based on joint works with

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## Long time accuracy for ergodic SDEs

$$dX(t) = f(X(t))dt + g(X(t))dW(t), \quad X(0) = x.$$

Under standard **ergodicity assumptions**,

$$\begin{aligned} \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \phi(X(t)) dt &= \int_{\mathbb{R}^d} \phi(y) d\mu(y) \\ \left| \mathbb{E}(\phi(X(t))) - \int_{\mathbb{R}^d} \phi(y) d\mu(y) \right| &\leq K(x, \phi) e^{-ct}, \quad \text{for all } t \geq 0. \end{aligned}$$

Two standard approaches using an ergodic integrator of **order  $p$** :

- Compute a single long trajectory  $\{X_n\}$  of length  $T = Nh$ ,

$$\frac{1}{N+1} \sum_{k=0}^N \phi(X_k) \simeq \int_{\mathbb{R}^d} \phi(y) d\mu(y), \quad \text{error } \mathcal{O}(h^p + T^{-1/2}),$$

- Compute many trajectories  $\{X_n^i\}$  of length of length  $t = Nh$ ,

$$\frac{1}{M} \sum_{i=1}^M \phi(X_N^i) \simeq \int_{\mathbb{R}^d} \phi(y) d\mu(y), \quad \text{error } \mathcal{O}(e^{-ct} + h^p + M^{-1/2}).$$

# A classical tool: the Fokker-Plank equation

$$dX(t) = f(X(t))dt + \sqrt{2}dW(t).$$

The density  $\rho(x, t)$  of  $X(t)$  at time  $t$  solves the parabolic problem

$$\partial_t \rho = \mathcal{L}^* \rho = -\operatorname{div}(f\rho) + \Delta \rho, \quad t > 0, x \in \mathbb{R}^d.$$

For ergodic SDEs, for any initial condition  $X(0) = X_0$ , as  $t \rightarrow +\infty$ ,

$$\mathbb{E}(\phi(X(t))) = \int_{\mathbb{R}^d} \phi(x) \rho(x, t) dx \longrightarrow \int_{\mathbb{R}^d} \phi(x) d\mu_\infty(x).$$

The invariant measure  $d\mu_\infty(x) \sim \rho_\infty(x)dx$  is a stationary solution ( $\partial_t \rho_\infty = 0$ ) of the Fokker-Plank equation

$$\mathcal{L}^* \rho_\infty = 0.$$

# Plan of the talk

- 1 Order conditions for the invariant measure
- 2 Postprocessed integrators for ergodic SDEs and SPDEs
- 3 Optimal explicit stabilized integrator
- 4 An algebraic framework based on exotic aromatic Butcher-series

# Order conditions for the invariant measure

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- A. Abdule, G. V., K. Zygalakis, *High order numerical approximation of ergodic SDE invariant measures*, *SIAM SINUM*, 2014.
- A. Abdule, G. V., K. Zygalakis, *Long time accuracy of Lie-Trotter splitting methods for Langevin dynamics*, *SIAM SINUM*, 2015.

# Asymptotic expansions

## Theorem (Talay and Tubaro, 1990, see also, Milstein, Tretyakov)

Assume that  $X_n \mapsto X_{n+1}$  (weak order  $p$ ) is *ergodic* and has a Taylor expansion  $\mathbb{E}(\phi(X_1))|X_0 = x = \phi(x) + h\mathcal{L}\phi + h^2A_1\phi + h^3A_2\phi + \dots$ . If  $\mu_\infty^h$  denotes the numerical invariant distribution, then

$$e(\phi, h) = \int_{\mathbb{R}^d} \phi d\mu_\infty^h - \int_{\mathbb{R}^d} \phi d\mu_\infty = \lambda_p h^p + \mathcal{O}(h^{p+1}),$$

$$\mathbb{E}(\phi(X_n)) - \int_{\mathbb{R}^d} \phi d\mu_\infty - \lambda_p h^p = \mathcal{O}(\exp(-cnh) + h^{p+1}),$$

where, denoting  $u(t, x) = \mathbb{E}\phi(X(t, x))$ ,

$$\begin{aligned}\lambda_p &= \int_0^{+\infty} \int_{\mathbb{R}^d} \left( A_p - \frac{\mathcal{L}^{p+1}}{(p+1)!} \right) u(t, x) \rho_\infty(x) dx dt \\ &= \int_0^{+\infty} \int_{\mathbb{R}^d} u(t, x) (A_p)^* \rho_\infty(x) dx dt.\end{aligned}$$

## High order approximation of the numerical invariant measure

Assume that  $X_n \mapsto X_{n+1}$  is **ergodic** with standard assumptions and

$$\mathbb{E}(\phi(X_1)|X_0 = x) = \phi(x) + h\mathcal{L}\phi + h^2A_1\phi + h^3A_2\phi + \dots$$

### Standard weak order condition.

If  $A_j = \frac{\mathcal{L}^j}{j!}$ ,  $1 \leq j < p$ , then (weak order  $p$ )

$$\mathbb{E}(\phi(X(t_n))) = \mathbb{E}(\phi(X_n)) + \mathcal{O}(h^p), \quad t_n = nh \leq T.$$

### Order condition for the invariant measure.

If  $A_j^* \rho_\infty = 0$ ,  $1 \leq j < p$ , then (order  $p$  for the invariant measure)

$$\lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=1}^N \phi(X_n) = \int_{\mathbb{R}^d} \phi(y) d\mu(y) + \mathcal{O}(h^p),$$
$$\mathbb{E}(\phi(X_n)) - \int_{\mathbb{R}^d} \phi d\mu_\infty = \mathcal{O}(\exp(-cnh) + h^p).$$

## Application: high order integrator based on modified equations

It is possible to construct integrators of weak order 1 that have order  $p$  for the invariant measure.

This can be done inspired by recent advances in modified equations of SDEs (see Shardlow 2006, Zygalakis, 2011, Debussche & Faou, 2011, Abdulle Cohen, V., Zygalakis, 2013).

### Theorem (Abdulle, V., Zygalakis)

Consider an ergodic integrator  $X_n \mapsto X_{n+1}$  (with weak order  $\geq 1$ ) for an ergodic SDE in the torus  $\mathbb{T}^d$  (with technical assumptions),

$$dX = f(X)dt + g(X)dW.$$

Then, for all  $p \geq 1$ , there exist a modified equations

$$dX = (f + hf_1 + \dots + h^{p-1}f_{p-1})(X)dt + g(X)dW,$$

such that the integrator applied to this modified equation has order  $p$  for the invariant measure of the original system  $dX = fdt + gdw$  (assuming ergodicity).



## Example of high order integrator for the invariant measure

### Theorem (Abdulle, V., Zygalkakis)

Consider the Euler-Maruyama scheme  $X_{n+1} = X_n + hf(X_n) + \sigma\Delta W_n$  applied to Brownian dynamics ( $f = -\nabla V$ ).

Then, the Euler-Maruyama scheme applied to **the modified SDE**

$$dX = (f + hf_1 + h^2 f_2)dt + \sigma\Delta W_n$$

$$f_1 = -\frac{1}{2}f'f - \frac{\sigma^2}{4}\Delta f,$$

$$f_2 = -\frac{1}{2}f'f'f - \frac{1}{6}f''(f, f) - \frac{1}{3}\sigma^2 \sum_{i=1}^d f''(e_i, f'e_i) - \frac{1}{4}\sigma^2 f' \Delta f,$$

has **order 3 for the invariant measure** (assuming ergodicity).

Remark 1: the **weak order** of accuracy is only 1.

Remark 2: **derivative free** versions can also be constructed.

# Postprocessed integrators for ergodic SDEs and SPDEs

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- G. V., *Postprocessed integrators for the high order integration of ergodic SDEs*, *SIAM SISC*, 2015.
- C.-E. Bréhier and G. V., *High-order integrator for sampling the invariant distribution of a class of parabolic SPDEs with additive space-time noise*, *SIAM SISC*, 2016.

# Postprocessed integrators for ergodic SDEs

Idea: extend to the context of ergodic SDEs the popular idea of effective order for ODEs from Butcher 69',

$$y_{n+1} = \chi_h \circ K_h \circ \chi_h^{-1}(y_n), \quad y_n = \chi_h \circ K_h^n \circ \chi_h^{-1}(y_0).$$

## Example based on the Euler-Maruyama method

for Brownian dynamics:  $dX(t) = -\nabla V(X(t))dt + \sigma dW(t)$ .

$$X_{n+1} = X_n - h \nabla V \left( X_n + \frac{1}{2} \sigma \sqrt{h} \xi_n \right) + \sigma \sqrt{h} \xi_n, \quad \bar{X}_n = X_n + \frac{1}{2} \sigma \sqrt{h} \xi_n.$$

$X_n$  has order 1 of accuracy for the invariant measure.

$\bar{X}_n$  has order 2 of accuracy for the invariant measure (postprocessor).

This method was first derived as a [non-Markovian method](#) by [Leimkhuler, Matthews, 2013], see [Leimkhuler, Matthews, Tretyakov, 2014],

$$\bar{X}_{n+1} = \bar{X}_n + hf(\bar{X}_n) + \frac{1}{2} \sigma \sqrt{h} (\xi_n + \xi_{n+1}).$$

# Postprocessed integrators

Postprocessing:  $\bar{X}_n = G_n(X_n)$ , with weak Taylor series expansion

$$\mathbb{E}(\phi(G_n(x))) = \phi(x) + h^p \bar{A}_p \phi(x) + \mathcal{O}(h^{p+1}).$$

## Theorem (V.)

*Under technical assumptions, assume that  $X_n \mapsto X_{n+1}$  and  $\bar{X}_n$  satisfy*

$$A_j^* \rho_\infty = 0 \quad j < p, \quad (\text{order } p \text{ for the invariant measure}),$$

$$\text{and} \quad (A_p + [\mathcal{L}, \bar{A}_p])^* \rho_\infty = (A_p + \mathcal{L} \bar{A}_p - \bar{A}_p \mathcal{L})^* \rho_\infty = 0,$$

*then (order  $p + 1$  for the invariant measure)*

$$\mathbb{E}(\phi(\bar{X}_n)) - \int_{\mathbb{R}^d} \phi d\mu_\infty = \mathcal{O}(\exp(-cnh) + h^{p+1}).$$

**Remark:** the postprocessing is needed only at the end of the time interval (not at each time step).

# New schemes based on the theta method

We introduce a modification of the  $\theta = 1$  method:

$$X_{n+1} = X_n - h\nabla V(X_{n+1} + a\sigma\sqrt{h}\xi_n) + \sigma\sqrt{h}\xi_n, \quad a = -\frac{1}{2} + \frac{\sqrt{2}}{2},$$

## A postprocessor of order 2

$$\bar{X}_n = X_n + c\sigma\sqrt{h}J_n^{-1}\xi_n, \quad c = \sqrt{2\sqrt{2} - 1}/2$$

The matrix  $J_n^{-1}$  is the inverse of  $J_n = I - hf'(X_n + a\sigma\sqrt{h}\xi_{n-1})$ .

## A postprocessor of order 2 (order 3 for linear problems)

$$\bar{X}_n = X_n - hb\nabla V(\bar{X}_n) + c\sigma\sqrt{h}\xi_n, \quad b = \sqrt{2}/2, \quad c = \sqrt{4\sqrt{2} - 1}/2.$$

# The SPDE case: the linear implicit Euler scheme

Stochastic evolution equation on the Hilbert space  $H$ :

$$du(t) = Au(t)dt + F(u(t))dt + dW^Q(t) \quad , \quad u(0) = u_0 \in H.$$

Euler scheme, with time-step size  $h$ :

$$\begin{aligned} v_{n+1} &= v_n + hAv_{n+1} + hF(v_n) + \sqrt{h}\xi_n^Q \\ &= J_1 v_n + hJ_1 F(v_n) + \sqrt{h}J_1 \xi_n^Q, \end{aligned}$$

where  $J_1 = (I - hA)^{-1}$  and  $\sqrt{h}\xi_n^Q = W^Q((n+1)h) - W^Q(nh)$ .

Order of convergence is  $\bar{s} - \varepsilon$  for all  $\varepsilon > 0$  (see Bréhier 2014):

$$\bar{s} = \sup \left\{ s \in (0, 1) ; \operatorname{Trace} \left( (-A)^{-1+s} Q \right) < +\infty \right\} > 0.$$

Example: for  $A = \frac{\partial^2}{\partial x^2}$ ,  $Q = I$  in dimension 1, we have  $\bar{s} = 1/2$ .

# The postprocessed scheme

Linear Euler scheme:

$$v_{n+1} = J_1 \left( v_n + hF(v_n) + \sqrt{h} \xi_n^Q \right).$$

## New postprocessed scheme

$$u_{n+1} = J_1 \left( u_n + hF \left( u_n + \frac{1}{2} \sqrt{h} J_2 \xi_n^Q \right) + \sqrt{h} \xi_n^Q \right)$$

Postprocessing:  $\bar{u}_n = u_n + \frac{1}{2} J_3 \sqrt{h} \xi_n^Q$ ,

with

$$J_1 = (I - hA)^{-1}, \quad J_2 = \left( I - \frac{3 - \sqrt{2}}{2} hA \right)^{-1}, \quad J_3 = \left( I - \frac{h}{2} A \right)^{-1/2}.$$

# Analysis of the postprocessed Euler method

## Theorem (Bréhier, V.)

- The Markov chain  $(u_n, \bar{u}_{n-1})_{n \in \mathbb{N}}$  is ergodic, with unique invariant distribution, and for any test function  $\varphi : H \rightarrow \mathbb{R}$  of class  $\mathcal{C}^2$ , with bounded derivatives,

$$\left| \mathbb{E}(\varphi(\bar{u}_n)) - \int_H \varphi(y) d\bar{\mu}_\infty^h(y) \right| = \mathcal{O} \left( \exp\left(-\frac{(\lambda_1 - L)}{1 + \lambda_1 h} nh\right) \right).$$

- Moreover, for *the case of a linear  $F$* , for any  $s \in (0, \bar{s})$ ,

$$\int_H \varphi(y) d\bar{\mu}_\infty^h(y) - \int_H \varphi(y) d\mu_\infty(y) = \mathcal{O}(h^{s+1}).$$

**Remark:** error for the standard linear Euler:  $\mathcal{O}(h^s)$ ,  $s \in (0, \bar{s})$ .



# Numerical experiments (stochastic heat equation)

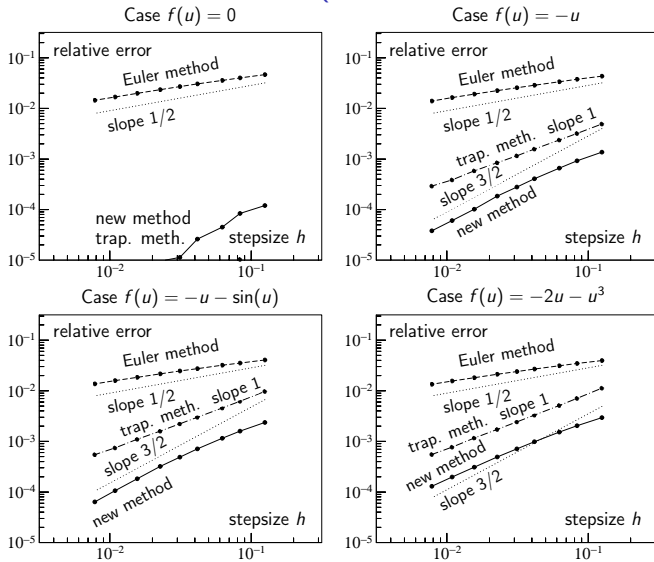


Figure: Orders of convergence, test function  $\varphi(u) = \exp(-\|u\|^2)$ .

# Optimal explicit stabilized integrator for stiff and ergodic SDEs

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A. Abdulle, I. Almuslimani, G. V., *Optimal explicit stabilized integrator of weak order one for stiff and ergodic stochastic differential equations*, [ArXiv, submitted](#), 2017.

See the poster of Ibrahim Almuslimani today at 6 p.m. (Atrium).

# Stability analysis of (deterministic) integrators

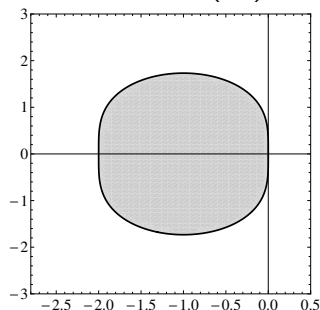
**Stability function.** Consider  $y'(t) = \lambda y(t)$ ,  $y(0) = 1$ .

A Runge-Kutta method with **stepsize**  $h$  yields  $y_{n+1} = R(h\lambda)y_n$ .

Stability domain  $\mathcal{S} := \{z \in \mathbb{C}; |R(z)| \leq 1\}$ .

**Stiff integrators.** If  $\mathbb{C}^- \subset \mathcal{S}$ , the method is called **A-stable**.

If in addition  $R(\infty) = 0$ , the method is called **L-stable**.



**Example: the Heun method (explicit)**

$$y_{n+1} = y_n + \frac{h}{2}f(y_n) + \frac{h}{2}f(y_n + hf(y_n)).$$

$$R(z) = 1 + z + \frac{z^2}{2}.$$

The stability condition  $-2 \leq h\lambda \leq 0$  becomes for diffusion problems  $h\Delta x^{-2} \leq C$  (severe stepsize restriction).

# Example: first order Chebyshev methods

An  $s$ -stage Runge-Kutta method  $y_0 \mapsto y_1$ .

$$K_1 = y_0 + \frac{h}{s^2} F(y_0), \quad K_0 = y_0,$$

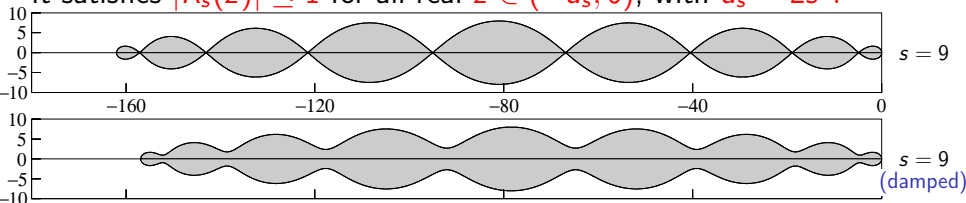
$$K_j = \frac{2h}{s^2} F(K_{j-1}) + 2K_{j-1} - K_{j-2}, \quad j = 2, \dots, s$$

$$y_1 = K_s$$

Stability function given by  $R_s(z) = T_s\left(1 + \frac{z}{s^2}\right)$  where

$T_s(\cos x) = \cos(sx)$  are the **Chebyshev polynomials**.

It satisfies  $|R_s(z)| \leq 1$  for all real  $z \in (-d_s, 0)$ , with  $d_s = 2s^2$ .

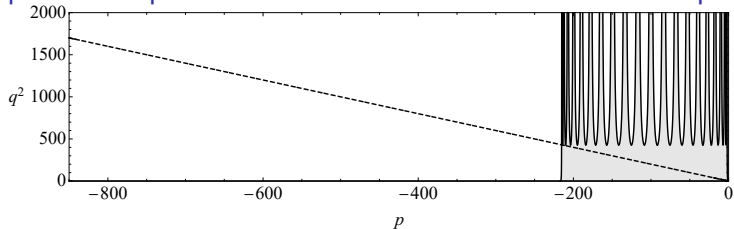


# Explicit stabilized integrators (Chebyshev methods)

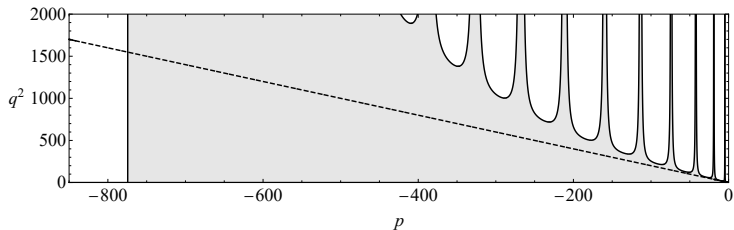
Yuan'Chzao Din (1958), Franklin (1959), Guillou, Lago (1960)...

- 1. RKC: Methods based on three-term recurrence relation (non-optimal) with  $d_s \simeq 0.66 \cdot s^2$   
van der Houwen, Shampine, Sommeijer, Verwer (RKC, IMEX extension IRKC, 1980-2007), Zbinden (PRKC 2011)
- 2. Methods based on composition (no-recurrence relation)  
Bogatyrev, Lebedev, Skvorstov, Medovikov (DUMKA 1976-2004), Jeltsch, Torrilhon 2007
- 3. ROCK methods (close to optimal stability for second order)  
Abdulle, Medovikov (ROCK2 2000-02) with  $d_s \simeq 0.81 \cdot s^2$   
Abdulle (ROCK4 2002-05) with  $d_s \simeq 0.35 \cdot s^2$
- 4. Extension to stiff stochastic problems: S-ROCK methods  
Weak order 1: Abdulle, Cirilli, Li, Hu (S-ROCK 2007-2009,  $\tau$ -ROCK methods 2010) with  $d_s \simeq 0.33 \cdot s^2$   
Weak order 2: Abdulle, Vilmart, Zygalkakis (S-ROCK2 preprint 2012) with  $d_s \simeq 0.43 \cdot s^2$

# New optimal explicit stabilized scheme for MS stiff problems



standard S-ROCK method (Abdulle and Li, 2008,  $s = 20$ ,  $\eta = 6.95$ )  
stability domain size  $d_s \simeq 0.33 \cdot s^2$ .

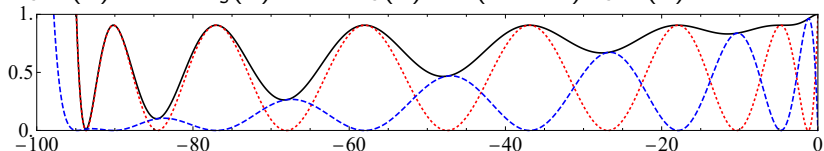


new SK-ROCK method ( $s = 20$ ,  $\eta = 0.05$ )  
stability domain size  $d_s \geq (2 - \frac{4}{3}\eta)s^2$ .

# New optimal explicit stabilized schemes:

Main idea: we use **second kind Chebyshev polynomials**

$U_{s-1}(x) = s^{-1} T'_s(x)$ , and  $T_s(x)^2 + (1 - x^2)U_{s-1}(x)^2 = 1$ .



$\mathbb{E}(|R(p, q, \xi)|^2) = A(p)^2 + B(p)^2 q^2$  as a function of  $p$  ( $q^2 = -2p$ ).

Features of the new **optimal second kind explicit Chebyshev methods**:

- Coincides with the **optimal deterministic Chebyshev method of order one** ( $d_s \geq (2 - \frac{4}{3}\eta) \cdot s^2$ ) for deterministic problems and inherits its optimal stability domain size.
- A **postprocessor of order two** is constructed for Brownian dynamics (for invariant measure sampling).

# An algebraic framework based on exotic aromatic Butcher-series

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A. Laurent, G. V., *Exotic aromatic B-series for the study of long time integrators for a class of ergodic SDEs*, [ArXiv, submitted](#), 2017.

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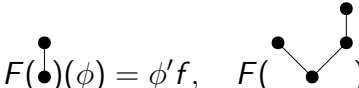


# Aromatic Butcher-series

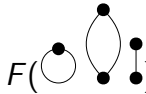
Stochastic case: Tree formalism for **strong and weak errors on finite time**: Burrage K., Burrage P.M., 1996; Komori, Mitsui, Sugiura, 1997; Rößler, 2004/2006, ...

Here we focus of the accuracy for the **invariant measure (long time)**.

We rewrite high-order differentials with trees. We denote  $F(\gamma)(\phi)$  the elementary differential of a tree  $\gamma$ .

$$F(\bullet)(\phi) = \phi, \quad F(\text{•}) (\phi) = \phi' f, \quad F(\text{•}) (\phi) = \phi''(f, f' f)$$


**Aromatic forests**: introduced for deterministic geometric integration by Chartier, Murua, 2007 (See also Bogfjellmo, 2015)

$$F(\text{•}) (\phi) = \text{div}(f) \times \left( \sum \partial_i f_j \partial_j f_i \right) \times \phi' f$$


# New exotic aromatic B-series: using lianas

Grafted aromatic forests: a random vector  $\xi \sim \mathcal{N}(0, I_d)$  is represented by crosses (in the spirit of P-series)

$$F(\text{diagram 1})(\phi) = \phi''(f'\xi, \xi) \quad \text{and} \quad F(\text{diagram 2})(\phi) = \phi' f''(\xi, \xi).$$

We also introduce **lianas** in our forests called **exotic aromatic forests**:

$$F(\text{diagram 3})(\phi) = \sum_i \phi''(f' e_i, e_i) = \mathbb{E}(\phi''(f'\xi, \xi)).$$

$$F(\text{diagram 4})(\phi) = \sum_i \phi''(e_i, e_i) = \Delta\phi = \mathbb{E}(\phi''(\xi, \xi)).$$

$$F(\text{diagram 5})(\phi) = \sum_{i,j} \phi''(e_i, f'''(e_j, e_j, e_i)) = \sum_i \phi''(e_i, (\Delta f)'(e_i)).$$

# Integration by parts using trees: examples

$$\begin{aligned}
 \int_{\mathbb{R}^d} F(\text{tree with 2 nodes, one dashed loop})(\phi) \rho_\infty dy &= \sum_{i,j} \int_{\mathbb{R}^d} \frac{\partial^3 \phi}{\partial x_i \partial x_j \partial x_j} f_i \rho_\infty dy \\
 &= - \sum_{i,j} \left[ \int_{\mathbb{R}^d} \frac{\partial \phi}{\partial x_i \partial x_j} \frac{\partial f_i}{\partial x_j} \rho_\infty dy + \int_{\mathbb{R}^d} \frac{\partial \phi}{\partial x_i \partial x_j} f_i \frac{\partial \rho_\infty}{\partial x_j} dy \right] \\
 &= - \int_{\mathbb{R}^d} F(\text{tree with 2 nodes, one dashed loop})(\phi) \rho_\infty dy - \frac{2}{\sigma^2} \int_{\mathbb{R}^d} F(\text{tree with 3 nodes, V-shape})(\phi) \rho_\infty dy.
 \end{aligned}$$

We obtain:

$$\text{tree with 2 nodes, one dashed loop} \sim - \text{tree with 2 nodes, one solid line} - \frac{2}{\sigma^2} \text{tree with 3 nodes, V-shape}.$$

Remark: the new exotic aromatic B-series satisfy an **isometric equivariance property** (see related work on characterizing **affine equivariant maps** by McLachlan, Modin, Munthe-Kaas, Verdier, 2016)

# Order conditions for the invariant measure

$$Y_i^n = X_n + h \sum_{j=1}^s a_{ij} f(Y_j^n) + d_i \sigma \sqrt{h} \xi_n, \quad i = 1, \dots, s,$$

$$X_{n+1} = X_n + h \sum_{i=1}^s b_i f(Y_i^n) + \sigma \sqrt{h} \xi_n,$$

## Theorem (Laurent, V., Conditions for order $p$ )

Order	Tree $\tau$	$F(\tau)(\phi)$	Order condition
1		$\phi' f$	$\sum b_i = 1$
2		$\phi' f' f$	$\sum b_i c_i - 2 \sum b_i d_i = -\frac{1}{2}$
		$\phi' \Delta f$	$\sum b_i d_i^2 - 2 \sum b_i d_i = -\frac{1}{2}$
3		$\phi' f' f' f$	$\sum b_i a_{ij} c_j - 2 \sum b_i a_{ij} d_j$ $+ \sum b_i c_i - (\sum b_i d_i)^2 = 0$
	...	...	...

# Summary

- Using tools from geometric integration, we presented **new order conditions** for the accuracy of ergodic integrators, with emphasis on **postprocessed integrators**.
- In particular, **high order in the deterministic or weak sense is not necessary** to achieve high order for the invariant measure.
- A new **high-order** method ( $\bar{s} + 1$  instead of  $\bar{s}$  for linearized Euler) for sampling the invariant distribution of parabolic SPDEs

$$du(t) = Au(t)dt + F(u(t))dt + dW^Q(t),$$

(proof in a simplified linear case).

- study of algebraic structures with **exotic aromatic Butcher trees**.

## Current works:

- analysis of the order of convergence in the general semilinear SPDE case.
- combination with Multilevel Monte-Carlo strategies.